Nuclei as Laboratories: Nuclear Tests of Fundamental Symmetries



M.J. Ramsey-Musolf

N. Bell

S. Page

- V. Cirigliano
- J. Erler
- B. Holstein
- A. Kurylov
- C. Lee

C. Maekawa

S. Profumo

G. Prezeau

- S. Tulin
- B. Van Kolck
- P. Vogel
- S. Zhu

Nuclear Science



The mission: Explain the origin, evolution, and structure of the baryonic matter of the Universe

Nuclear Science



Three frontiers:

- Fundamental symmetries & neutrinos
- Nuclei and nuclear astrophysics
- QCD





Puzzles the Standard Model can't solve



What are the new fundamental symmetries?

• Why is there more matter than antimatter in the present universe?

Electric dipole moment searches

• What are the unseen forces that disappeared from view as the universe cooled?

Precision electroweak: weak decays, scattering, LFV

 What are the masses of neutrinos and how have they shaped the evolution of the universe?

Neutrino oscillations, $0\nu\beta\beta$ -decay, θ_{13} , ...

Tribble report

What is the origin of baryonic matter ?



What are the quantitative implications of new EDM experiments for explaining the origin of the baryonic component of the Universe ?

What is the origin of baryonic matter?



EW Baryogenesis: Standard Model

Sakharov:



EW Baryogenesis: Standard Model

Shaposhnikov



- CP-violation too weak
- EW PT too weak

$$J = s_{12} s_{13} s_{23} c_{12} c_{13}^2 c_{23} \sin \delta_{13}$$

= (2.88 ± 0.33) × 10⁻⁵
$$\frac{m_t^4}{M_W^4} \frac{m_b^4}{M_W^4} \frac{m_c^2}{M_W^2} \frac{m_s^2}{M_W^2} \approx 3 \times 10^{-13}$$



Baryogenesis: New Electroweak Physics

90's: Cohen, Kaplan, Nelson Joyce, Prokopec, Turok



EDM Probes of New CP Violation



EDM constraints & SUSY CPV





Weak decays & new physics CKM unitarity ?

$$d \to u e^{-} \overline{v}_{e}$$
$$s \to u e^{-} \overline{v}_{e}$$
$$b \to u e^{-} \overline{v}_{e}$$



R Parity Violation Kurylov, R-M, Su



No long-lived LSP or SUSY DM



Nuclear structure effects?

Weak decays & new physics Correlations

$$d \rightarrow u e^{-} \overline{v}_{e}$$

$$s \rightarrow u e^{-} \overline{v}_{e}$$

$$b \rightarrow u e^{-} \overline{v}_{e}$$

$$\downarrow^{\nu} \qquad \downarrow^{\nu} \qquad$$

ŦL

SUSY

 e^{-}

 \tilde{V}_{e}

 \overline{V}_{c}

 ${ ilde \chi}^0$

 $\tilde{\chi}^{-}$

 \mathcal{V}_{μ}

 μ^{-}

ũ

$$dW \propto 1 + a \frac{\overrightarrow{p_e} \cdot \overrightarrow{p_v}}{E_e E_v} + A \overrightarrow{\sigma_n} \cdot \frac{\overrightarrow{p_e}}{E_e} + \mathbf{L}$$

Non (V-A) x (V-A) interactions: m_e/E

 $\begin{array}{cccc}
V_{us} & V_{ub} \\
V_{cs} & V_{cb} \\
V_{ts} & V_{tb} \\
\end{array} \begin{pmatrix} d \\ s \\ b \\ b \\
\end{array}$

 β -decay at RIA?

Neutrinos ?

Are they their own antiparticles?LFV & LNV ?Why are their masses so small?What is m_v ?Can they have magnetic moments?What is m_v ?Implications of m_v for neutrino interactions ?



0ν $\beta\beta$ - decay probes the charge conjugation properties of the neutrino



$0v \beta\beta$ - decay: heavy particle exchange

$$\frac{M_H}{M_L} \sim \frac{M_W^4 \bar{k}^2}{\Lambda_{\beta\beta}^5 m_{\beta\beta}} \sim O(1)$$
$$m_{\beta\beta} \sim 0.1 \ eV$$
$$\Lambda_{\beta\beta} \sim 1 \ TeV$$
$$\bar{k} \sim 50 \ MeV$$

How do we compute & separate heavy particle exchange effects?



LF and LN: symmetries of the early universe?



LF and LN: symmetries of the early universe?



$0v \beta\beta$ - decay: heavy particle exchange

How do we compute & separate heavy particle exchange effects?





We have a clear separation of scales



Operator classification



Operator classification

$$\mu = M_{WEAK}$$

$$\mathcal{L}(\boldsymbol{q},\boldsymbol{e}) = \frac{G_F^2}{\Lambda_{\beta\beta}} \sum_{j=1}^{14} C_j(\mu) \, \hat{O}_j^{++} \, \overline{e} \Gamma_j e^c + h.c.$$

e.g.

$$\hat{O}_{1+}^{ab} = \overline{q}_L \gamma^\mu \tau^a q_L \ \overline{q}_R \gamma_\mu \tau^b q_R$$

0ν ββ - decay: a = b = +

Operator classification

$$\mu = M_{WEAK}$$

$$\hat{O}_{1+}^{ab} = \overline{q}_L \gamma^\mu \tau^a q_L \ \overline{q}_R \gamma_\mu \tau^b q_R$$

Chiral transformations: $SU(2)_L \times SU(2)_R$

$$\begin{array}{ccc} q_L \to L \, q_L & L \\ q_R \to R \, q_R & R \end{array} = \exp \left(\begin{array}{ccc} \mathbf{R} & \mathbf{T} \\ i \, \theta_L & \mathbf{T} \\ R & \mathbf{T} \end{array} \right)$$

$$\hat{O}_{1+}^{ab} \in (\mathcal{B}_L, \mathcal{B}_R)$$

Parity transformations: q_L \$ q_R

0ν ββ - decay: a = b = +

$$\hat{O}_{1+}^{++} \leftrightarrow \hat{O}_{1+}^{++}$$



An open question

Is the power counting of operators sufficient to understand weak matrix elements in nuclei ?



An open question

Is the power counting of operators sufficient to understand weak matrix elements in nuclei ?

$$\begin{array}{c} \mathbf{L} = \mathbf{0}, \mathbf{K}, \mathbf{9} \\ \hline D_{0} \mathbf{\nu} \mathbf{\beta} \mathbf{\beta} \\ \mathbf{M}_{fi} \sim p^{0} \\ \mathbf{L} = \mathbf{L}' = \mathbf{0} \\ \mathbf{M}_{fi} \sim p^{2} \\ \mathbf{L} = 2, \mathbf{L}' = \mathbf{0} \\ \hline D_{0} \mathbf{\mu} \mathbf{\beta} \mathbf{\beta} \\ \hline D_{fi} \sim p^{2} \\ \mathbf{L} = 0, \mathbf{L}' = 2 \\ \hline D_{0} \mathbf{\mu} \mathbf{\beta} \mathbf{\beta} \\ \hline D_{fi} \sim p^{4} \\ \mathbf{L} = 4, \mathbf{L}' = \mathbf{0} \\ \hline D_{0} \mathbf{\mu} \mathbf{\beta} \mathbf{\beta} \\ \hline D_{0} \mathbf{\mu} \mathbf{\beta} \mathbf{\beta} \\ \mathbf{M}_{fi} \approx p^{4} \\ \hline \mathbf{L} = 4, \mathbf{L}' = \mathbf{0} \\ \hline D_{0} \mathbf{\mu} \mathbf{\beta} \mathbf{\beta} \\ \mathbf{M}_{0} \mathbf{\mu} \mathbf{\beta} \mathbf{\beta} \\ \mathbf{M}_{0} \mathbf{\beta} \\ \mathbf{M}_{0} \mathbf{\beta} \mathbf{\beta} \\ \mathbf{M}_{0} \mathbf{\beta} \mathbf{\beta} \\ \mathbf{M}_{0} \mathbf{\beta} \\ \mathbf{M}_{0} \mathbf{\beta} \\ \mathbf{M}_{0} \mathbf{\beta} \mathbf{\beta} \\ \mathbf{M}_{0} \mathbf{\beta} \mathbf{\beta} \\ \mathbf{M}_{0} \mathbf{\beta} \mathbf{\beta} \\ \mathbf{M}_{0} \mathbf{\beta} \\ \mathbf{M}_{0} \mathbf{\beta} \mathbf{\beta} \\ \mathbf{M}_{0} \mathbf{\beta} \\ \mathbf{M}_{0} \mathbf{\beta} \mathbf{\beta} \\ \mathbf{M}_{0} \mathbf{\beta} \\ \mathbf{$$

naive

An open question

Complications:

- Bound state wavefunctions (*e.g., h.o.*) don't obey simple power counting
- Configuration mixing is important in heavy nuclei

Is the power counting of operators sufficient to understand weak matrix elements in nuclei ?

- More theoretical study required (RIA)
- Hadronic PV may provide an empirical test



The weak qq force is short range



Use parity-violation to filter out EM & strong interactions

Desplanques, Donoghue, &Holstein (DDH)







N



Hadronic PV: Effective Field Theory

PV Potential







Long Range

 $k_{\pi NN}^{1a}$

O(p)

 $h^1_{\pi NN}$

 $O(p^{-1})$

Short Range

 $\lambda_s^{1,2,3}, \lambda_t, \rho_t$

O(p)

Medium Range



O(p)

A program of few-body measurements

Pionless theory

Ab initio few-body calcs



A program of few-body measurements

Complete determination of PV NN & γ NN interactions through O(p)

Attempt to understand the λ^{i} , h_{π} etc. from QCD

Are the PV LEC's "natural"?

Attempt to understand *nuclear* PV observables systematically

Does EFT power counting work in nuclei ?

Hadronic PV in n-rich nuclei ?

Hadronic PV as a probe



- Determine V_{PV} through O(p)from PV low-energy *few-body* studies where power counting works
- Re-analyze *nuclear* PV observables using this V_{PV}



Conclusions

- Nuclei provide unique and powerful laboratories in which to probe the fundamental symmetries of the early universe
- RIA will provide opportunities to carry out new and complementary experiments whose impact can live on well into the LHC era
- A number of theoretical challenges remain to be addressed at the level of field theory, QCD, and nuclear structure
- New experimental and theoretical efforts in nuclear structure physics are a key component of this quest