



*Pion and Kaon Structure at an EIC
Argonne National Lab, June 2, 2017*

π/K structure & sea quarks in the proton: constraints from chiral EFT & global QCD analysis

Wally Melnitchouk



Collaborators

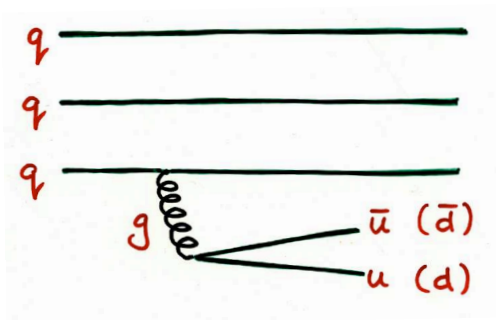
- Patrick Barry (North Carolina State U)
- Tim Hobbs (U Washington)
- Chueng Ji (North Carolina State U)
- Josh McKinney (U North Carolina)
- Yusupujiang Salamu (IHEP Beijing)
- Nobuo Sato (U Connecticut / JLab)
- Tony Thomas (U Adelaide)
- Ping Wang (IHEP Beijing)
- Xuangong Wang (U Adelaide)

Outline

- Motivation: $\bar{d} - \bar{u}$ asymmetry
- PDF constraints from chiral symmetry in QCD / chiral EFT
- Tagged neutron DIS — implications for pion models and pion PDF extraction
- Strange quarks and hyperon production
- Outlook

Light quark sea

- From text-books: perturbative QCD expected to generate symmetric $q\bar{q}$ sea via gluon radiation into $q\bar{q}$ pairs

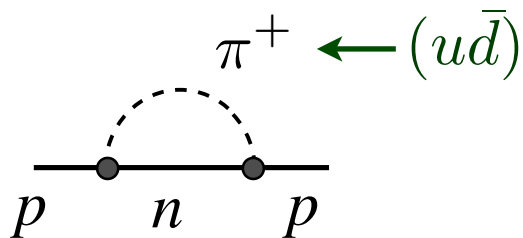


→ since u and d quarks nearly degenerate, expect flavor-symmetric light-quark sea

$$\bar{d} \approx \bar{u}$$

Ross, Sachrajda (1979)

- (Almost) from text-books: Thomas suggested that chiral symmetry of QCD (“low energy”) should have consequences for antiquark PDFs in the nucleon (“high energy”)

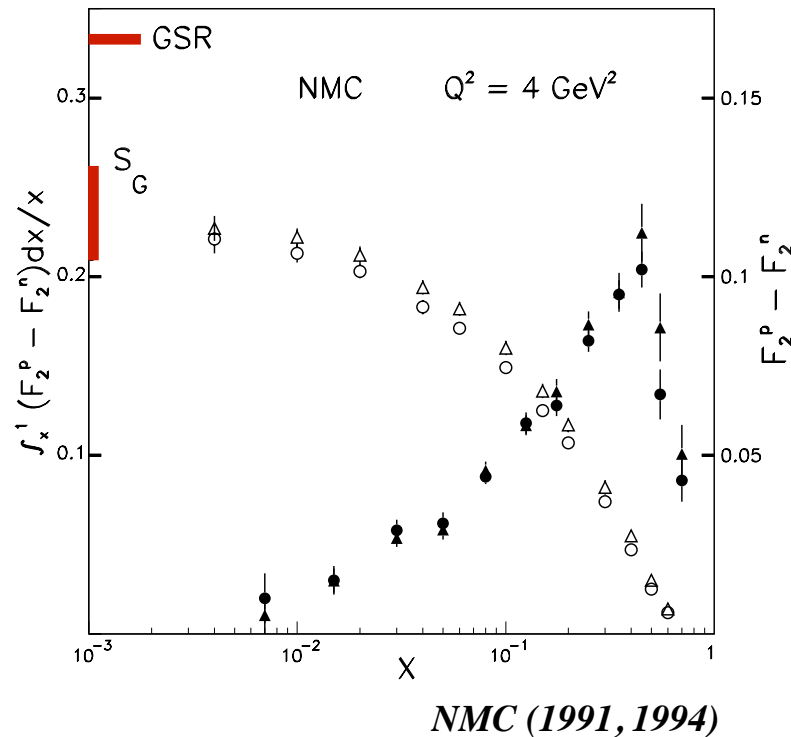


$$\rightarrow \bar{d} > \bar{u}$$

Thomas (1984)

Light quark sea

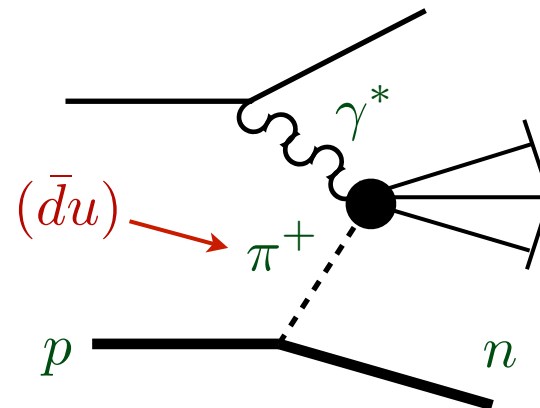
- First clear experimental support for $\bar{d} \neq \bar{u}$ came from measurement of Gottfried sum observed by NMC



$$\int_0^1 \frac{dx}{x} (F_2^p - F_2^n) = \frac{1}{3} - \frac{2}{3} \int_0^1 dx (\bar{d} - \bar{u}) = 0.235(26)$$

→ violation of “Gottfried sum rule”!

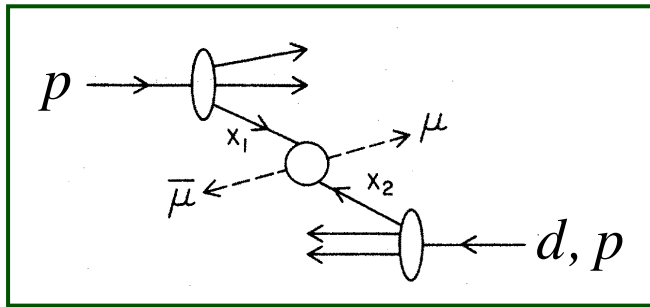
→ Sullivan process — DIS from pion cloud of the nucleon



Sullivan (1972)

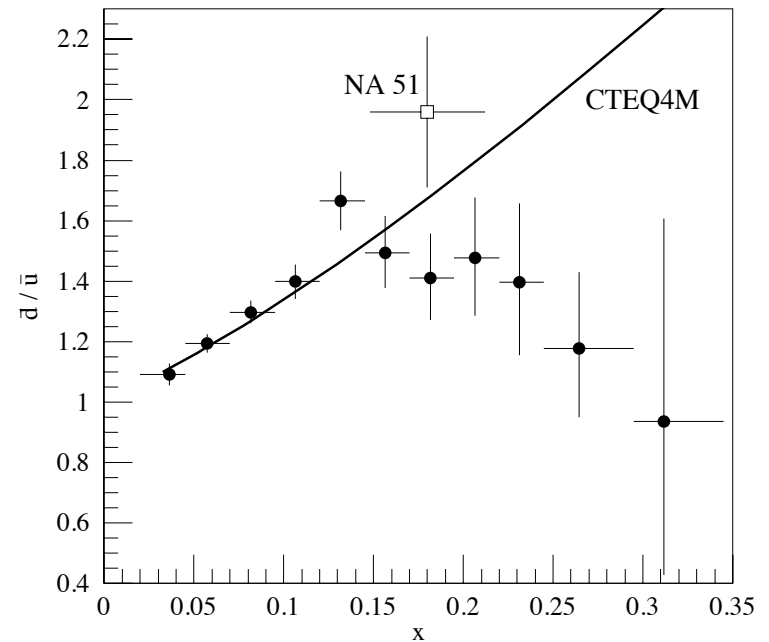
Light quark sea

- x dependence of $\bar{d} - \bar{u}$ asymmetry established in Fermilab E866 pp/pd Drell-Yan experiment



$$\frac{d\sigma}{dx_1 dx_2} \sim \sum_q e_q^2 q(x_1) \bar{q}(x_2) + (x_1 \leftrightarrow x_2)$$

$$\frac{\sigma^{pd}}{\sigma^{pp}} \approx 1 + \frac{\bar{d}(x_2)}{\bar{u}(x_2)} \quad \text{for } x_1 \gg x_2$$

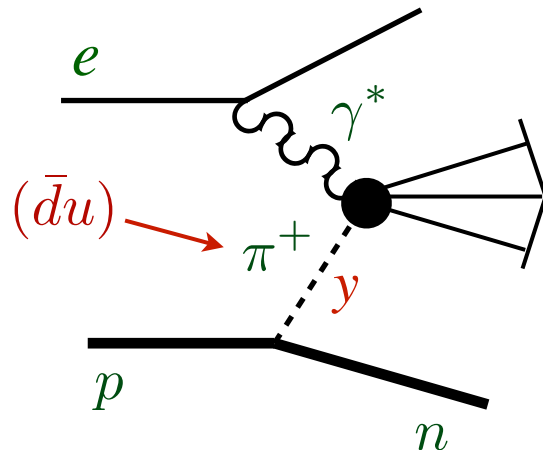


E866 (2001)

- strong enhancement of \bar{d} at $x \sim 0.1 - 0.2$
- intriguing behavior at large x hinting at possible sign change of $\bar{d} - \bar{u}$

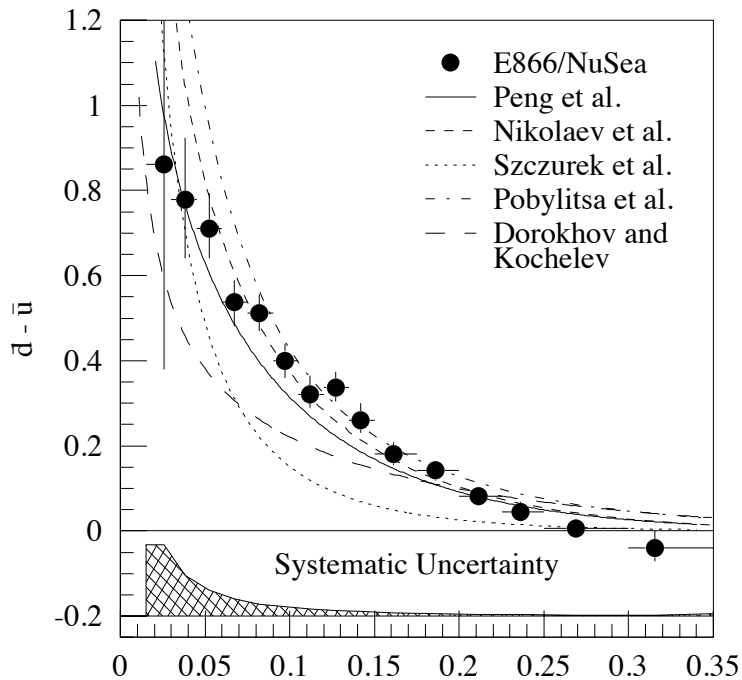
Light quark sea

■ General agreement with pion loop models



$$(\bar{d} - \bar{u})(x) = \int_x^1 \frac{dy}{y} f_{\pi+n}(y) \bar{q}_v^\pi(x/y)$$

$p \rightarrow \pi+n$ splitting function
("flux factor")



→ shape qualitatively reproduced by most models (except at high x),
— but is there a direct connection with QCD?

Chiral effective field theory

■ Rigorous connection with QCD established via chiral EFT

$$\mathcal{L}_{\text{eff}} = \frac{g_A}{2f_\pi} \bar{\psi}_N \gamma^\mu \gamma_5 \vec{\tau} \cdot \partial_\mu \vec{\pi} \psi_N - \frac{1}{(2f_\pi)^2} \bar{\psi}_N \gamma^\mu \vec{\tau} \cdot (\vec{\pi} \times \partial_\mu \vec{\pi}) \psi_N$$

Weinberg (1967)

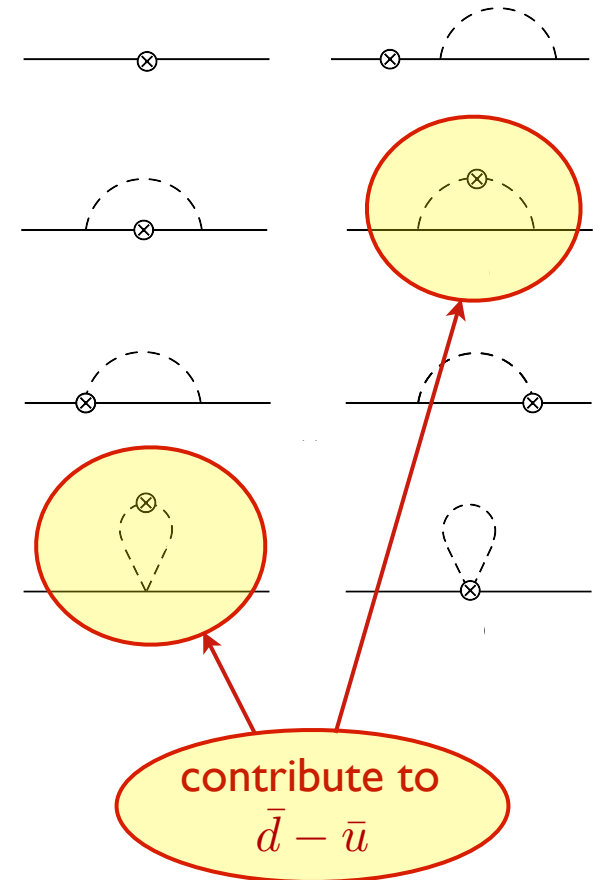
→ lowest order πN interaction includes pion rainbow and tadpole contributions

→ matching quark- and hadron-level operators

$$\mathcal{O}_q^{\mu_1 \dots \mu_n} = \sum_h c_{q/h}^{(n)} \mathcal{O}_h^{\mu_1 \dots \mu_n}$$

yields convolution representation

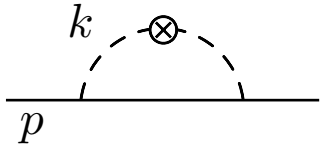
$$q(x) = \sum_h \int_x^1 \frac{dy}{y} f_h(y) q_v^h(x/y)$$



Ji, WM, Thomas (2013)

Chiral effective field theory

- Splitting functions for various diagrams computed in chiral theory
e.g. pion rainbow diagram



$$f_\pi(y) = f^{(\text{on})}(y) + f^{(\delta)}(y)$$

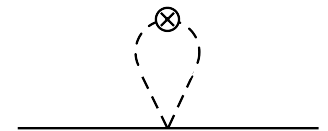
has on-shell ($y = k^+/p^+ > 0$)
and $\delta(y)$ contributions!

$$f^{(\text{on})}(y) = \frac{g_A^2 M^2}{(4\pi f_\pi)^2} \int dk_\perp^2 \frac{y(k_\perp^2 + y^2 M^2)}{[k_\perp^2 + y^2 M^2 + (1-y)m_\pi^2]^2} \mathcal{F}^2$$

$$f^{(\delta)}(y) = \frac{g_A^2}{4(4\pi f_\pi)^2} \int dk_\perp^2 \log\left(\frac{k_\perp^2 + m_\pi^2}{\mu^2}\right) \delta(y) \mathcal{F}^2$$

- Bubble diagram contributes only at $y=0$ (hence $x=0$)

$$f^{(\text{bub})}(y) = \frac{8}{g_A^2} f^{(\delta)}(y)$$



→ contributes to lowest moment, but not at $x > 0$

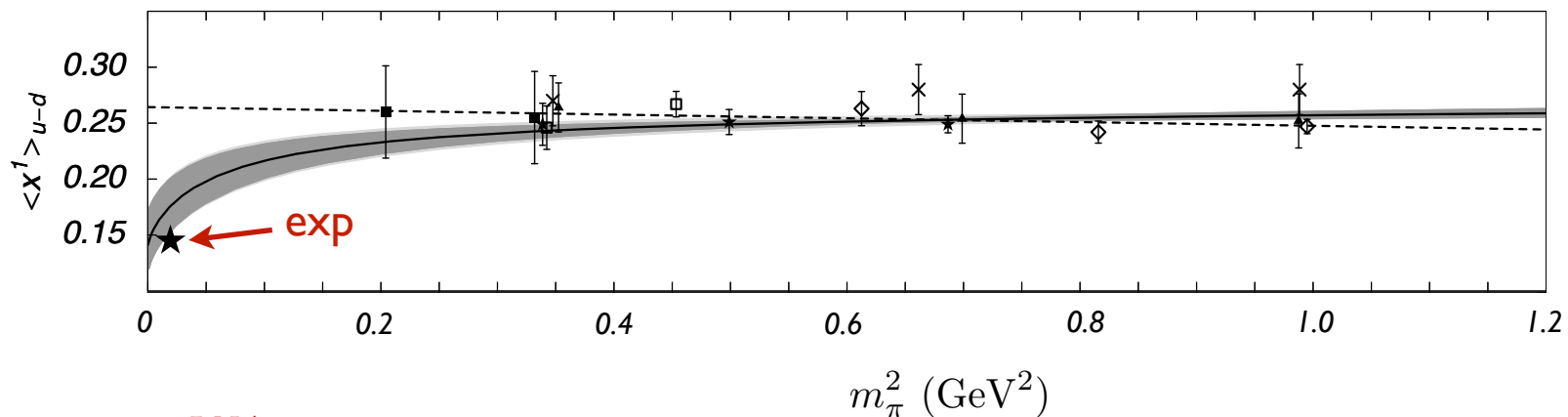
Chiral effective field theory

- Expand moments of PDFs (splitting functions) in powers of m_π
 - coefficients of leading nonanalytic (LNA) terms, reflecting infrared behavior, are *model-independent!*
 - QCD therefore *predicts* a nonzero asymmetry from π loops

$$\int_0^1 dx (\bar{d} - \bar{u}) = \frac{(3g_A^2 - 1)}{(4\pi f_\pi)^2} m_\pi^2 \log(m_\pi^2/\mu^2) + \text{analytic in } m_\pi^2$$

Thomas, WM, Steffens (2000)

- nonanalytic behavior vital for understanding lattice data on PDF moments at low m_π



$$\langle x \rangle_{u-d}^{\text{LNA}} \sim m_\pi^2 \log m_\pi^2$$

Detmold, WM, Thomas (2001)

Chiral effective field theory

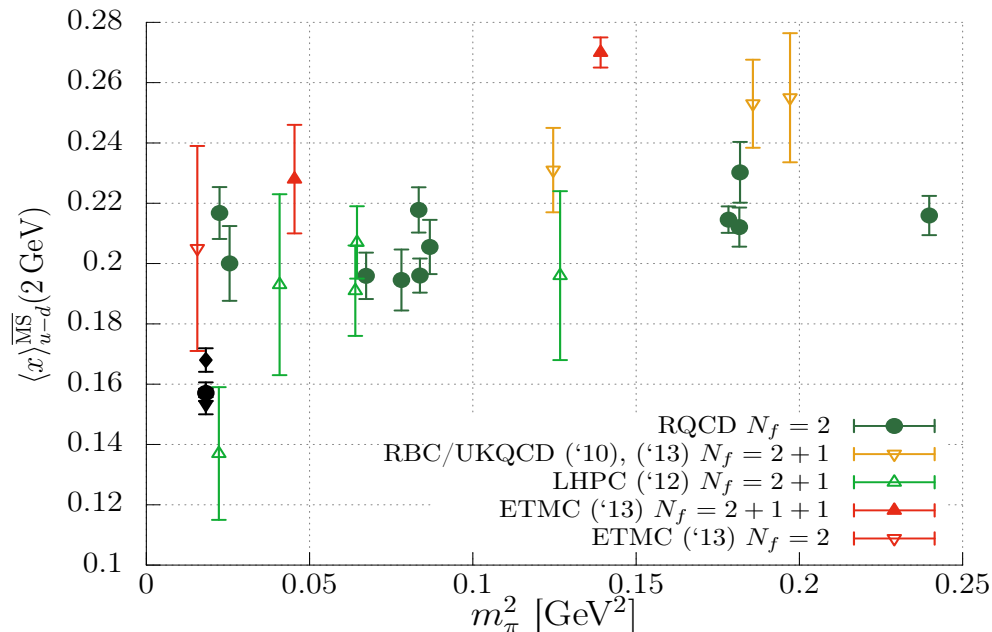
■ Expand moments of PDFs in powers of m_π

- coefficients of leading nonanalytic (LNA) terms, reflecting infrared behavior, are *model-independent!*
- QCD therefore *predicts* a nonzero asymmetry from π loops

$$\int_0^1 dx (\bar{d} - \bar{u}) = \frac{(3g_A^2 - 1)}{(4\pi f_\pi)^2} m_\pi^2 \log(m_\pi^2/\mu^2) + \text{analytic in } m_\pi^2$$

Thomas, WM, Steffens (2000)

- nonanalytic behavior vital for understanding lattice data on PDF moments at low m_π



Bali et al. (2014)

Chiral effective field theory

■ For point-like nucleons and pions, integrals divergent

→ finite size of nucleon provides natural regularization scale
(but does not prescribe form of regularization)

$$\mathcal{F} = \Theta(\Lambda^2 - k_{\perp}^2) \quad k_{\perp} \text{ cutoff}$$

$$\mathcal{F} = \left(\frac{\Lambda^2 - m_{\pi}^2}{\Lambda^2 - t} \right) \quad t \text{ monopole}$$

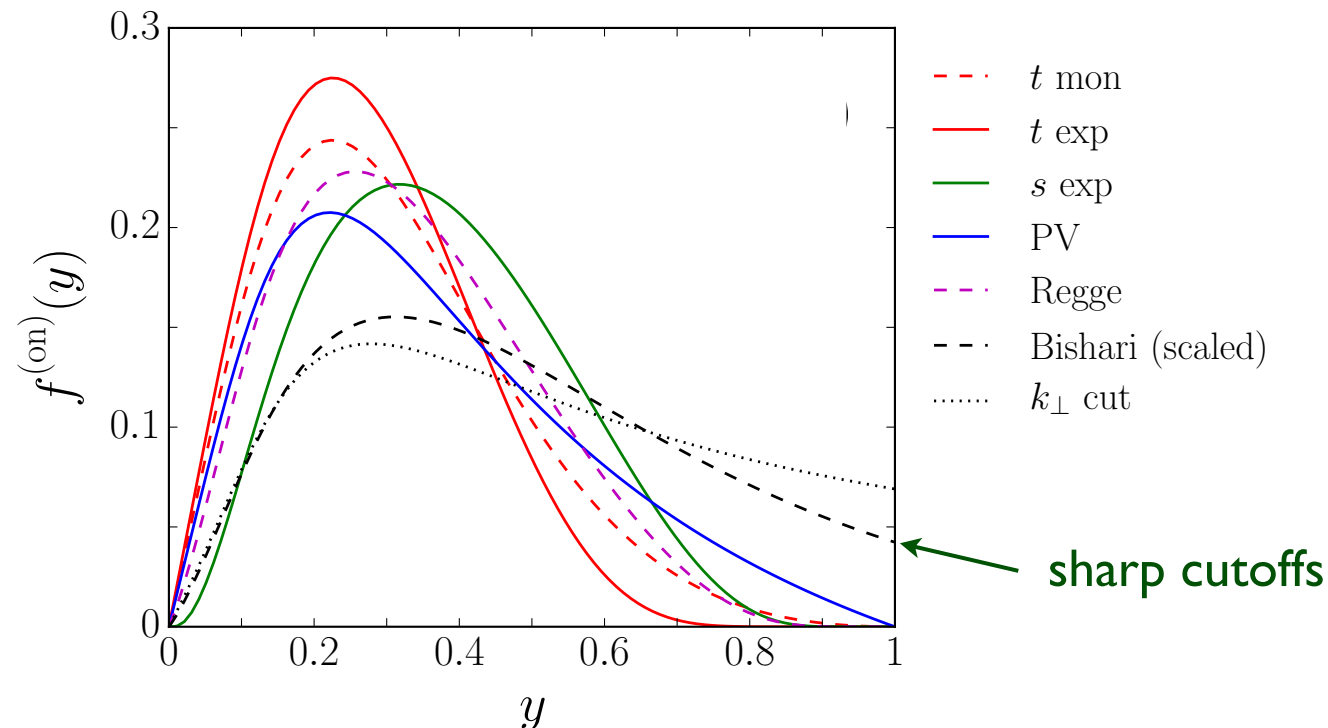
$$\mathcal{F} = \exp \left[(t - m_{\pi}^2) / \Lambda^2 \right] \quad t \text{ exponential}$$

$$\mathcal{F} = \exp \left[(M^2 - s) / \Lambda^2 \right] \quad s\text{-dep. exponential}$$

$$\mathcal{F} = \left[1 - \frac{(t - m_{\pi}^2)^2}{(t - \Lambda^2)^2} \right]^{1/2} \quad \text{Pauli-Villars}$$

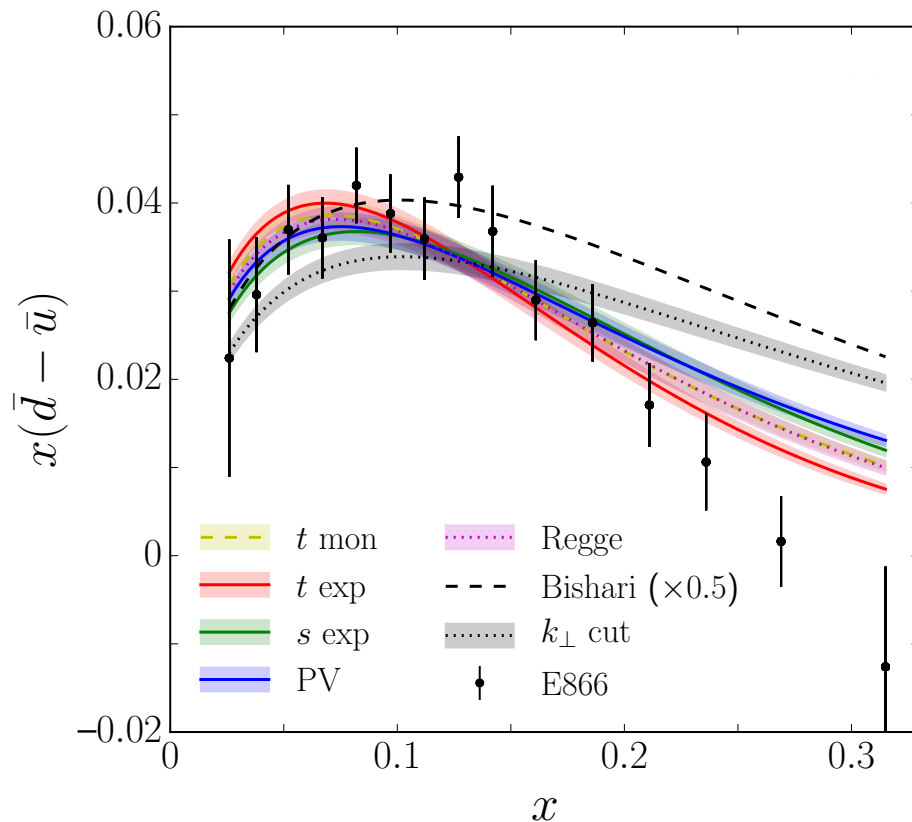
$$\mathcal{F} = y^{-\alpha_{\pi}(t)} \exp \left[(t - m_{\pi}^2) / \Lambda^2 \right] \quad \text{Regge}$$

e.g. on-shell
function



Chiral effective field theory

- E866 $\bar{d} - \bar{u}$ data can be fitted with range of regulators



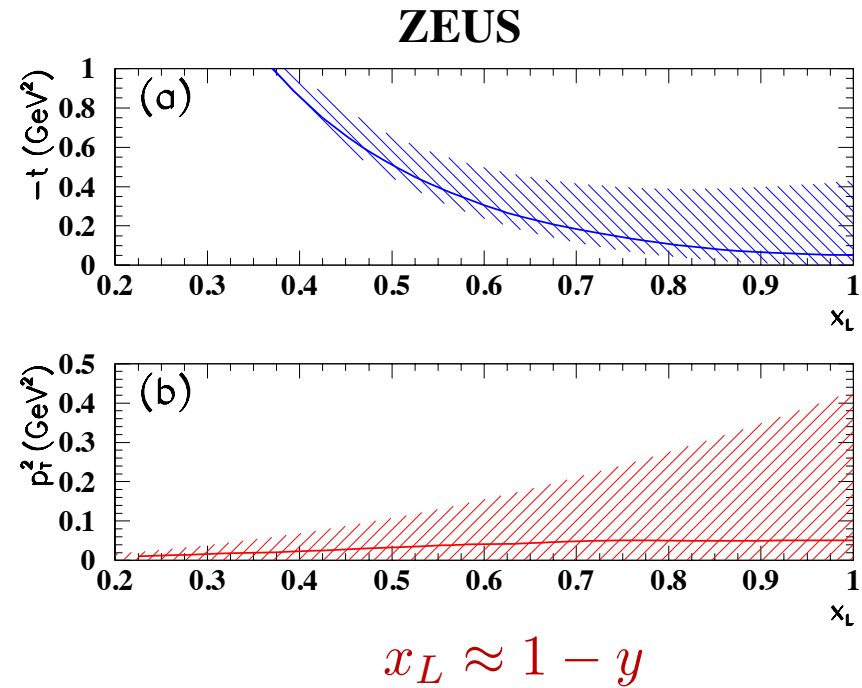
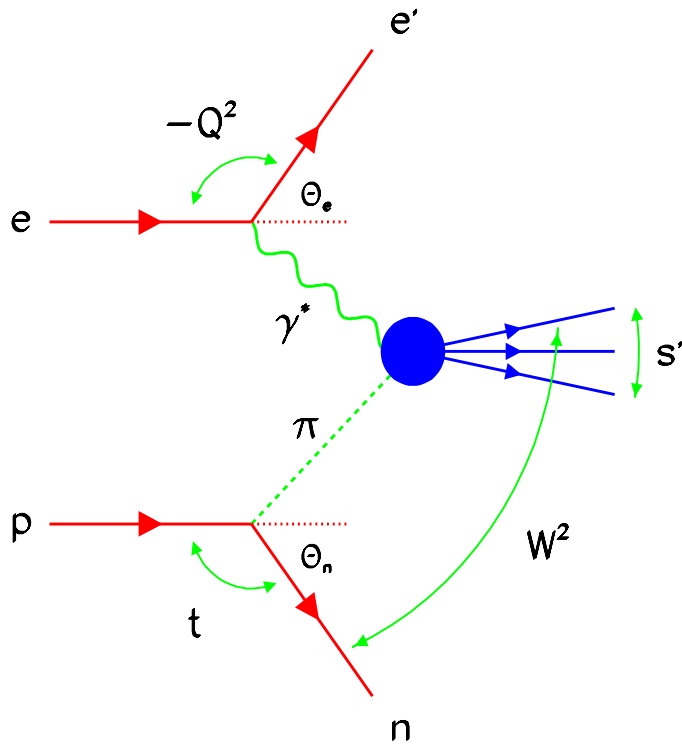
average pion “multiplicity”

$$\langle n \rangle_{\pi N} = 3 \int_0^1 dy f_N^{(\text{on})}(y) \\ \sim 0.25 - 0.3$$

- with exception of k_{\perp} cutoff and Bishari models, all others give reasonable fits, $\chi^2 \lesssim 1.5$
- are there other data that can be more discriminating?

Leading neutron production at HERA

- ZEUS & H1 collaborations measured spectra of neutrons produced at very forward angles, $\theta_n < 0.8$ mrad



- can data be described within same framework as E866 asymmetry?
- simultaneous fit never previously been performed!

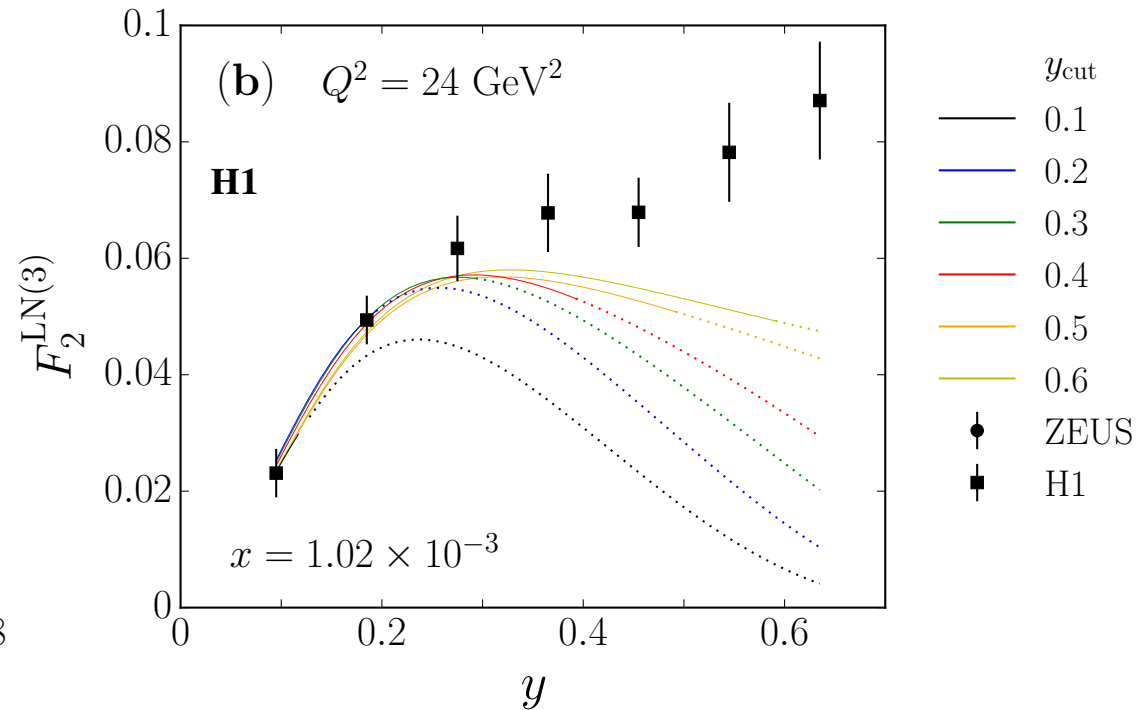
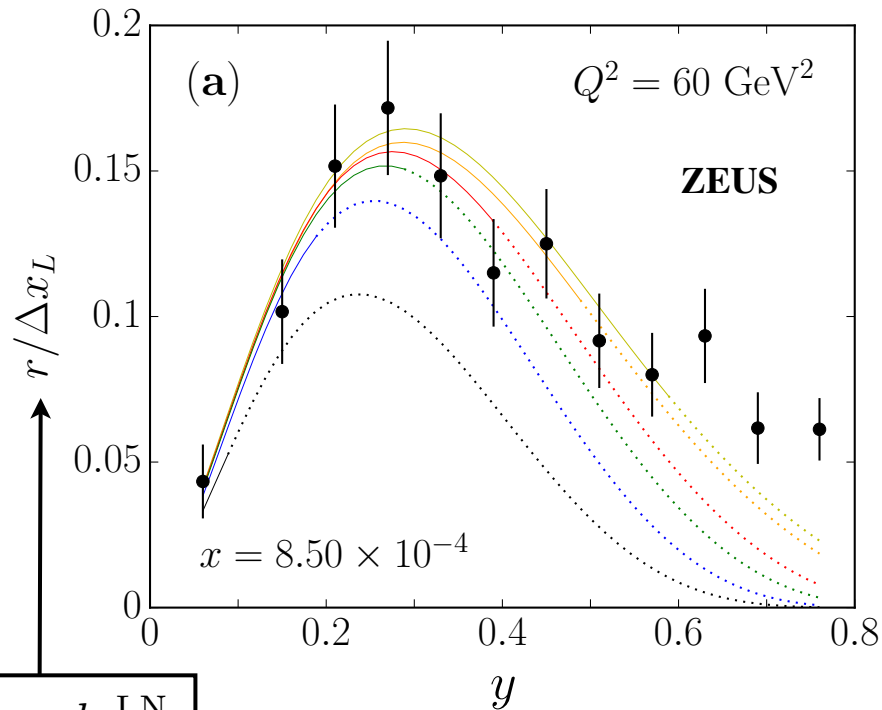
Leading neutron production at HERA

- Measured LN differential cross section (integrated over p_{\perp})

$$\frac{d^3\sigma^{\text{LN}}}{dx dQ^2 dy} \sim F_2^{\text{LN}(3)}(x, Q^2, y)$$

$$2f_N^{(\text{on})}(y) F_2^{\pi}(x/y, Q^2) \text{ for } \pi \text{ exchange}$$

e.g.

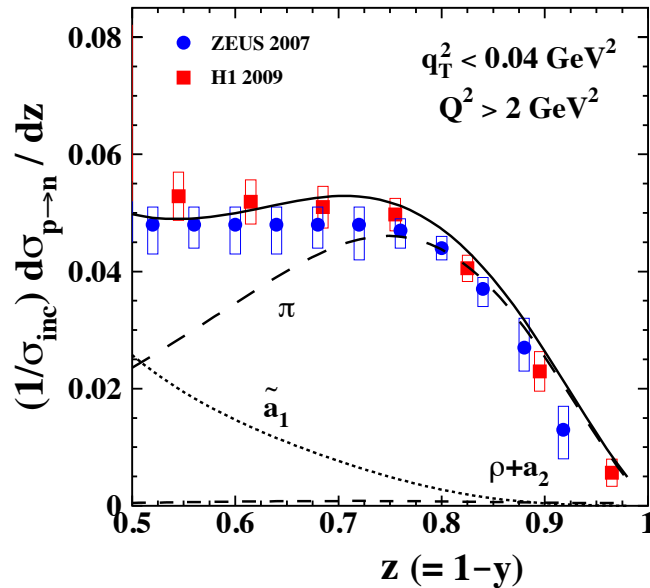


$$r = \frac{d\sigma^{\text{LN}}}{d\sigma^{\text{inc}}}$$

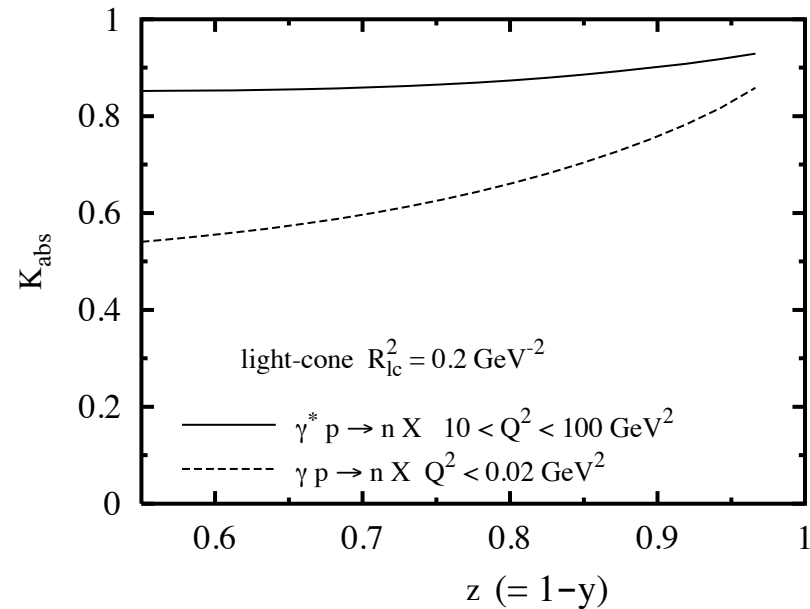
→ quality of fit depends on range of y fitted

Leading neutron production at HERA

- At large y , non-pionic mechanisms contribute (*e.g.* heavier mesons, absorption)



Kopeliovich et al. (2012)

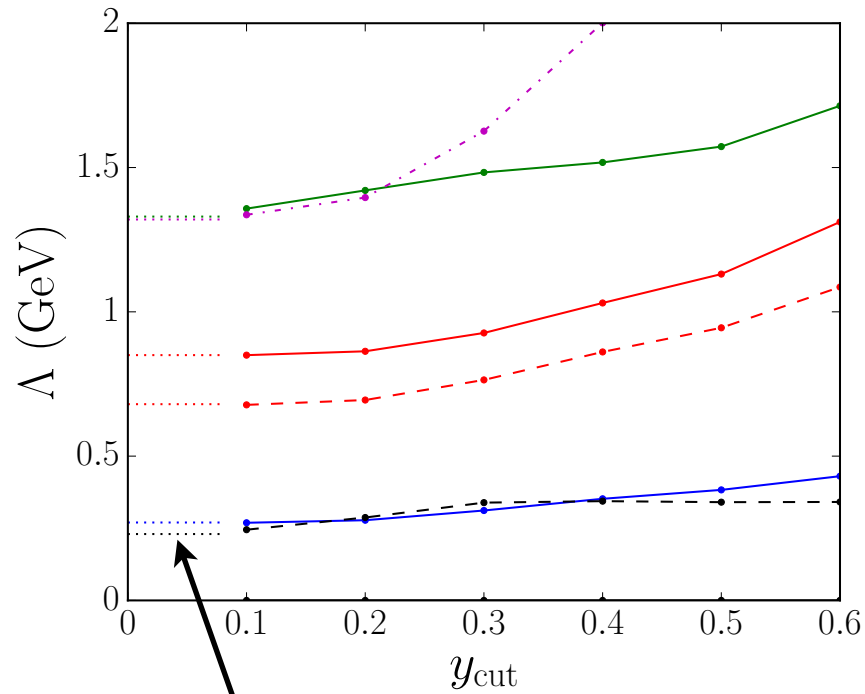


D'Alesio, Pirner (2000)

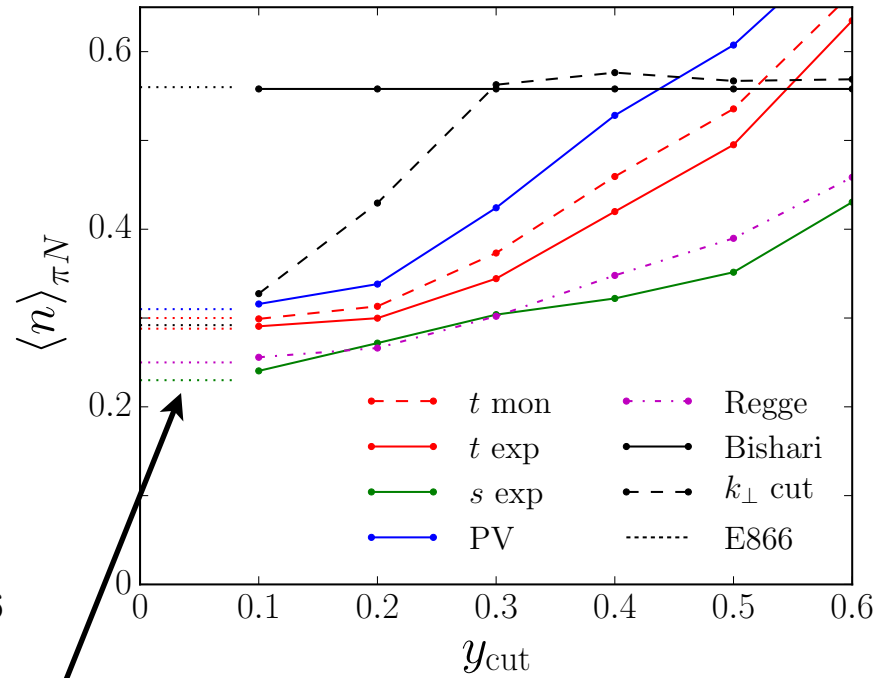
- To reduce model dependence, fit the value of y_{cut} up to which data can be described in terms of π exchange

Leading neutron production at HERA

- Fit requires higher momentum pions with increasing y_{cut}



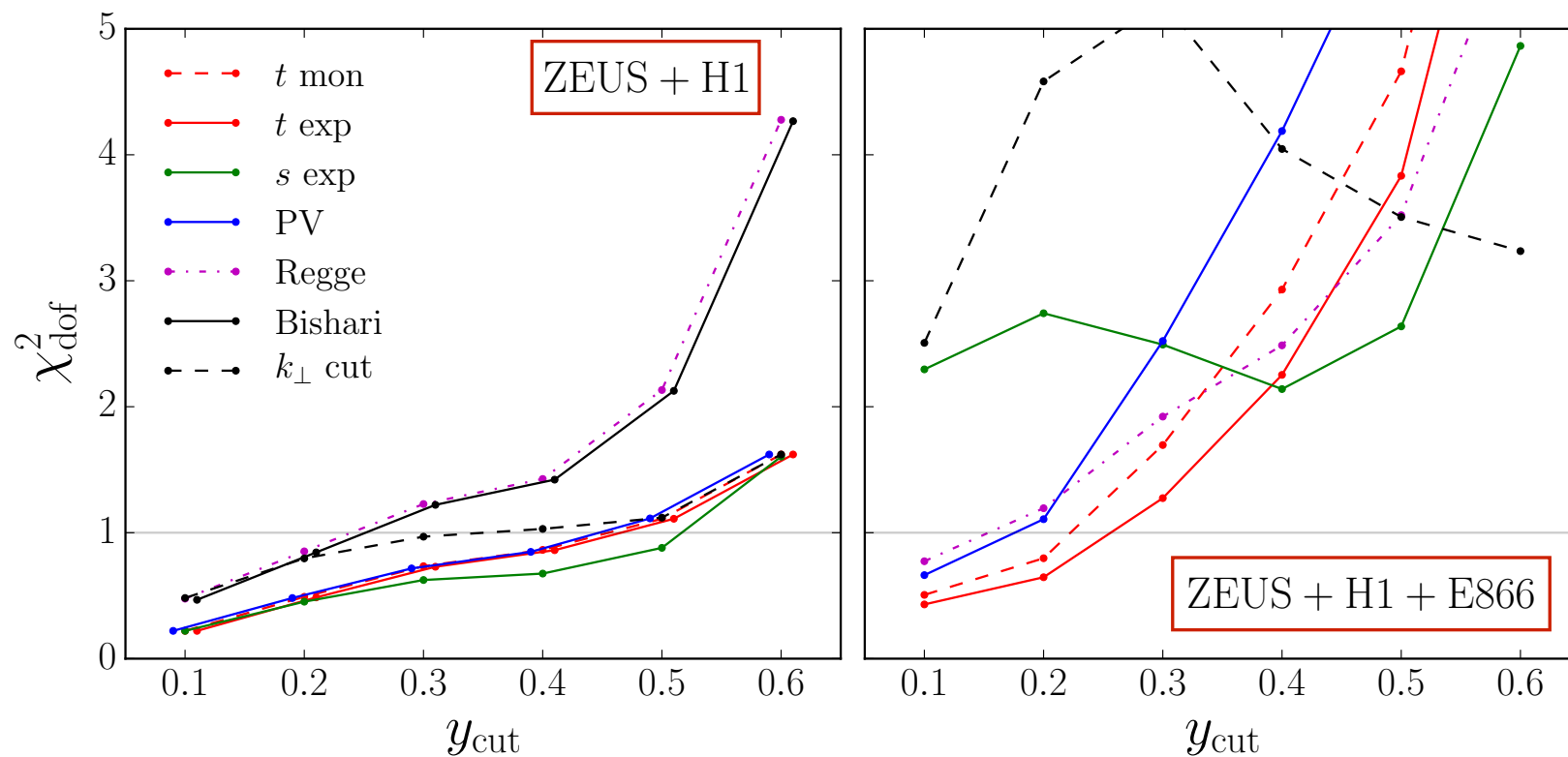
values from fit to E866 data only



→ larger values of y_{cut} more in conflict with E866 data

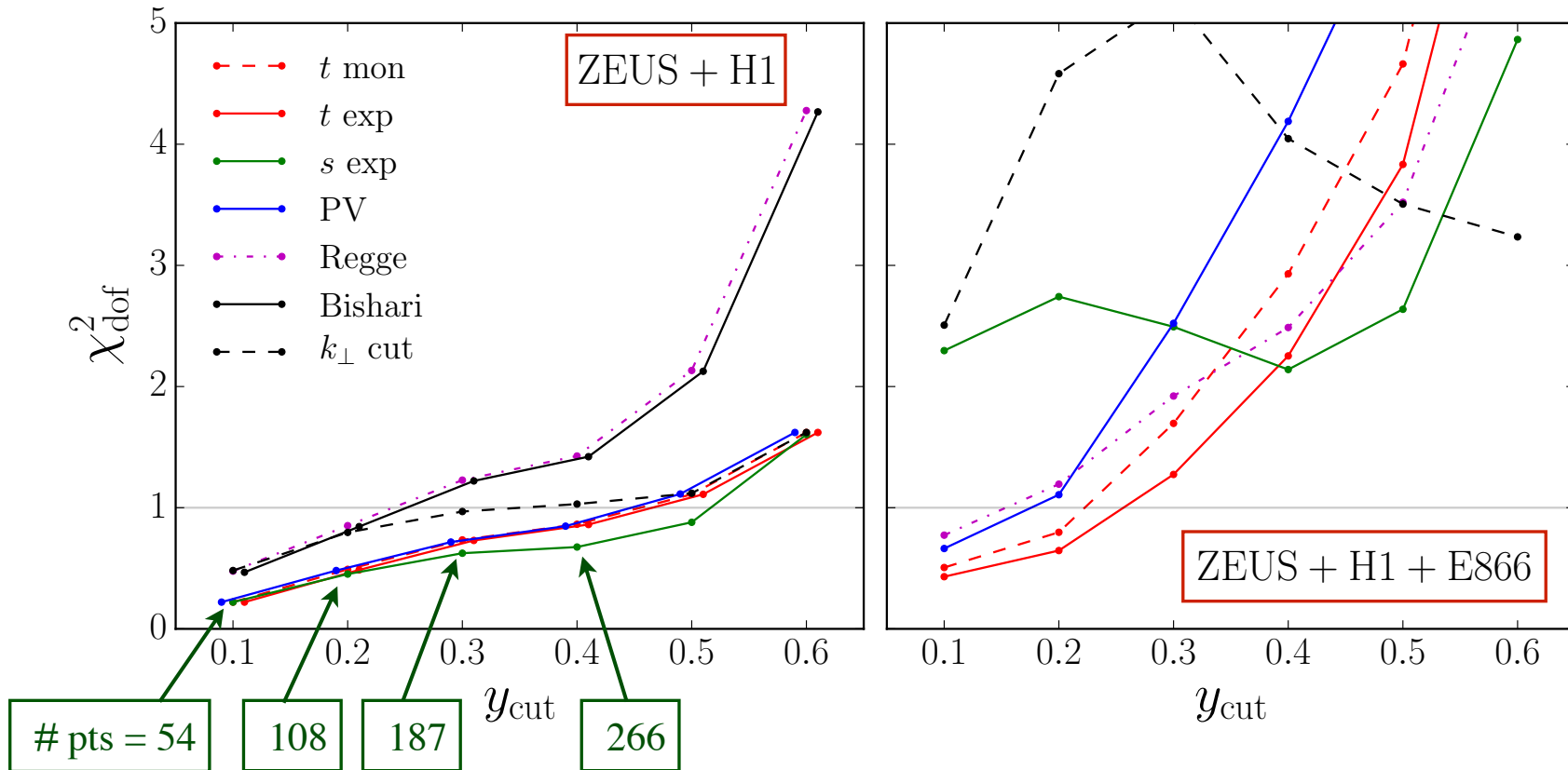
Leading neutron production at HERA

■ Combined fit to HERA LN and E866 Drell-Yan data



Leading neutron production at HERA

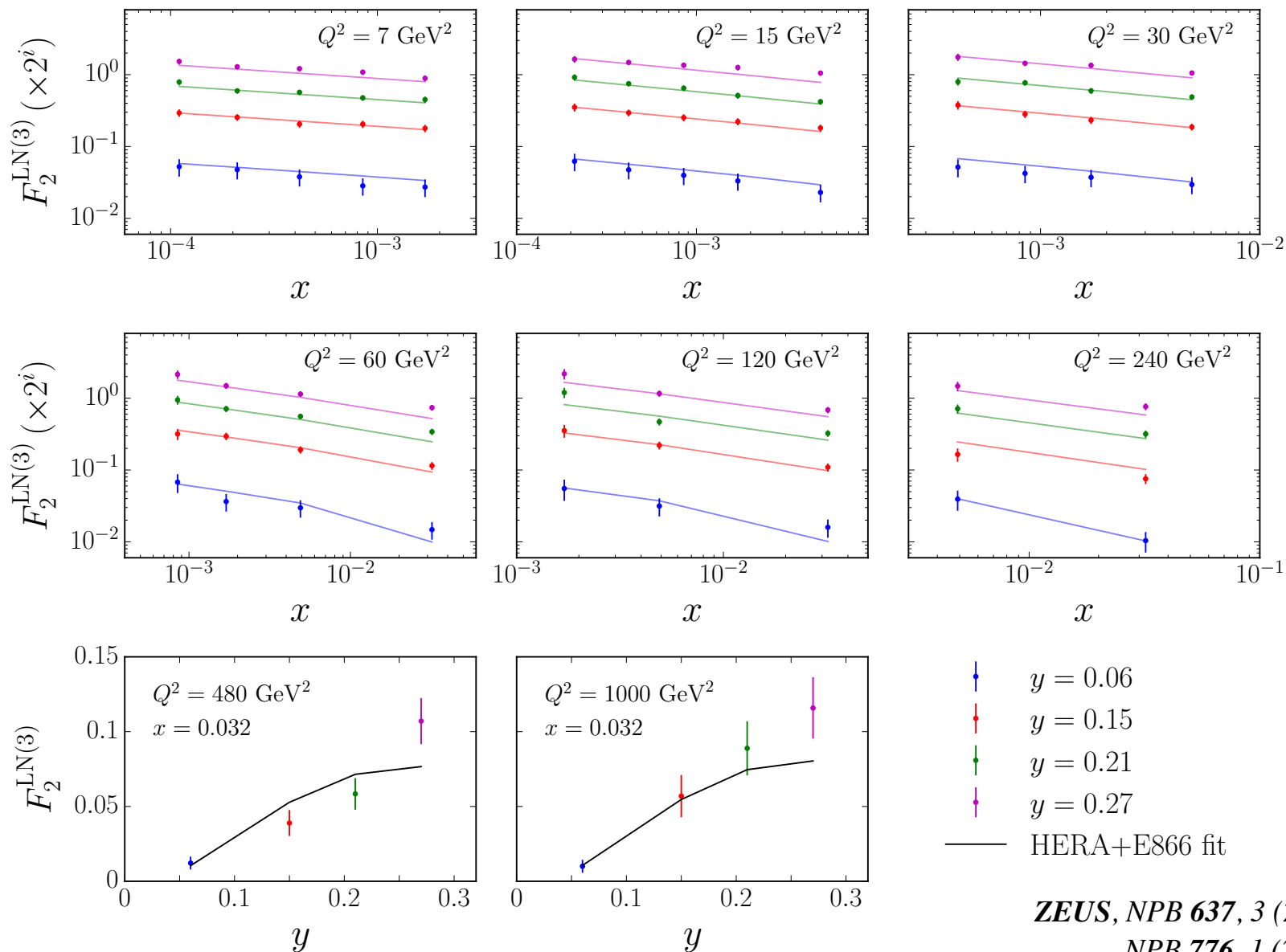
■ Combined fit to HERA LN and E866 Drell-Yan data



→ best fits for largest number of points afforded by t -dependent exponential (and t monopole) regulators

Leading neutron production at HERA

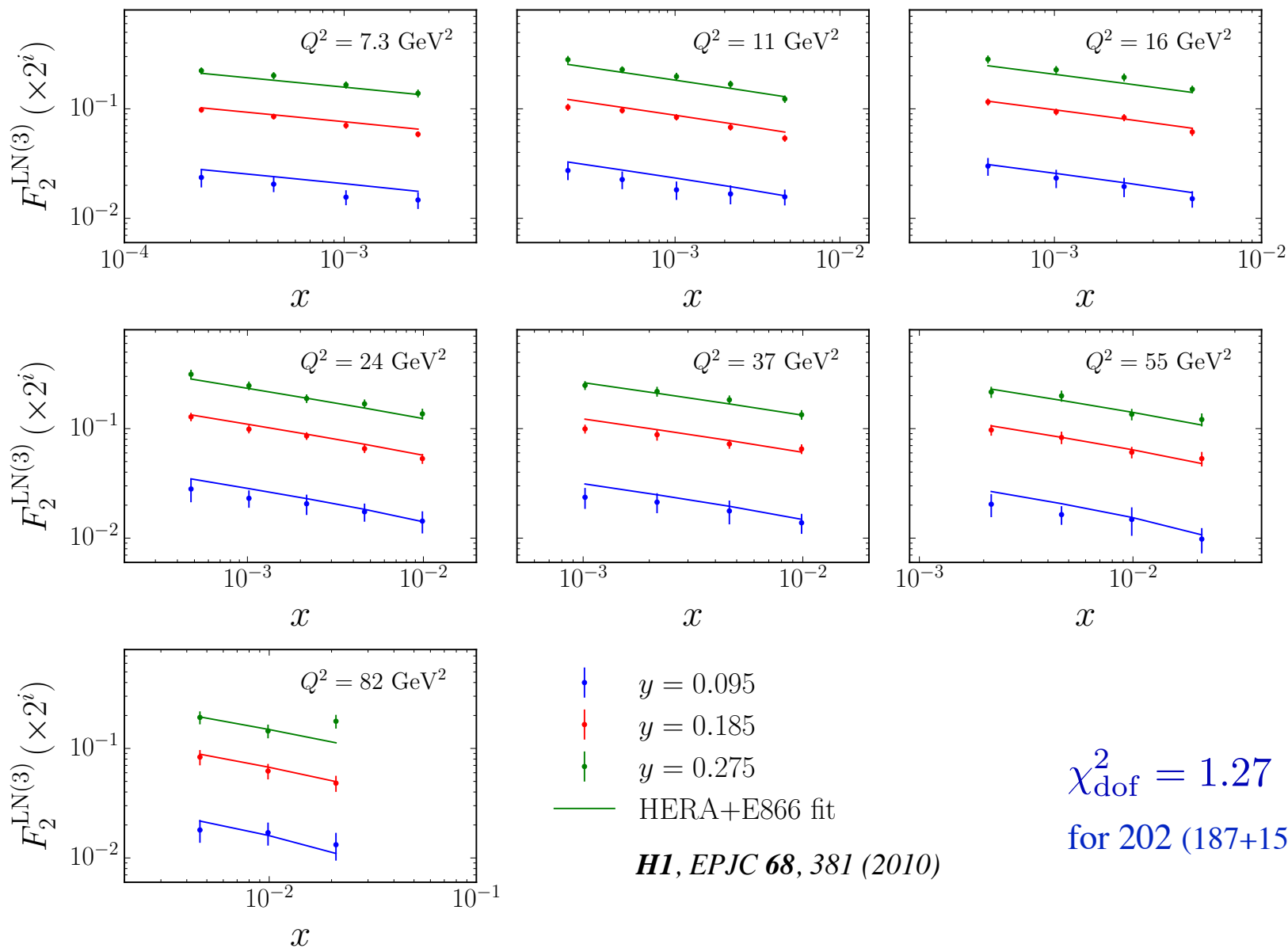
Fit to ZEUS LN spectra for $y_{\text{cut}} = 0.3$ (t -dependent exponential)



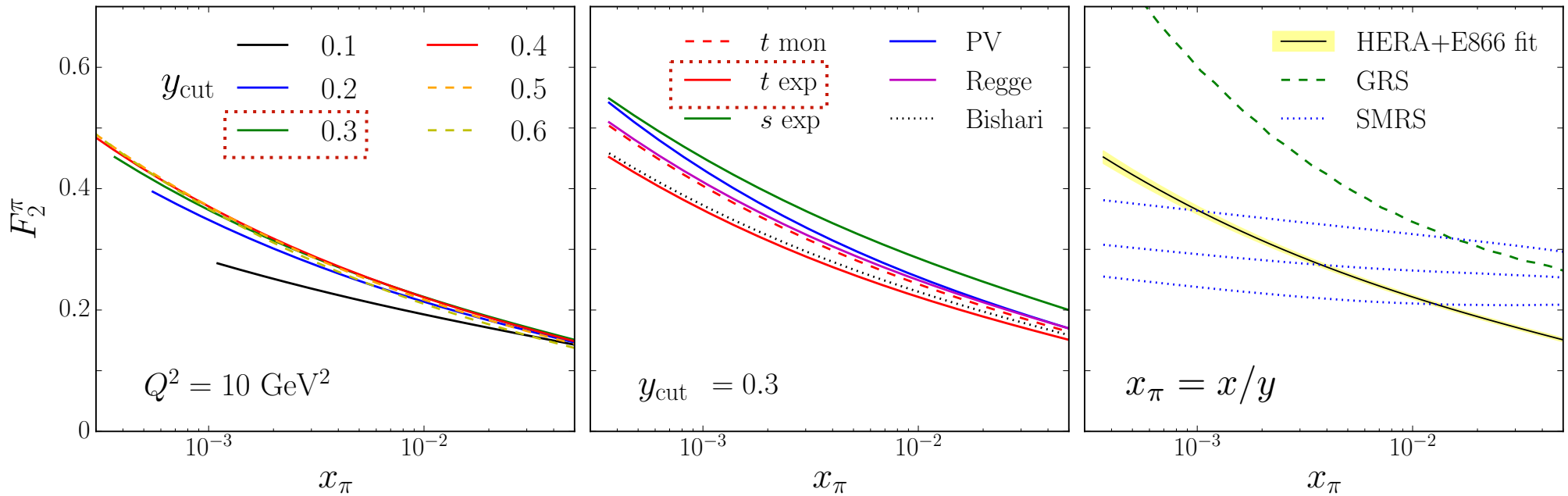
ZEUS, NPB 637, 3 (2002)
NPB 776, 1 (2007)

Leading neutron production at HERA

- Fit to H1 LN spectra for $y_{\text{cut}} = 0.3$ (t -dependent exponential)



Extracted pion structure function

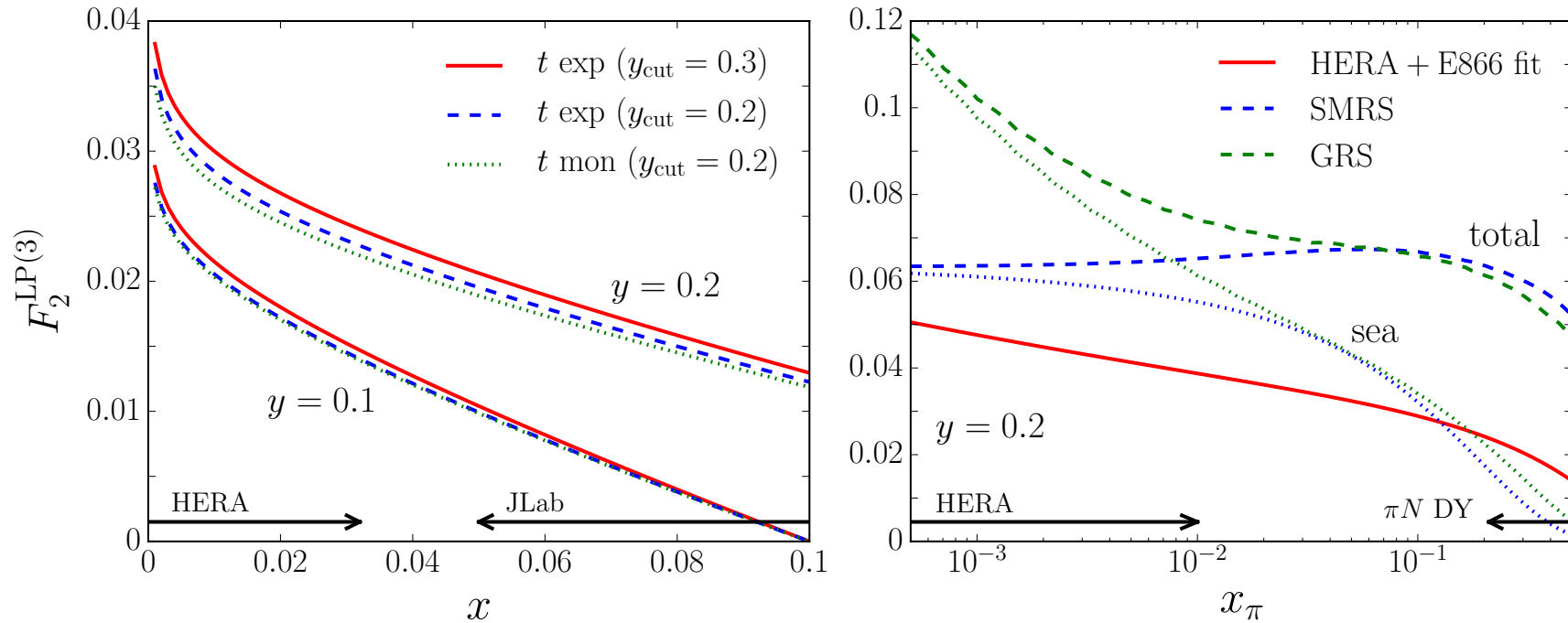


$$F_2^\pi = N x_\pi^a (1 - x_\pi)^b, \quad a = a_0 + a_1 \eta$$

$$\eta \sim \log(\log Q^2)$$

- stable values of F_2^π at $4 \times 10^{-4} \lesssim x_\pi \lesssim 0.03$ from combined fit
- shape similar to GRS fit to πN Drell-Yan data (for $x_\pi \gtrsim 0.2$), but smaller magnitude

Predictions at TDIS kinematics



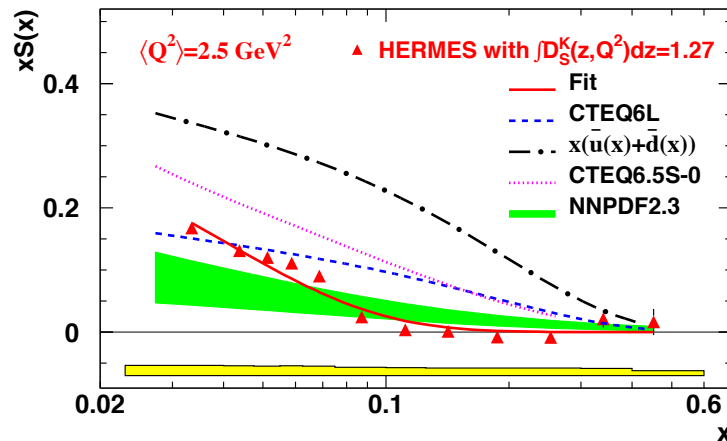
→ JLab TDIS experiment can fill gap in x_π coverage between HERA and πN Drell-Yan kinematics

Strange quarks

☀ Traditionally, strange quark PDFs most directly determined from $\mu^+\mu^-$ production in $\nu(\bar{\nu})A$ DIS ($W^+s \rightarrow c / W^-\bar{s} \rightarrow \bar{c}$)

→ but significant uncertainty from nuclear corrections, semileptonic branching ratio uncertainty

→ tension with HERMES semi-inclusive K -production data?



historically, strange to nonstrange ratio

$$\kappa = \frac{s + \bar{s}}{\bar{u} + \bar{d}} \sim 0.2 - 0.5$$

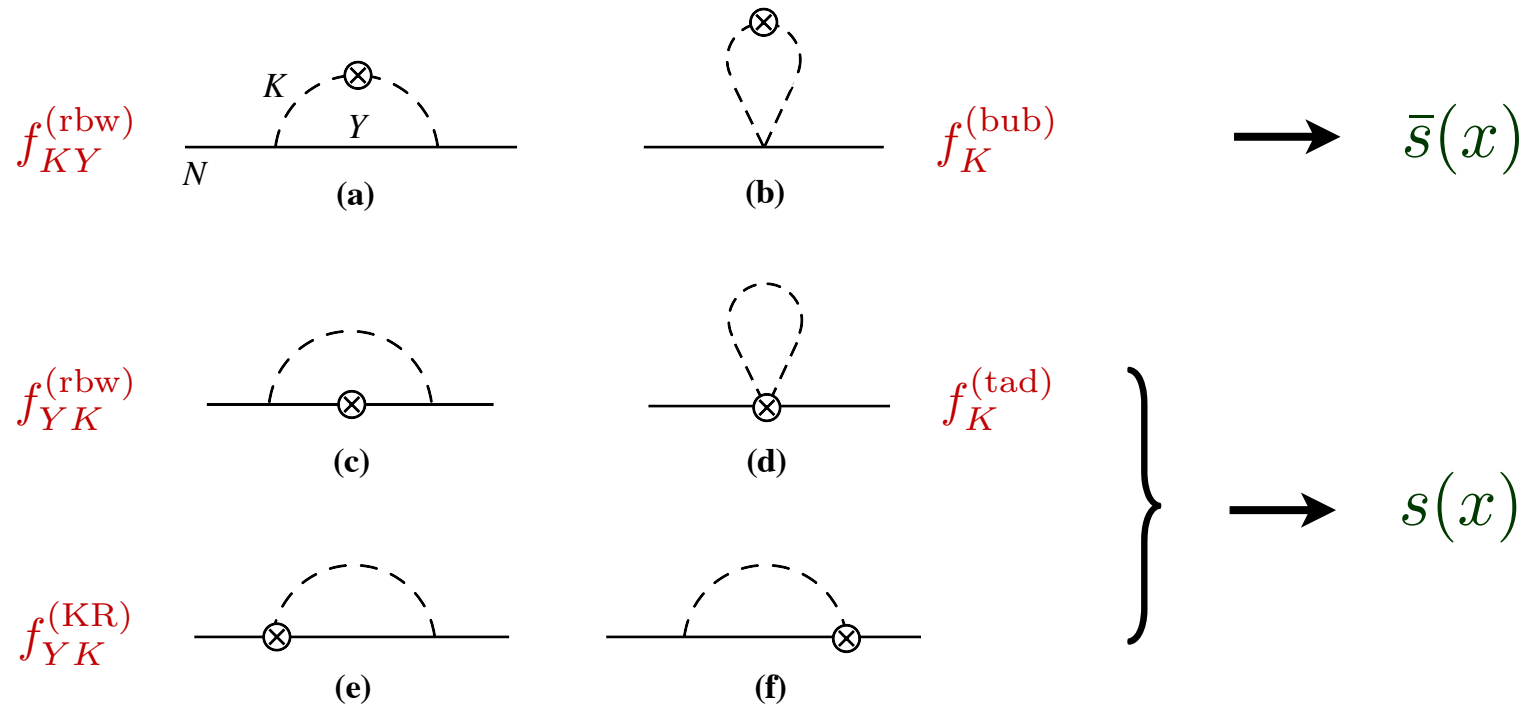
☀ Some indication of strange–antistrange asymmetry from $\nu/\bar{\nu}$ DIS data

$$S^- = \int_0^1 dx x(s - \bar{s}) = (2.0 \pm 1.4) \times 10^{-3}$$

NuTeV (2007)

Strange quarks

☀ Chiral SU(3) effective theory analysis suggests natural mechanism for generating strange asymmetry



→ gauge invariance requires the relations

$$f_{YK}^{(rbw)} + f_{YK}^{(KR)} = f_{KY}^{(rbw)}$$

$$f_K^{(tad)} + f_K^{(bub)} = 0$$

Strange quarks

☀ Convolution representation

$$\bar{s} = \left(f_{KY}^{(\text{rbw})} + f_K^{(\text{bub})} \right) \otimes \bar{s}_K$$

$$s = \left(\bar{f}_{YK}^{(\text{rbw})} \otimes s_Y + \bar{f}^{(\text{KR})} \otimes s_Y^{(\text{KR})} \right) + \bar{f}_K^{(\text{tad})} \otimes s_K^{(\text{tad})}$$

$$\bar{f}(y) \equiv f(1-y)$$

$$\sim \Delta u, \Delta d$$

$$\sim u, d$$

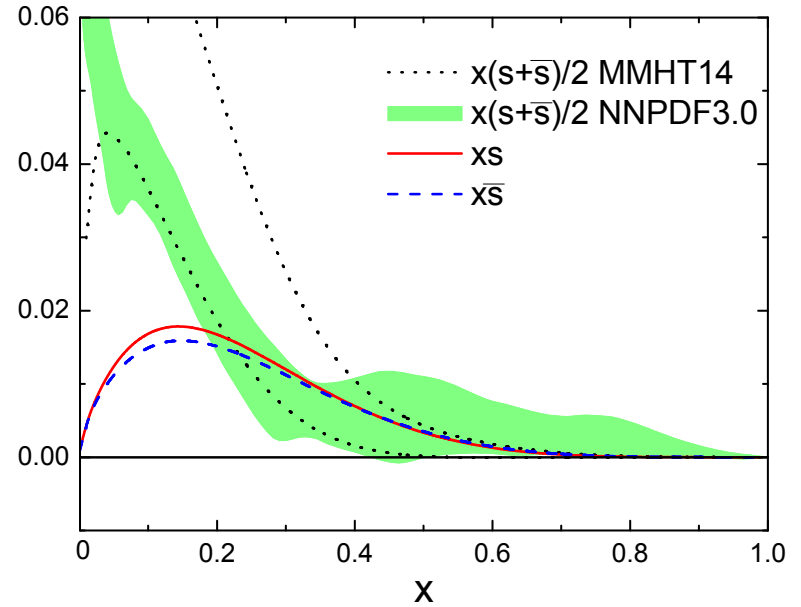
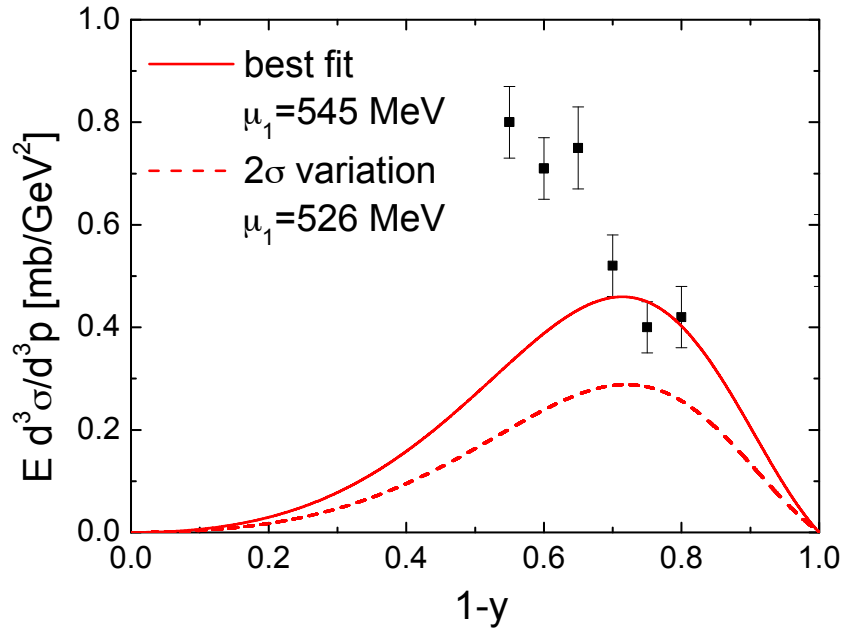
- KY splitting functions regularized using Pauli-Villars regularization
- δ -function term requires 2 subtractions (parameters μ_1, μ_2)
- since $f_K^{(\text{tad})}(y) \sim \delta(y)$, tadpole term generates *valence-like* strange-quark PDF

$$\sim s_K^{(\text{tad})}(x)$$

Strange quarks



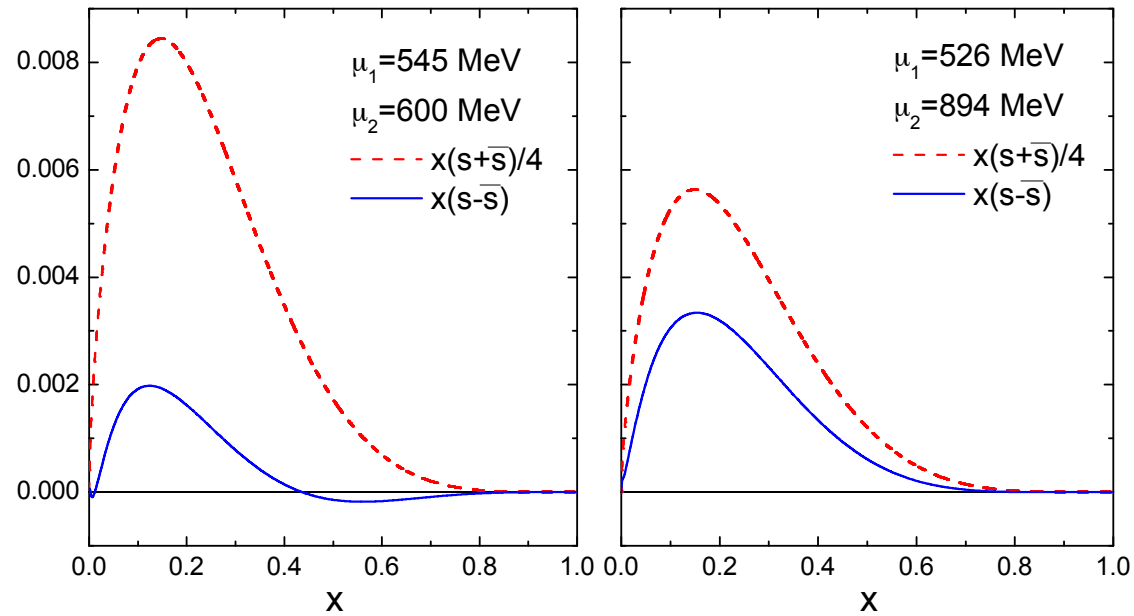
Constraints on cutoff parameters from $pp \rightarrow \Lambda X$
and total $(s + \bar{s})_{\text{loops}} \leq (s + \bar{s})_{\text{total}}$



Strange quarks



Gives rise to small but (mostly) positive $s - \bar{s}$ distribution



→ x -weighted difference $S^- = (0.4 - 1.1) \times 10^{-3}$

Outlook

- Eagerly await final SeaQuest data!
 - settle question of sign change in $\bar{d} - \bar{u}$ at high x
- Combine “leading neutron” analysis with πN Drell-Yan data to constrain pion PDFs at low and high x
 - talk by Nobuo Sato
- Compare chiral SU(3) predictions for $s - \bar{s}$ asymmetry with future lattice and experimental (SIDIS?) data