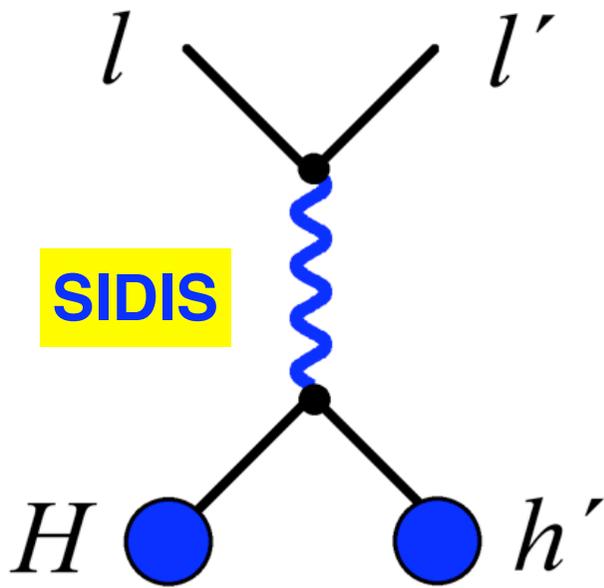


Drell-Yan:  
The Missing Spin Programme



## Leptons: clean, surgical tools

SIDIS

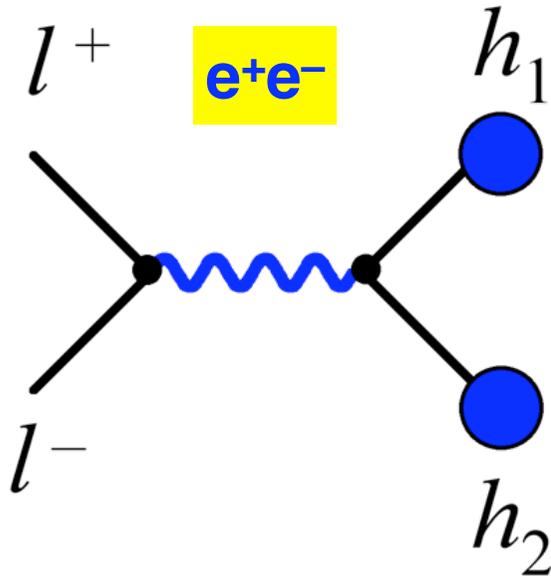


$$\sum_q e_q^2 f_q^{(H)}(x) D_q^{h'}(z)$$

- Disentangle **distribution** ( $f$ ) and **fragmentation** ( $D$ ) functions  $\rightarrow$  measure **all process**
- Disentangle **quark flavours**  $q \rightarrow$  measure as many **hadron species**  $H, h$  as possible

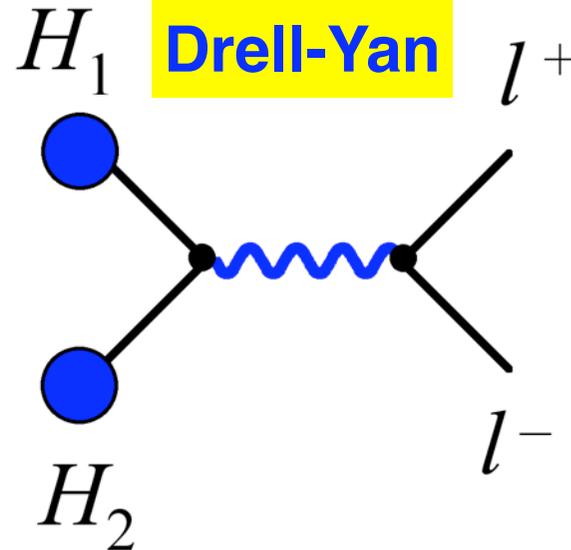
These are the **only** processes where TMD factorization is proven

$e^+e^-$



$$\sum_q e_q^2 D_q^{h_1}(z_1) D_q^{h_2}(z_2)$$

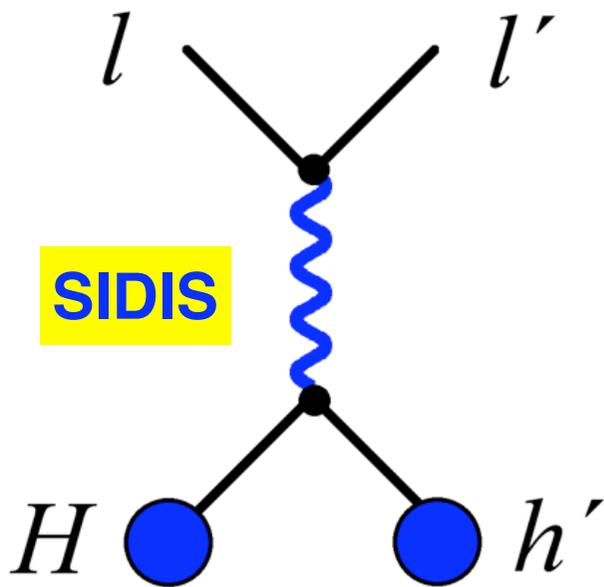
Drell-Yan



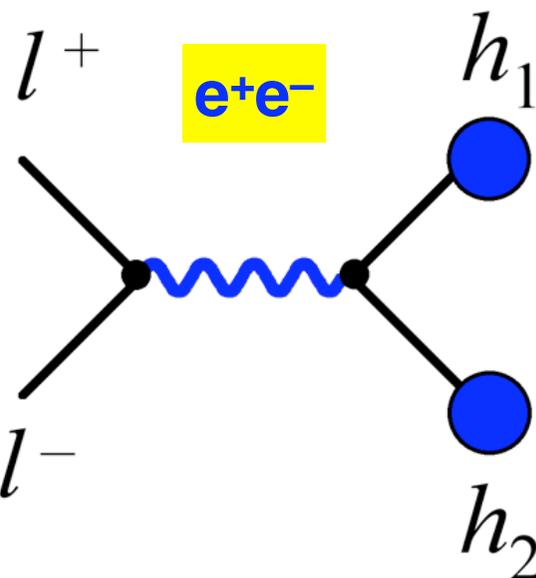
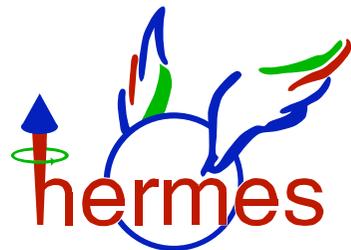
$$\sum_q e_q^2 f_q^{(H_1)}(x_1) f_{\bar{q}}^{(H_2)}(x_2)$$

# Leptons: clean, surgical tools

SIDIS

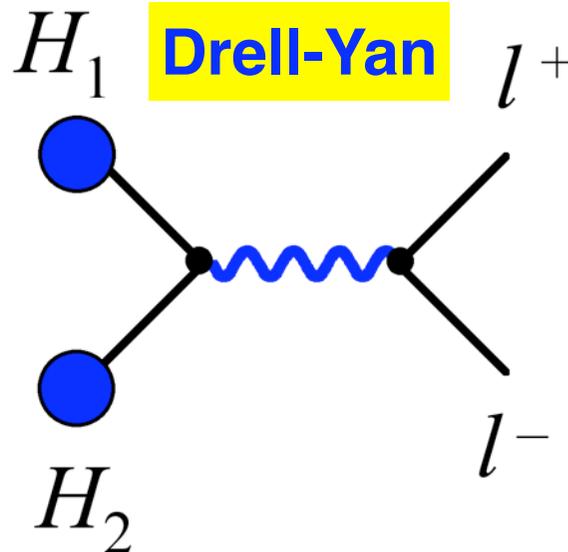


$$\sum_q e_q^2 f_q^{(H)}(x) D_q^{h'}(z)$$



e<sup>+</sup>e<sup>-</sup>

$$\sum_q e_q^2 D_q^{h_1}(z_1) D_{\bar{q}}^{h_2}(z_2)$$



Drell-Yan

$$\sum_q e_q^2 f_q^{(H_1)}(x_1) f_{\bar{q}}^{(H_2)}(x_2)$$



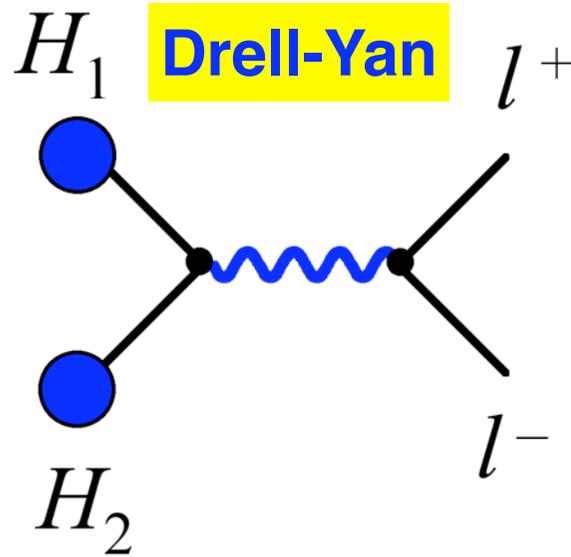
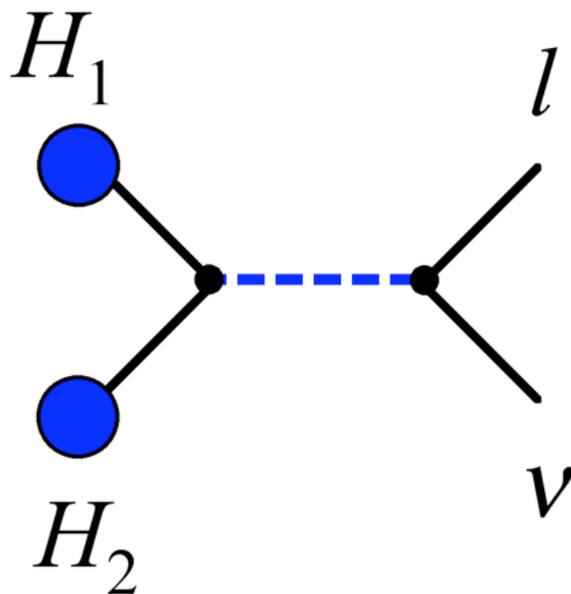
# The Missing Spin Program: Drell-Yan



$$\sum_q e_q^2 \mathbf{f}_q^{(H_1)}(x_1) \mathbf{f}_{\bar{q}}^{(H_2)}(x_2)$$



## W production



- Clean access to **sea quarks**  
e.g.  $\Delta\bar{u}(x), \Delta\bar{d}(x)$  at RHIC
- Crucial test of **TMD formalism**  
→ **sign change** of T-odd functions
- A **complete** spin program requires multiple hadron species  
→ **nucleon & meson beams**

# SeaQuest

2032





## Sea-Spin: The Final Frontier



- Baseline: **Sign Change** for T-Odd TMD( $u$ )
  - as much **u-quark dominance** as possible
  - **COMPASS-II  $\pi p \uparrow$**  first: **approved** and best kinematic match to SIDIS
  - SeaQuest with beam  $\uparrow$  maybe next
- The Sea: **Sivers & Boer-Mulders** for **ubar**
  - **$\pi^+$  or p beam** at COMPASS-II
  - possible fixed target:  **$N \uparrow$  target**  
@ SeaQuest, PANDA
  - possible **collider**: RHIC, NICA
- **Full Programme** of Spin-DY: TMD ultimate global fit
  - **many beam species** (  $p$ ,  $pbar$ ,  $d$ , meson ): flavor separation
  - **L and T** polarization: TMD separation



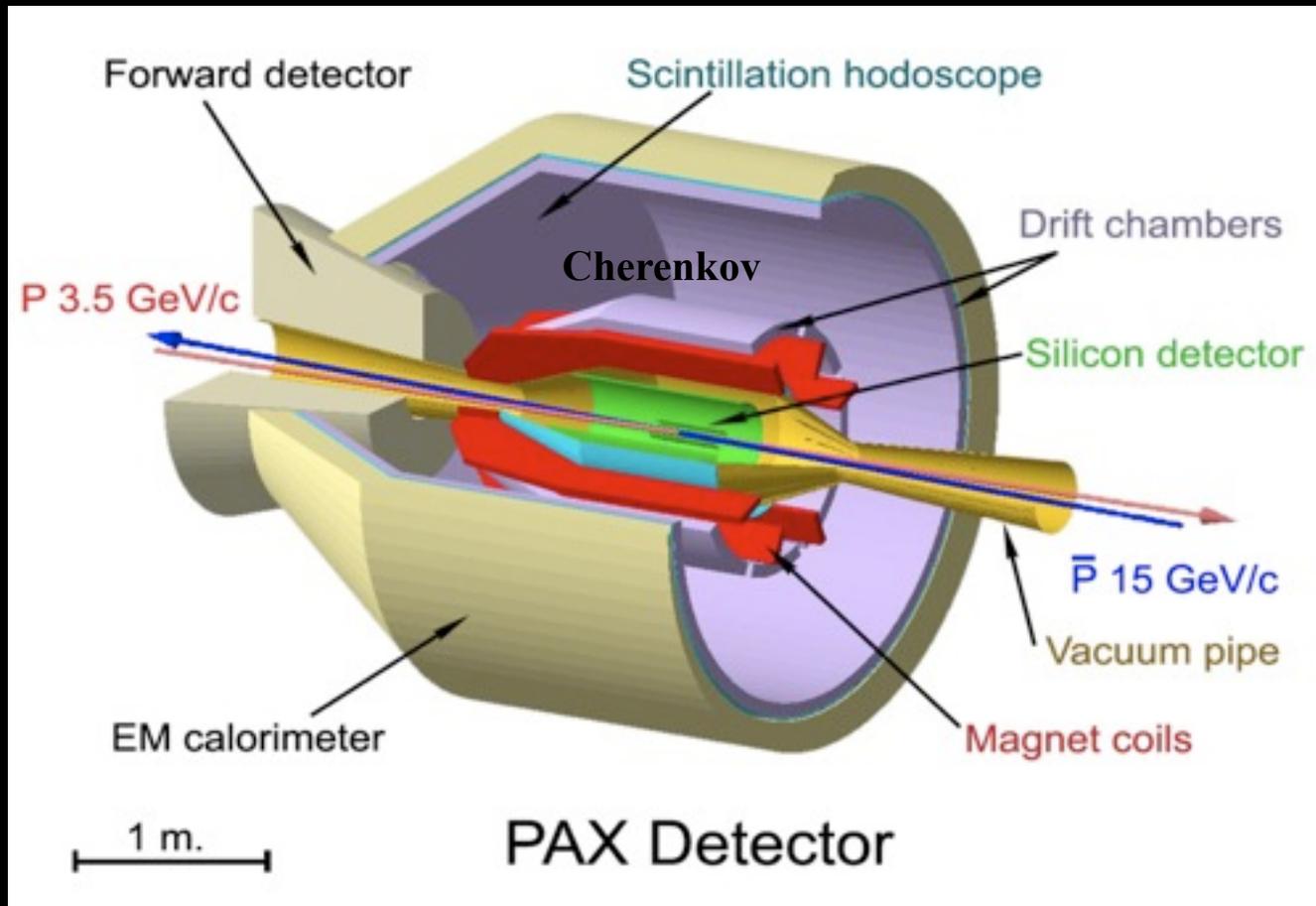
## Sea-Spin: The Final Frontier



- Baseline: **Sign Change** for T-Odd TMD(u)
  - as much **u-quark dominance** as possible
  - **COMPASS-II  $\pi p \uparrow$**  first: approved and best kinematic match to SIDIS
  - SeaQuest with beam  $\uparrow$  maybe next
- The Sea: **Sivers & Boer-Mulders** for **ubar**
  - **$\pi^+$  or p beam** at COMPASS-II
  - possible fixed target:  **$N \uparrow$  target** @ SeaQuest, PANDA
  - possible **collider**: RHIC, NICA
- **Full Programme** of Spin-DY: TMD ultimate global fit
  - **many beam species** ( p, pbar, d, meson ): flavor separation
  - **L and T** polarization: TMD separation



# The Experiments



# The Experiments

 approved  
 unpolarized

experiment		particles	energy	$x_b$ or $x_t$	luminosity	timeline
COMPASS (CERN)	[22]	$\pi^\pm + p^\uparrow$	160 GeV $\sqrt{s} = 17.4$ GeV	$x_t = 0.2 - 0.3$	$1 \times 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$	2014
PAX (GSI)	[23]	$p^\uparrow + \bar{p}$	collider $\sqrt{s} = 14$ GeV	$x_b = 0.1 - 0.9$	$2 \times 10^{30} \text{ cm}^{-2} \text{ s}^{-1}$	>2017
PANDA (GSI)	[24]	$\bar{p} + p^\uparrow$	15 GeV $\sqrt{s} = 5.5$ GeV	$x_t = 0.2 - 0.4$	$2 \times 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$	>2016
NICA (JINR)	[25]	$p^\uparrow + p$	collider $\sqrt{s} = 20$ GeV	$x_b = 0.1 - 0.8$	$1 \times 10^{30} \text{ cm}^{-2} \text{ s}^{-1}$	>2014
PHENIX (BNL)	[26]	$p^\uparrow + p$	collider $\sqrt{s} = 200$ GeV	$x_b = 0.05 - 0.1$	$2 \times 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$	>2018
RHIC internal target phase-1	[27]	$p^\uparrow + p$	250 GeV $\sqrt{s} = 22$ GeV	$x_b = 0.25 - 0.4$	$2 \times 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$	>2015
RHIC internal target phase-2	[27]	$p^\uparrow + p$	250 GeV $\sqrt{s} = 22$ GeV	$x_b = 0.25 - 0.4$	$3 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$	>2018
E-906/SeaQuest (FNAL)	[12]	$p + p$	120 GeV $\sqrt{s} = 15$ GeV	$x_b = 0.3 - 0.9$ $x_t = 0.1 - 0.45$	$2.0 \times 10^{35} \text{ cm}^{-2} \text{ s}^{-1}$	2011
pol. SeaQuest (FNAL)		$p^\uparrow + p$	120 GeV $\sqrt{s} = 15$ GeV	$x_b = 0.3 - 0.9$	$1 \times 10^{36} \text{ cm}^{-2} \text{ s}^{-1}$	>2014

“valence”  $x$   
 $\approx 0.25$   
 → match to SIDIS

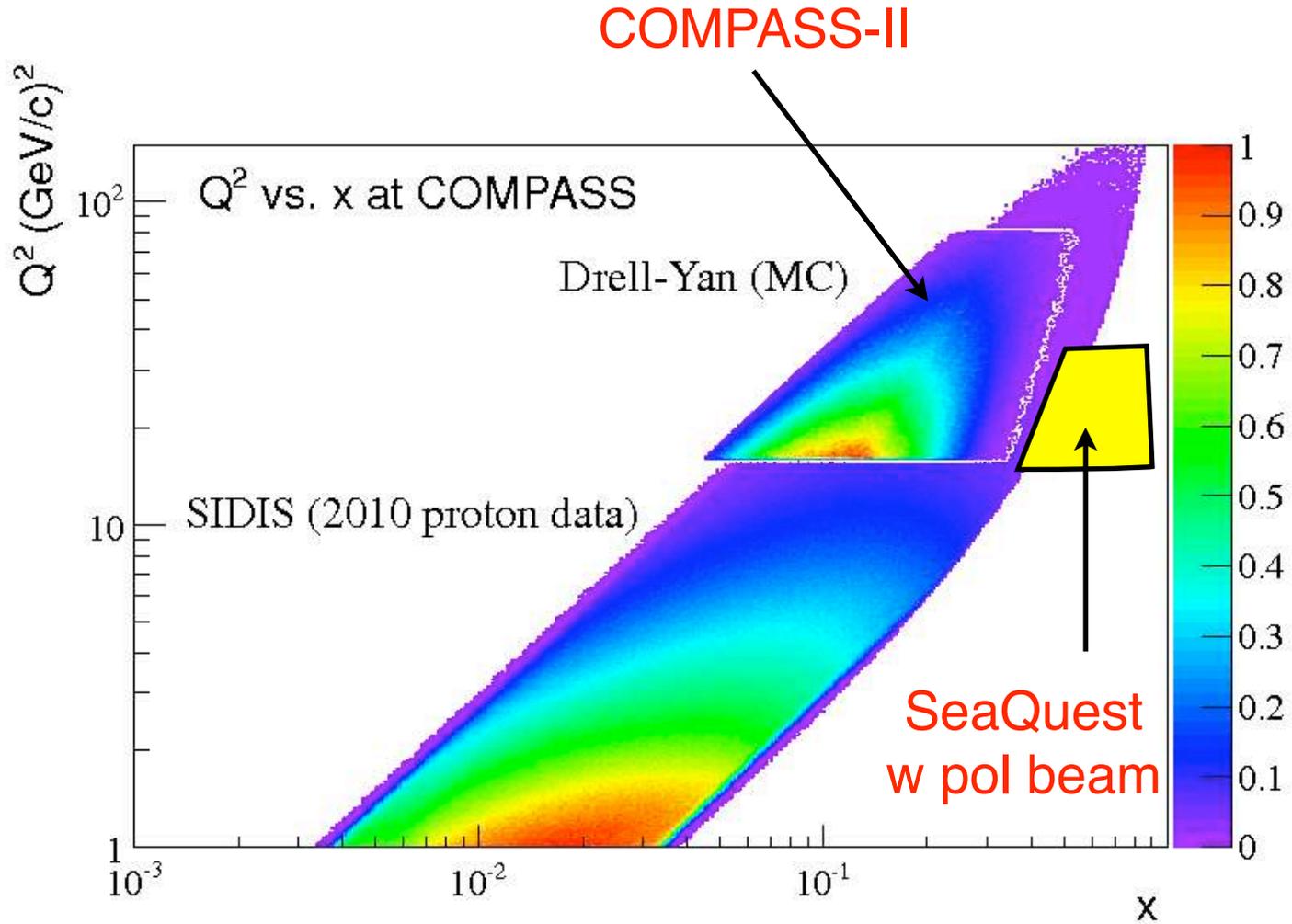
high  $\sqrt{s}$   
 → low  $x$ ,  
 low lumi

low mass  
 $M_{\mu\mu} < 2.5$



“Valence x” → where the action is

# SIDIS and DY



The COMPASS SIDIS and DY experimental measurements have an overlapping region.

# AnDY - Staging



Polarized proton runs at  $\sqrt{s} = 500$  GeV

- 2011 HCal + newly constructed BBC at IP2 to establish the impact of a 3<sup>rd</sup> IR operation and to demonstrate the calibration of HCal to get first data constraints on charged hadron backgrounds
- 2012 HCal + EmCal + neutral/charge veto + BBC for zero-field data sample with  $L_{\text{int}} \approx 150 \text{ pb}^{-1}$  and  $P \approx 50\%$  to observe di-leptons from  $J/\psi$ ,  $\Upsilon$ , and intervening continuum.
- 2013 HCal + EmCal + neutral/charge veto + BBC + split-dipole for data sample with  $L_{\text{int}} \approx 150 \text{ pb}^{-1}$  and  $P = 50\%$  to observe di-leptons from  $J/\psi$ ,  $\Upsilon$ , and intervening continuum to address whether charge sign discrimination is required

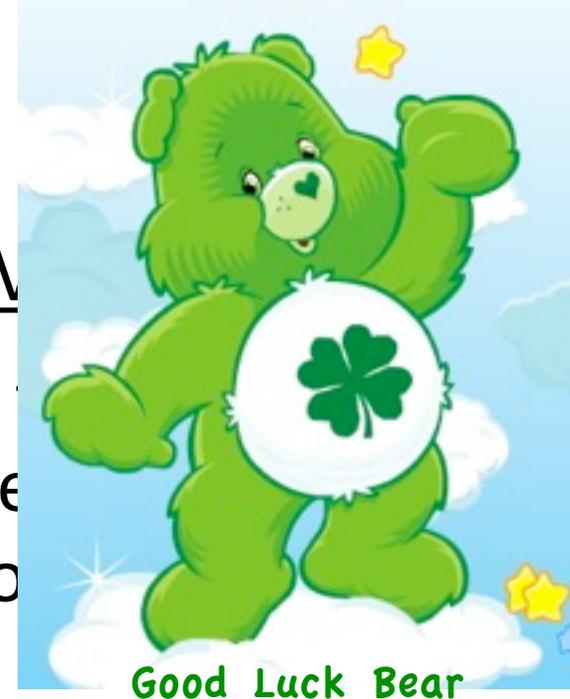
# AnDY - Staging



Polarized proton runs at  $\sqrt{s} = 500$  GeV

2011

HCal + newly constructed BBC at IP2  
 impact of a 3<sup>rd</sup> IR operation and to de  
 calibration of HCal to get first data co  
 charged hadron backgrounds



2012

HCal + EmCal + neutral/charge veto + BBC for zero-  
 field data sample with  $L_{\text{int}} \approx 150 \text{ pb}^{-1}$  and  $P \approx 50\%$  to  
 observe di-leptons from  $J/\psi$ ,  $\Upsilon$ , and intervening  
 continuum.

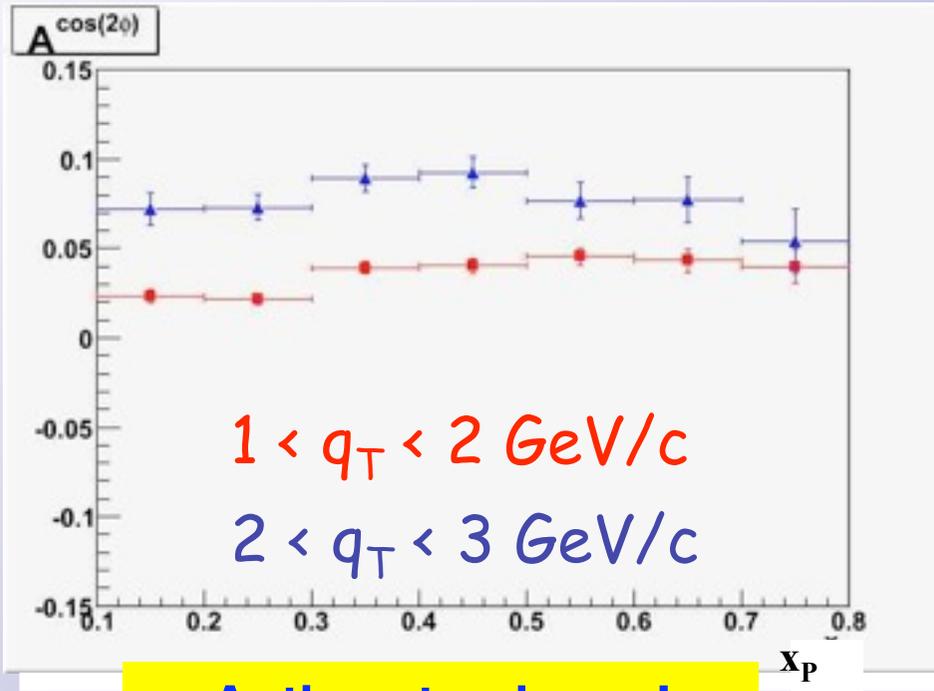
2013

HCal + EmCal + neutral/charge veto + BBC + split-  
 dipole for data sample with  $L_{\text{int}} \approx 150 \text{ pb}^{-1}$  and  $P = 50\%$   
 to observe di-leptons from  $J/\psi$ ,  $\Upsilon$ , and intervening  
 continuum to address whether charge sign  
 discrimination is required



DY @ 15 GeV/c —  $\bar{p}p \rightarrow \mu^+\mu^-X$   $\sqrt{s} = 5.5$  GeV

### UNPOLARISED



$1 < q_T < 2$  GeV/c  
 $2 < q_T < 3$  GeV/c

Anti-proton beam!  
 PAX → polarize it ...

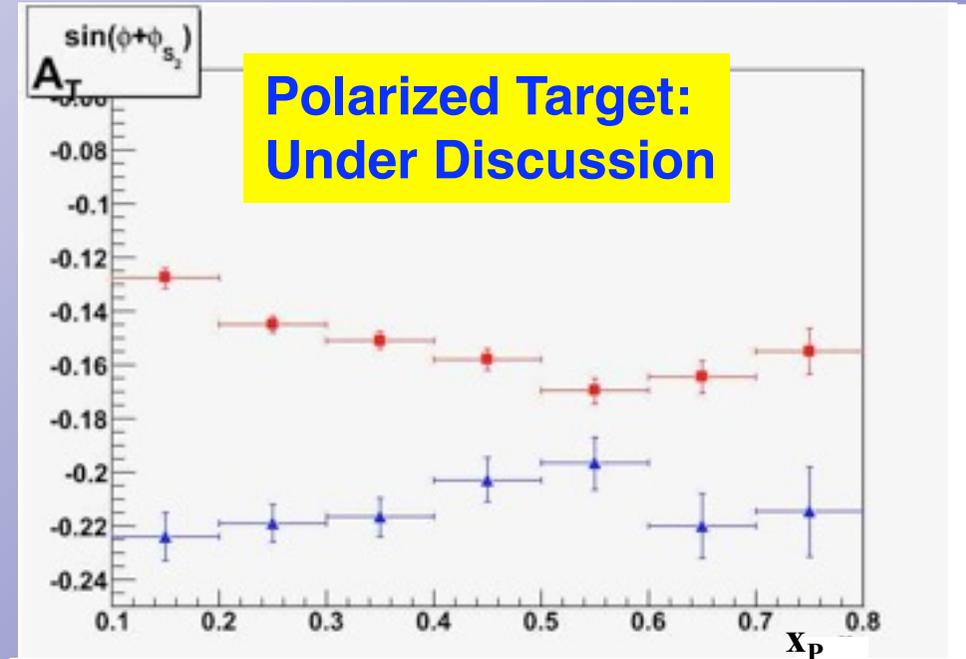
500KEv included in  
 asymmetries

Acceptance corrections

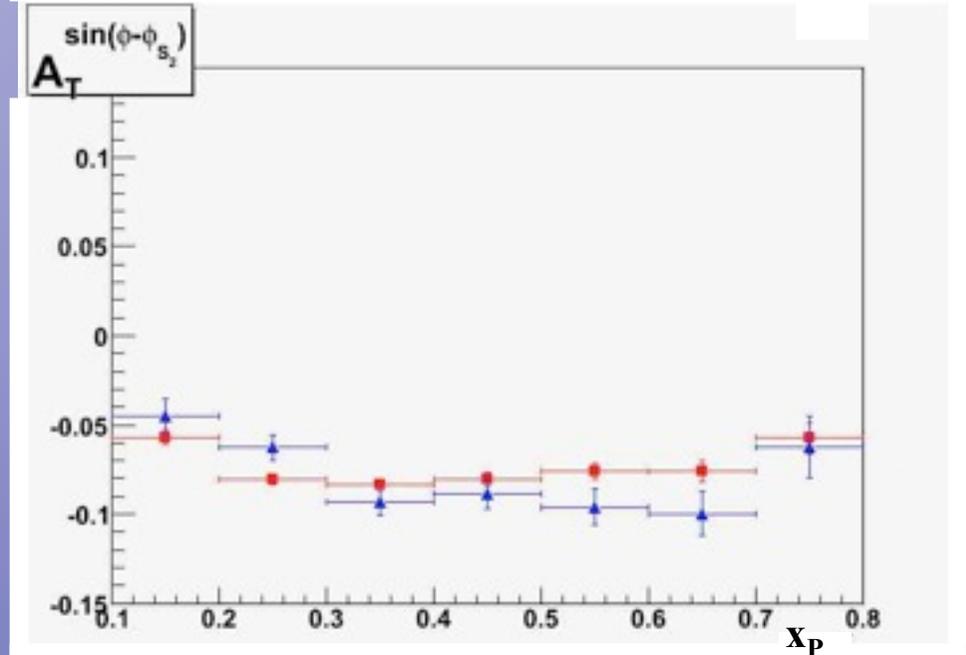
Low-Mass  $M_{\mu\mu} < 2.5$  GeV

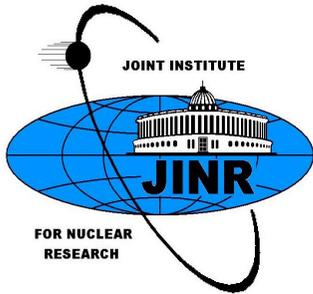
crucial!

### SINGLE-POLARISED



Polarized Target:  
 Under Discussion





# SPD EXPERIMENT AT NICA. PROPOSED MEASUREMENTS.



Extraction of unknown (poor known) parton distribution functions (PDFs):

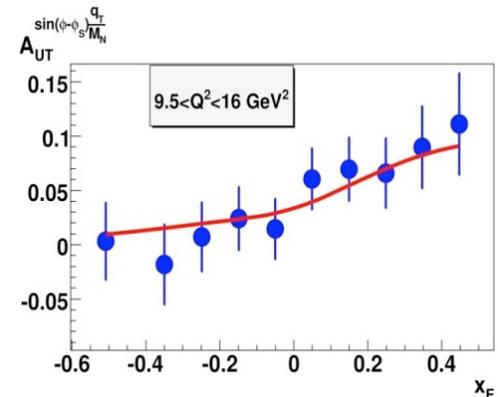
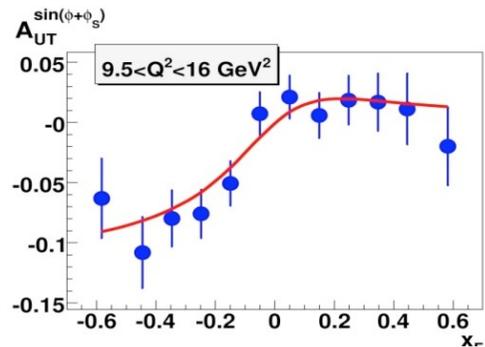
$$p(D)p(D) \rightarrow \gamma^* X \rightarrow l^+l^- X \quad \text{Boer-Mulders PDF}$$

$$p^\uparrow(D^\uparrow)p(D) \rightarrow \gamma^* X \rightarrow l^+l^- X \quad \text{Sivers PDFs (Efremov,... PLB 612 (2005), PRD 73(2006));}$$

$$p^\uparrow(D^\uparrow)p^\uparrow(D^\uparrow) \rightarrow \gamma^* X \rightarrow l^+l^- X \quad \text{Transversity PDF (Anselmino, Efremov, ...)}$$

$$p^\uparrow(D^\uparrow)p(D) \rightarrow \gamma^* X \rightarrow l^+l^- X \quad \text{Transversity and first moment of Boer-Mulders PDFs (Sissakian, Shevchenko, Nagaytsev, Ivanov, PRD 72(2005), EPJ C46, 2006 C59, 2009)}$$

$$p^\rightarrow(D^\rightarrow)p^\leftarrow(D^\leftarrow) \rightarrow \gamma^* X \rightarrow l^+l^- X \quad \text{Longitudinally polarized sea and strange PDFs and tensor deuteron structure (Teryaev, ...)}$$



The same PDFs from  $J/\psi$  production processes ( $\sqrt{s} \leq 10 \text{ GeV}$ ).

"Polarization effects in Drell-Yan processes",

**Sissakian A. N.**, Shevchenko O. Yu., Nagaitsev A. P., Ivanov O. N.  
Physics of Particles and Nuclei, Volume 41, Issue 1, pp.64-100,  
MAIK award for 2010.

**Low-Mass**  
 **$M_{\mu\mu} < 2.5 \text{ GeV}$**

**Polarized everything! Including  $d!$**

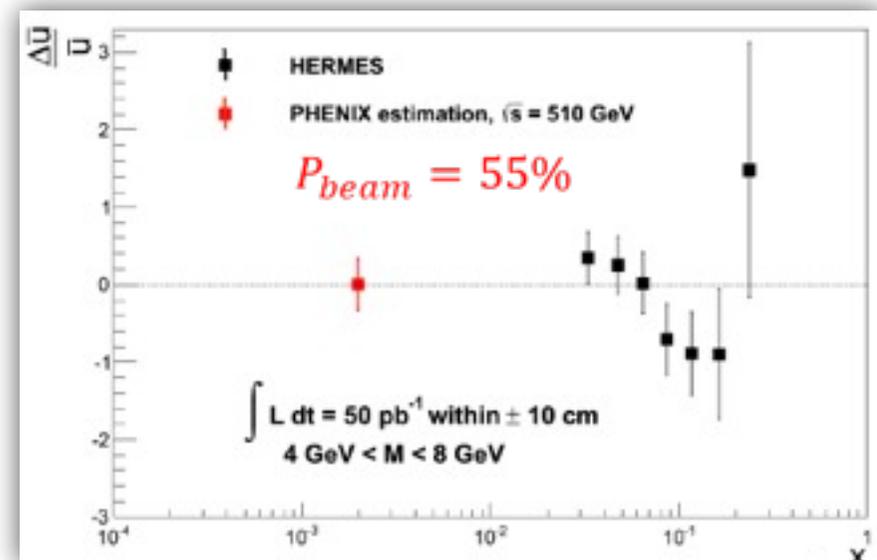
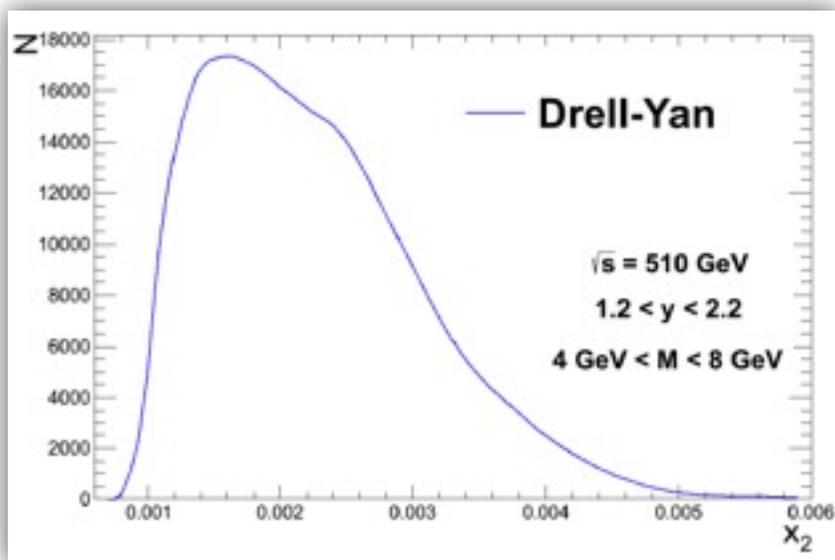
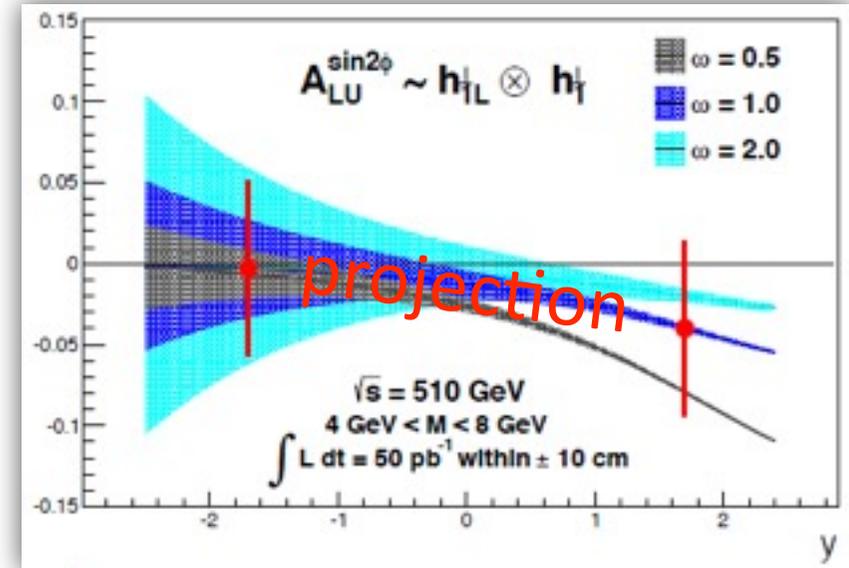
# Highest $\sqrt{s}$ polarized $\rightarrow$ RHIC collider

## Polarized Drell-Yan with PHENIX?

arXiv:1108.4974 (Lu, Ma, Zhu)

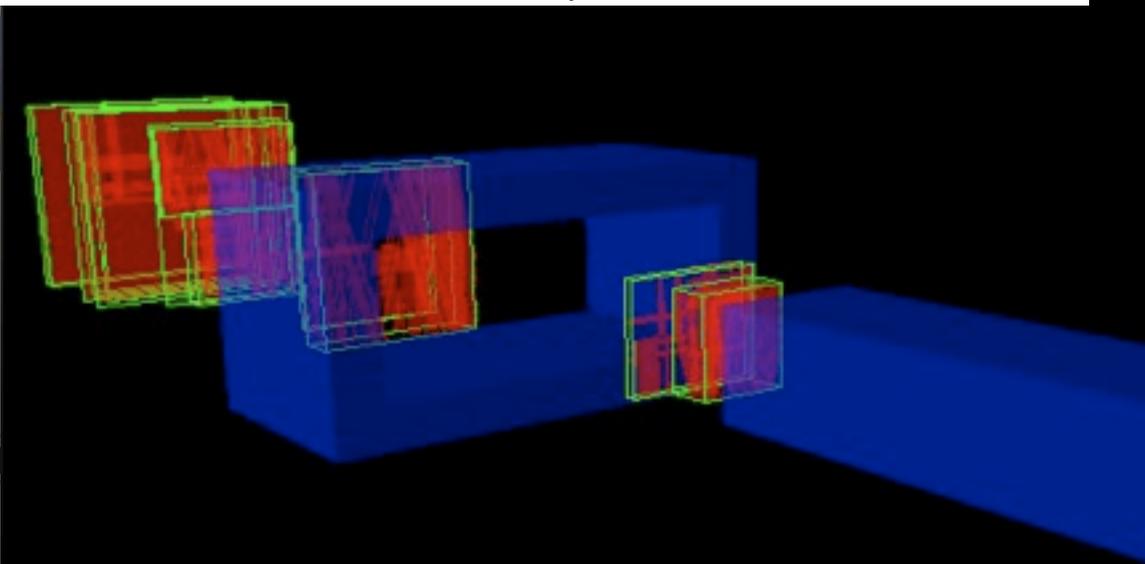
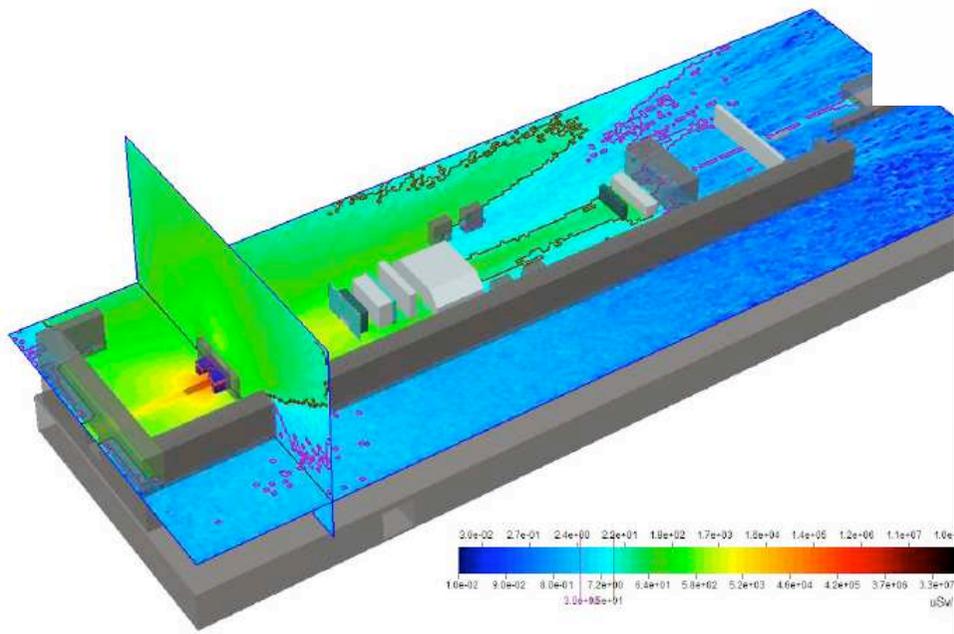
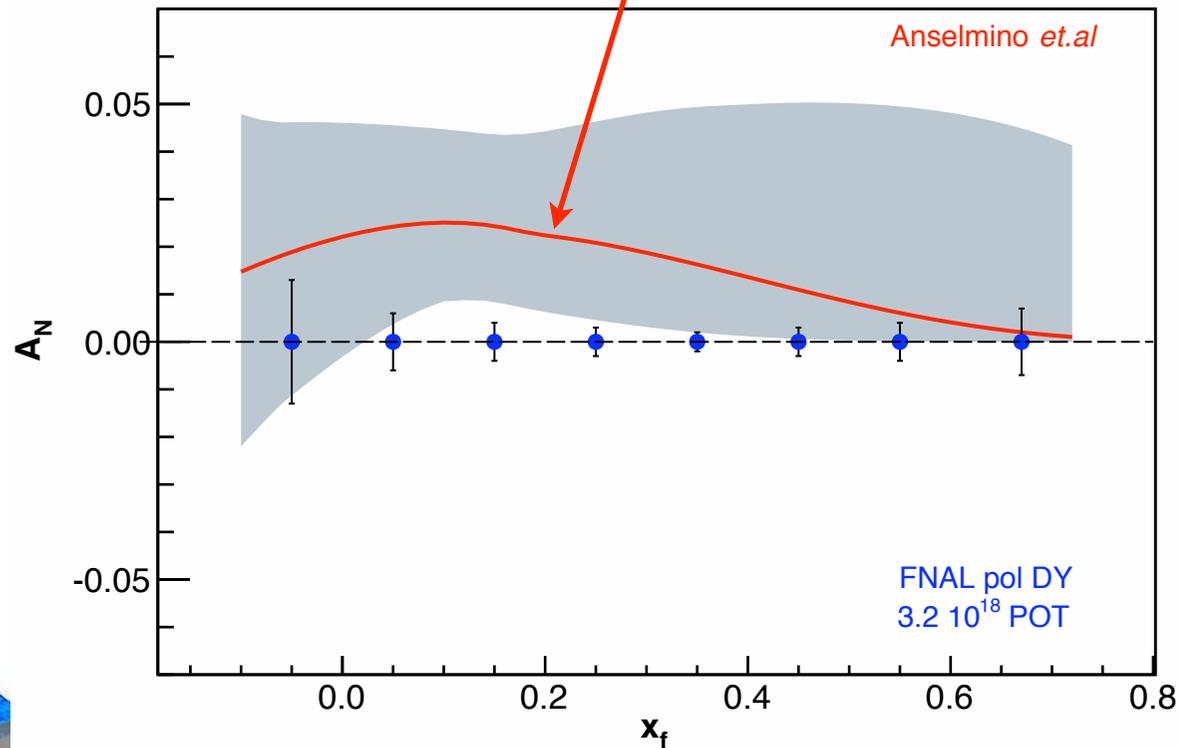
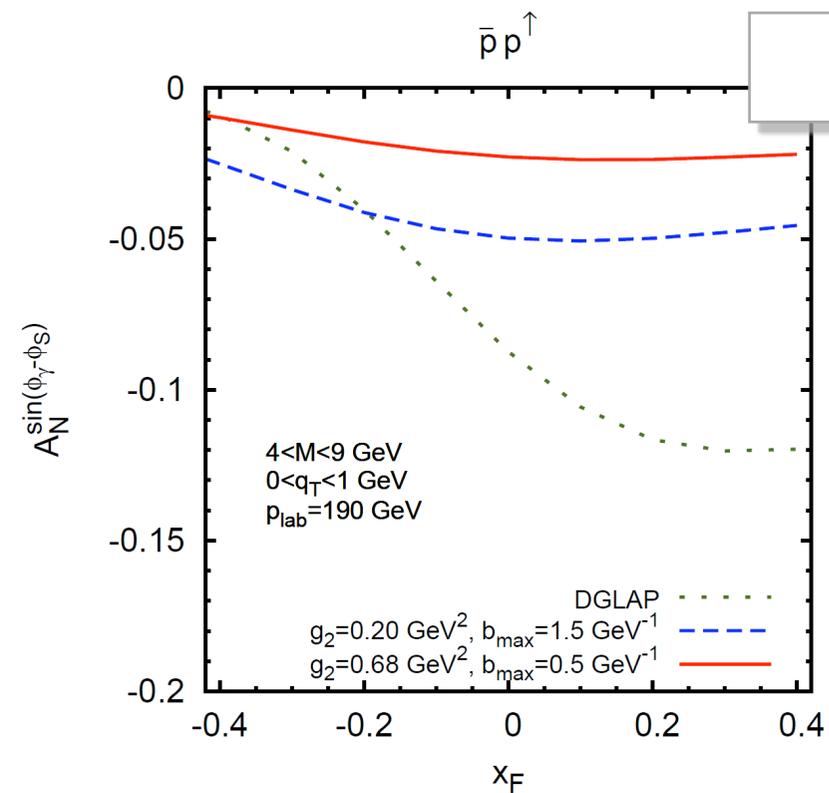
○ : Nucleon Spin    ↔ : Quark Spin

		Quark polarization		
		Unpolarized (U)	Longitudinally Polarized (L)	Transversely Polarized (T)
Nucleon Polarization	U	$f_1 =$		$h_1^\perp =$ - Boer-Mulders
	L		$g_{1L} =$ -	$h_{1L}^\perp =$ -
	T	$f_{1T}^\perp =$ -	$g_{1T} =$ -	$h_1 =$ - Transversity $h_{1T}^\perp =$ - Pretzelosity



This Week

will get smaller



Which dysfunctional care bear are you?



## YOU ARE NIHILIST BEAR!!

*You don't really like anything. Nothing matters enough for you to like it.*

*The only thing you even remotely like is the idea that nothing is worth liking.*

*As the antithesis of the typical Care Bear, you tend to have a lot of existential angst. You're an interesting mix of Goth and a philosopher.*

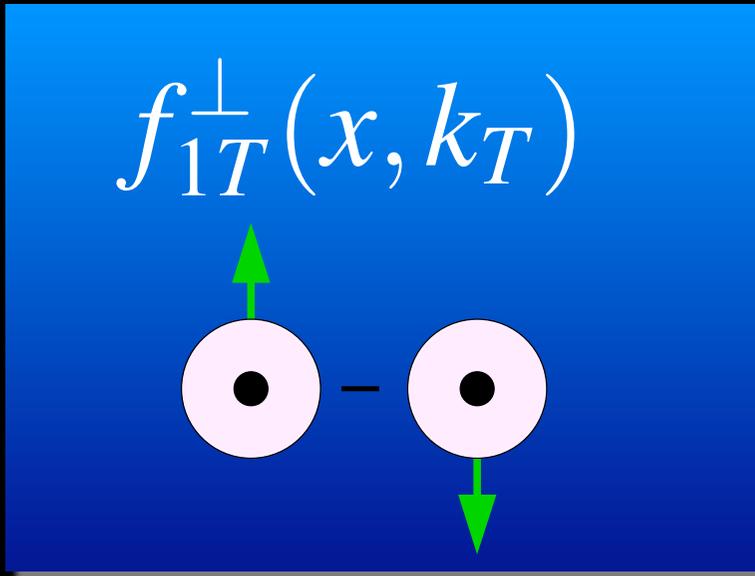
*You're the most intellectual of the Care Bears and can often be found brooding over the state of things.*

*Because of this, you find it very hard to care about things. Even fluffy kitties.*

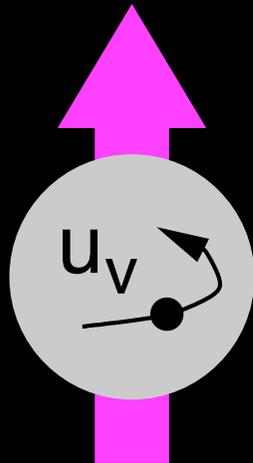
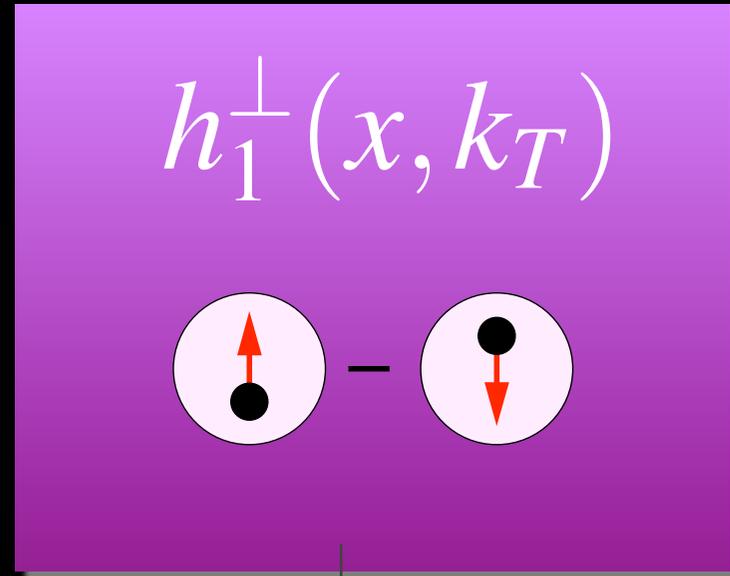
In Search of  $\underline{L}^*$

TMDs + Models give

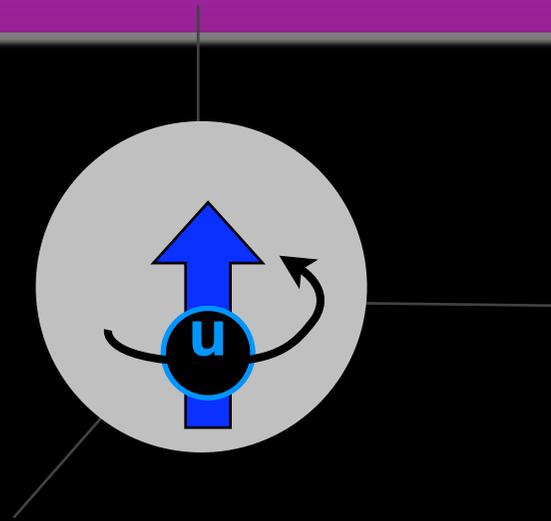
Sivers



Boer-Mulders

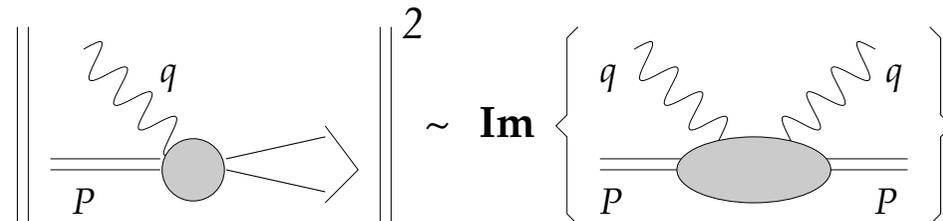


\*undefined  
but beloved

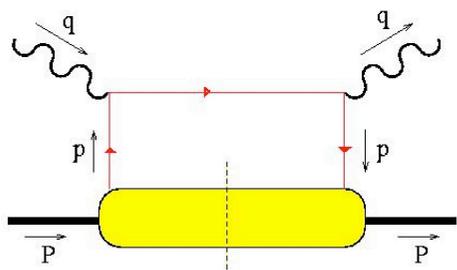


# T-odd TMDs → gauge links and L

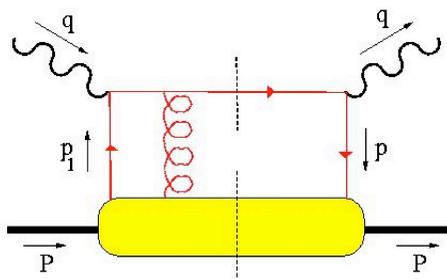
A T-odd function like  $f_{1T}^\perp$  **must** arise from **interference** ... but a distribution function is just a forward scattering amplitude, how can it contain an interference?



## Brodsky, Hwang, & Schmidt 2002



can interfere with

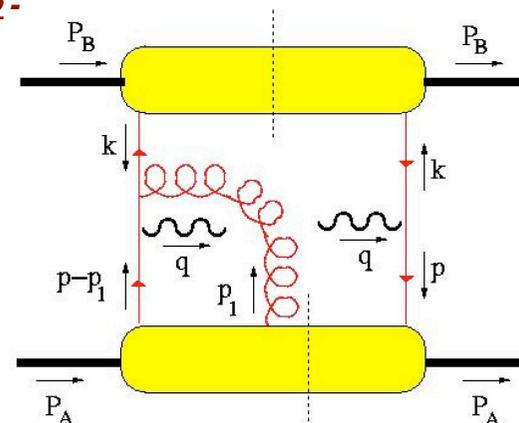


and produce a T-odd effect!  
(also need  $L_z \neq 0$ )

**It looks like higher-twist ... but no, these are soft gluons: “gauge links” required for color gauge invariance**

Such soft-gluon reinteractions with the soft wavefunction are **final / initial state interactions** ... and **process-dependent** ...

e.g. **Drell-Yan**: →  
Sivers effect should have **opposite sign**  
cf. SIDIS

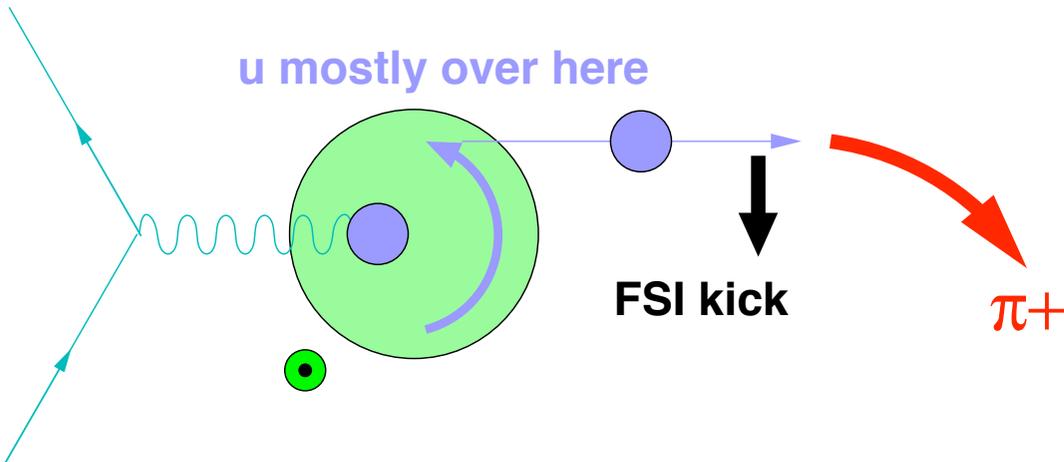


# Phenomenology: Sivers Mechanism

Assuming  $L_u > 0$  ...

## M. Burkardt: Chromodynamic lensing

**Electromagnetic coupling**  $\sim (J_0 + J_3)$  **stronger for *oncoming* quarks**



We observe  $\langle \sin(\phi_h^l - \phi_S^l) \rangle_{\pi^+} > 0$   
 (and opposite for  $\pi^-$ )  
 $\therefore$  for  $\phi_S^l = 0$ ,  $\phi_h^l = \pi/2$  preferred

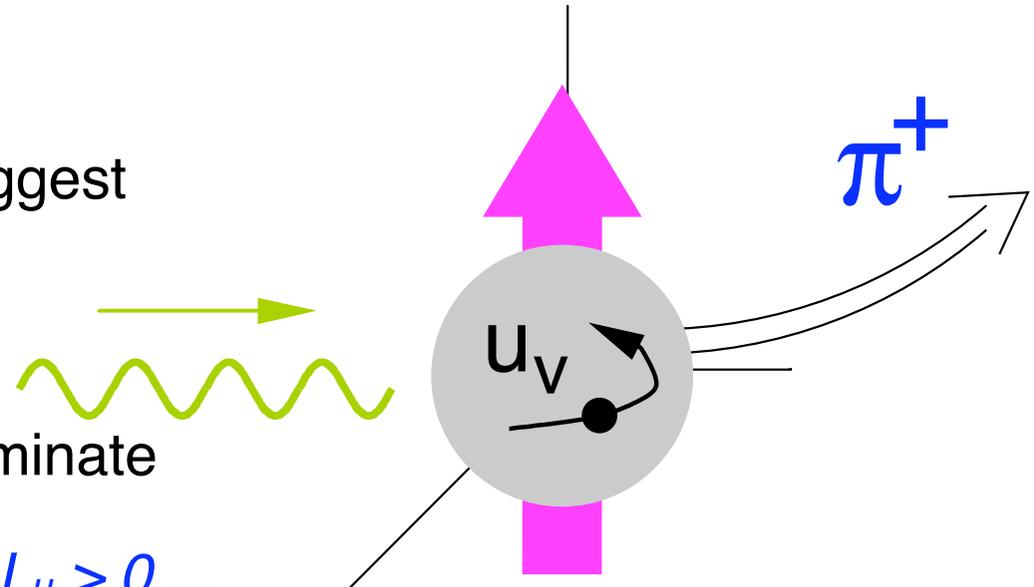
*Model agrees!*

## D. Sivers: Jet Shadowing

Parton energy loss considerations suggest **quenching of jets** from **“near” surface of target**

→ quarks from “far” surface should dominate

Opposite sign to data ... *assuming  $L_u > 0$  ...*

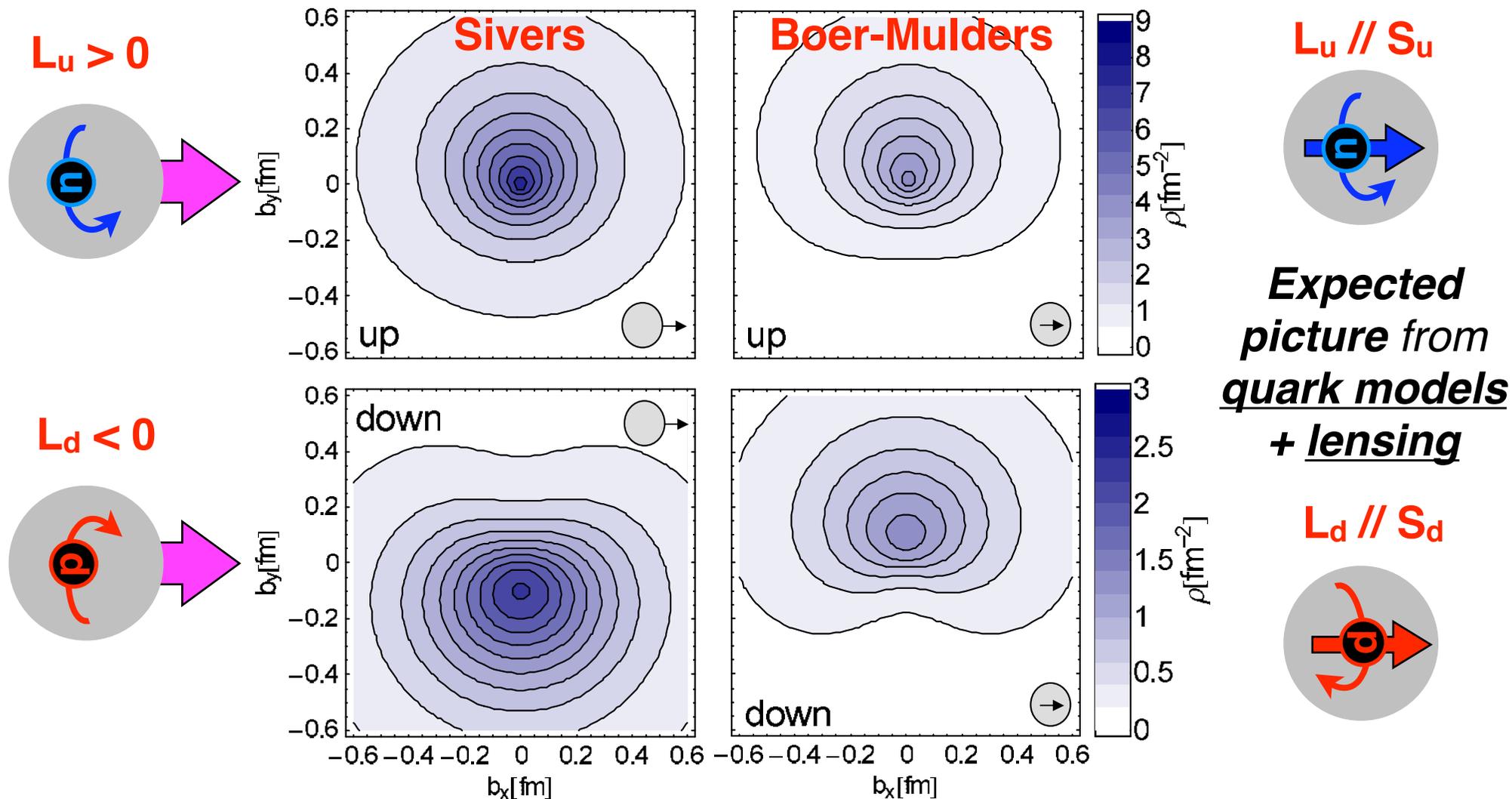


# Transverse spin on the lattice

Hagler et al,  
PRL98 (2007)

Compute **quark densities** in **impact-parameter space** via **GPD formalism**

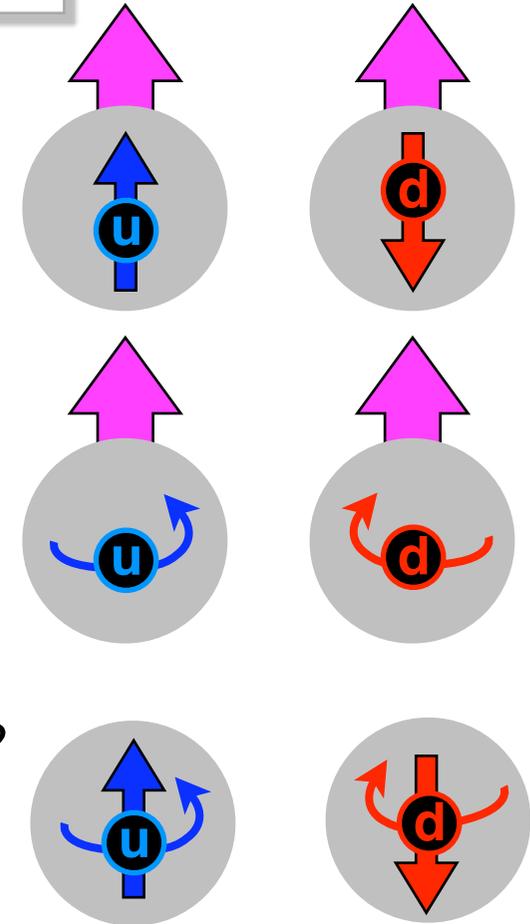
**nucleon coming out of page ... observe spin-dependent *shifts* in quark densities:**



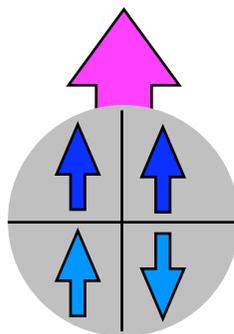
is it a **HAPPY** picture?

## A Tantalizing Picture from SIDIS + Lensing Models

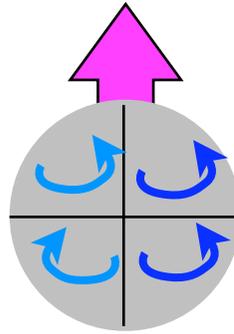
- **Transversity:**  $h_{1,u} > 0$   $h_{1,d} < 0$   
 → same as  $g_{1,u}$  and  $g_{1,d}$  in NR limit
- **Sivers:**  $f_{1T^\perp,u} < 0$   $f_{1T^\perp,d} > 0$   
 → relat<sup>n</sup> to **anomalous magnetic moment\***  
 $f_{1T^\perp,q} \sim \kappa_q$  where  $\kappa_u \approx +1.67$   $\kappa_d \approx -2.03$   
 values achieve  $\kappa^{p,n} = \sum_q e_q \kappa_q$  with  $u,d$  only
- **Boer-Mulders:** follows that  $h_{1^\perp,u}$  and  $h_{1^\perp,d} < 0$  ?  
 → **results** on  $\langle \cos(2\Phi) \rangle_{UU}$  suggest yes:



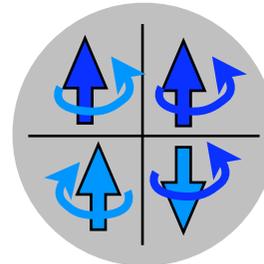
*N.B. these TMDs are all independent*



$$\langle \vec{s}_u \cdot \vec{S}_p \rangle = +0.5$$

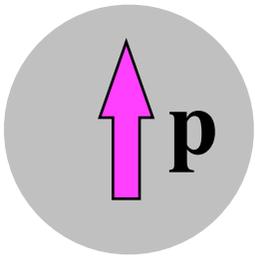


$$\langle \vec{l}_u \cdot \vec{S}_p \rangle = +0.5$$



$$\langle \vec{s}_u \cdot \vec{l}_u \rangle = 0$$

\* Burkardt PRD72 (2005) 094020;  
 Barone et al PRD78 (1008) 045022;



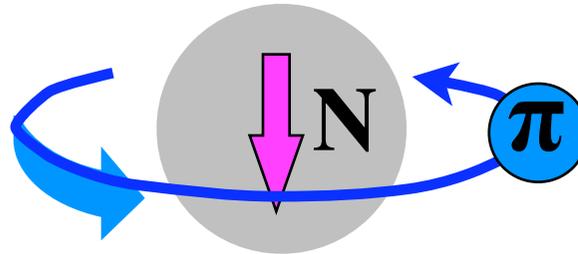
# Meson Cloud on an Envelope → It ORBITS

$|p\rangle = p + N\pi + \Delta\pi + \dots$       Pions have  $J^P = 0^- = \text{negative parity} \dots$   
 → **NEED  $L=1$**  to get proton's  $J^P = 1/2^+$

## $N\pi$ cloud:



$2/3 \quad n \pi^+$   
 $1/3 \quad p \pi^0$

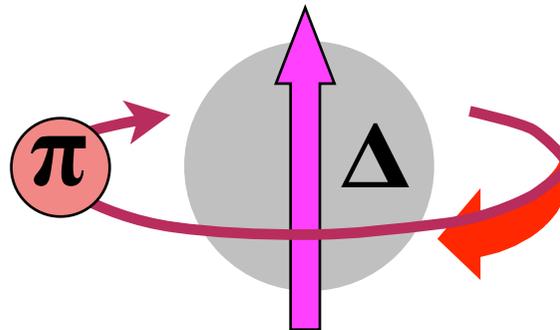


$2/3 \quad L_z = +1$   
 $1/3 \quad L_z = 0$

## $\Delta\pi$ cloud:



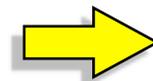
$1/2 \quad \Delta^{++} \pi^-$   
 $1/3 \quad \Delta^+ \pi^0$   
 $1/6 \quad \Delta^0 \pi^+$



$1/2 \quad L_z = -1$   
 $1/3 \quad L_z = 0$   
 $1/6 \quad L_z = +1$

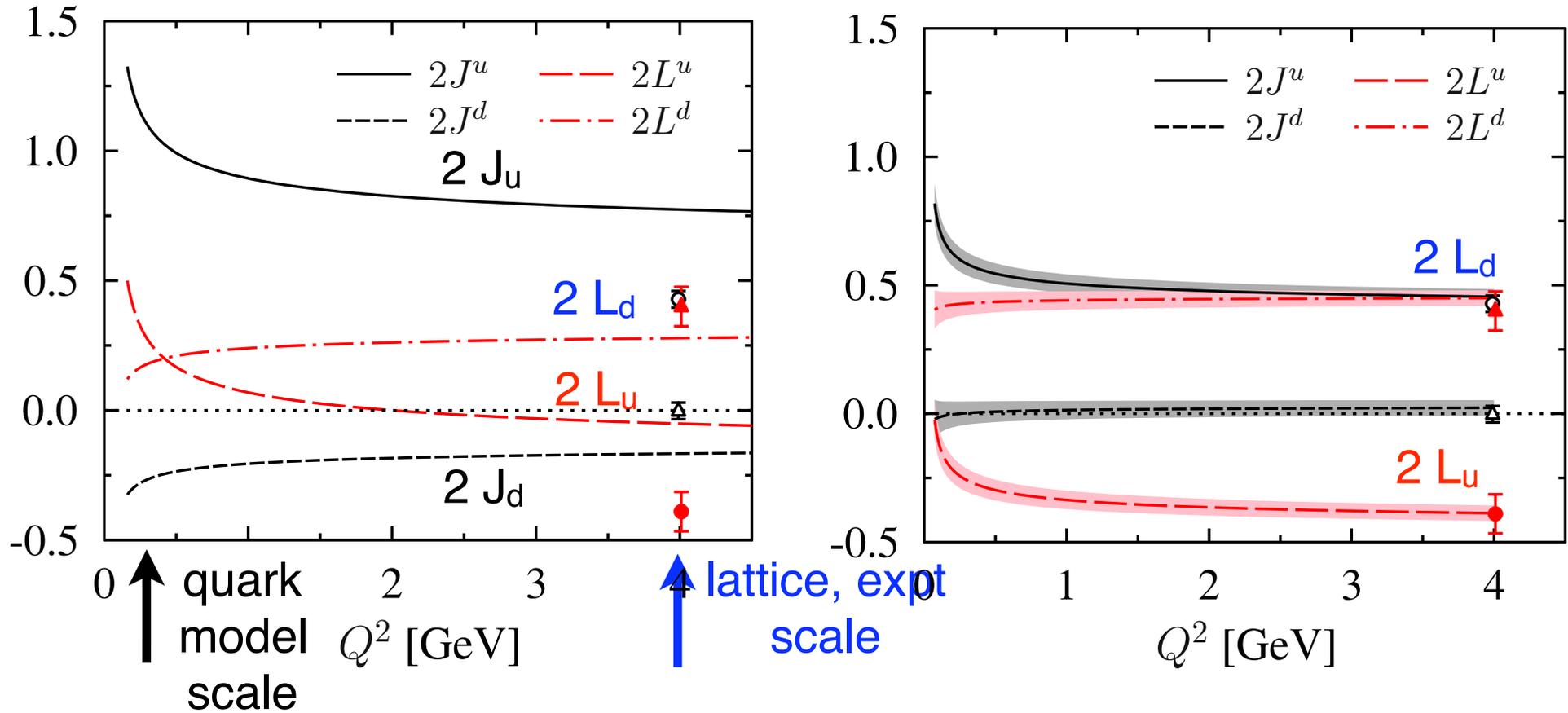
**Dominant source of:**

orbiting u:  $n \pi^+$  with  $L_z(\pi) > 0$   
 orbiting d:  $\Delta^{++} \pi^-$  with  $L_z(\pi) < 0$



$L_u > 0$   
 $L_d < 0$   
 $L_{qbar} \neq 0$

Thomas: **cloudy bag model** evolved up to  $Q^2$  of expt / lattice



→ lattice shows  $L_u < 0$  and  $L_d > 0$  in longitudinal case at expt'al scales!

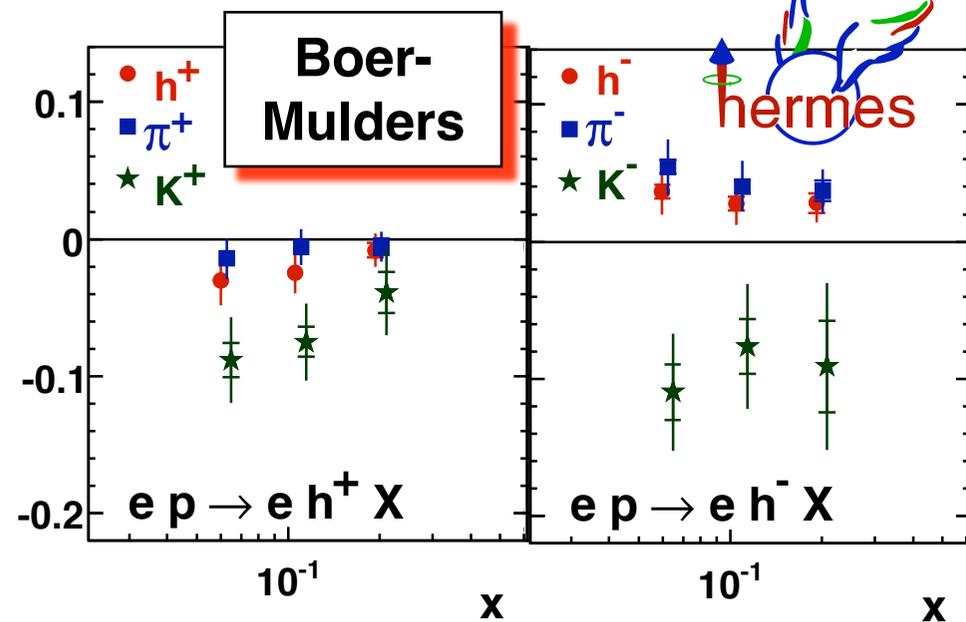
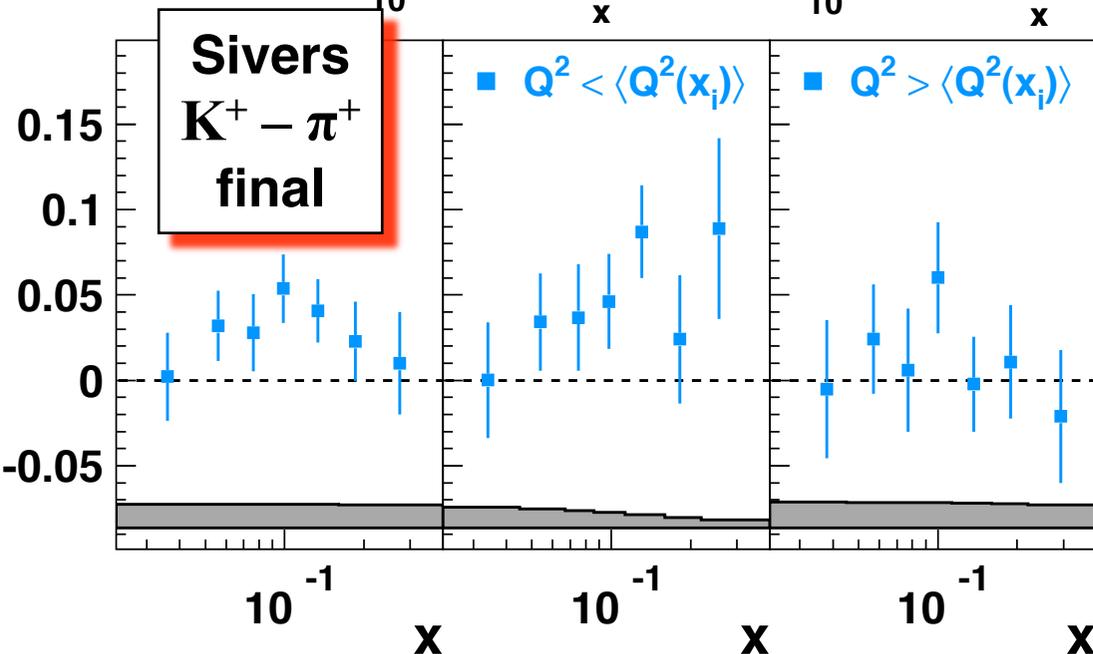
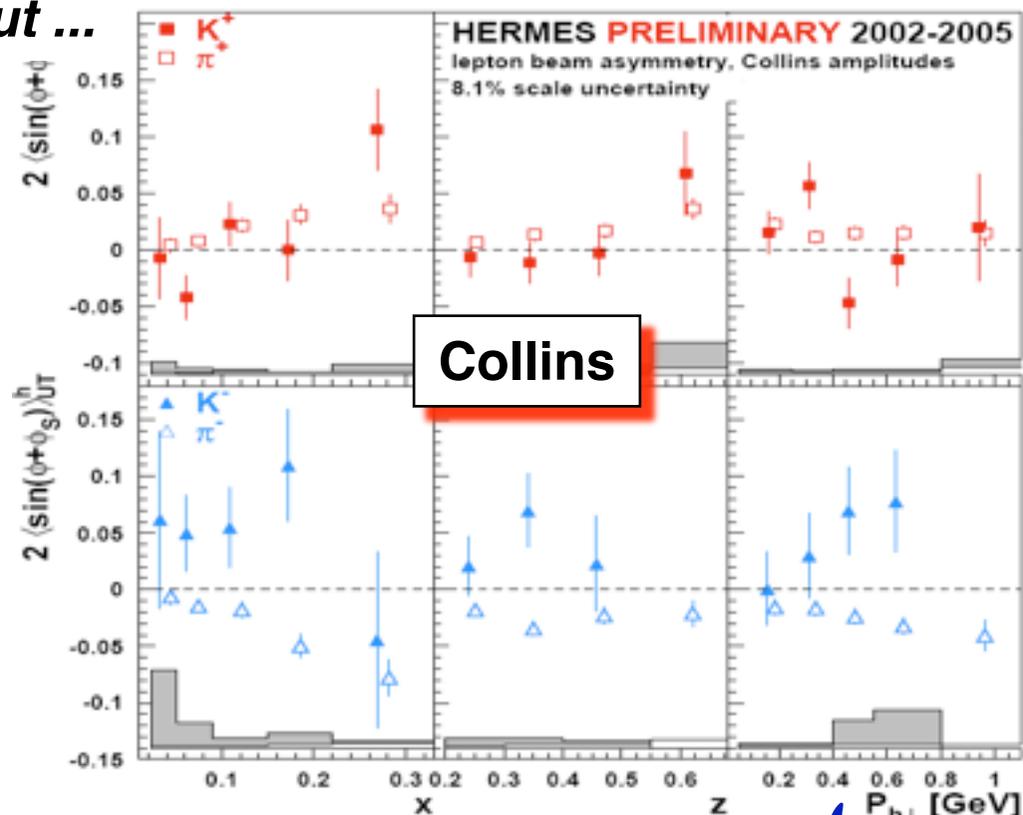
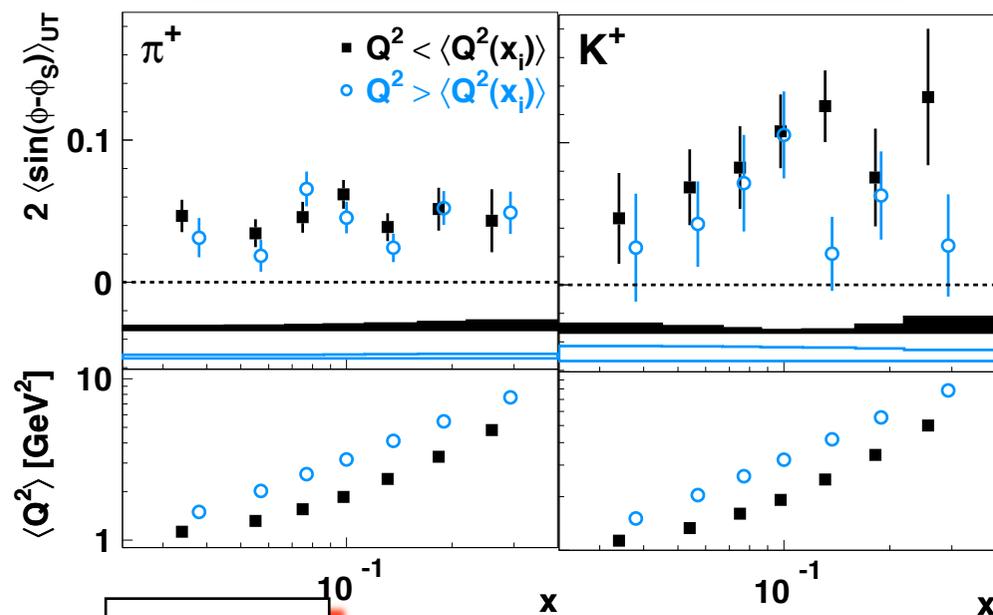
Evolution might explain disagreement with quark models ...

**or not.** Wakamatsu evolves down → insensitive to uncertain scale of quark models

ENTER THE SEA

New **Sivers** fits give  $\approx 0$  for antiquarks, but ...

## The Kaon Collection



... and **BRAHMS** SSA's for kaons, never explained ...

# Flavor-singlet $g_A$

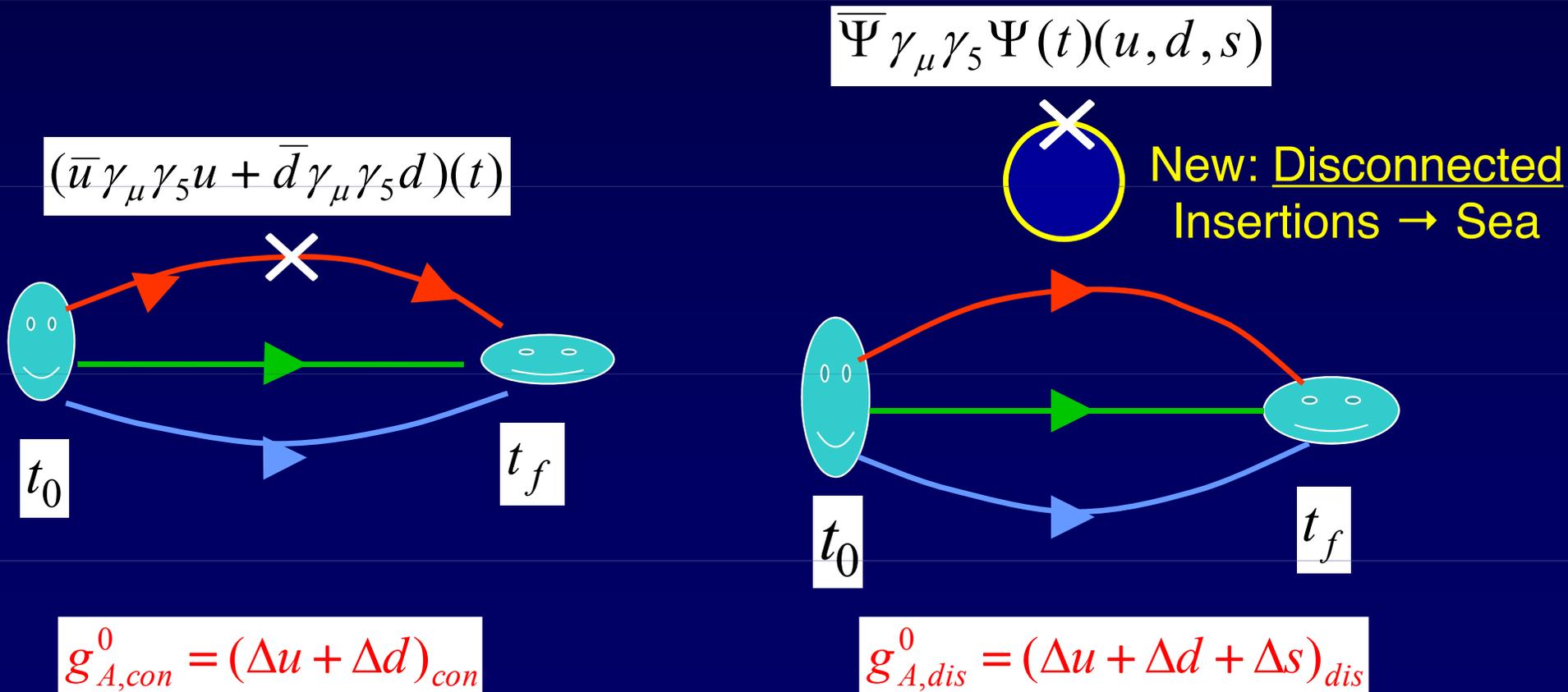
KehFeh Liu,  
INT Workshop, Feb 2012

- Quark spin puzzle (dubbed 'proton spin crisis')

$$g_A^0 = \Delta u + \Delta d + \Delta s = \begin{cases} 1 & \text{NRQM} \\ 0.75 & \text{RQM} \end{cases}$$

– Experimentally (EMC, SMC, ...)

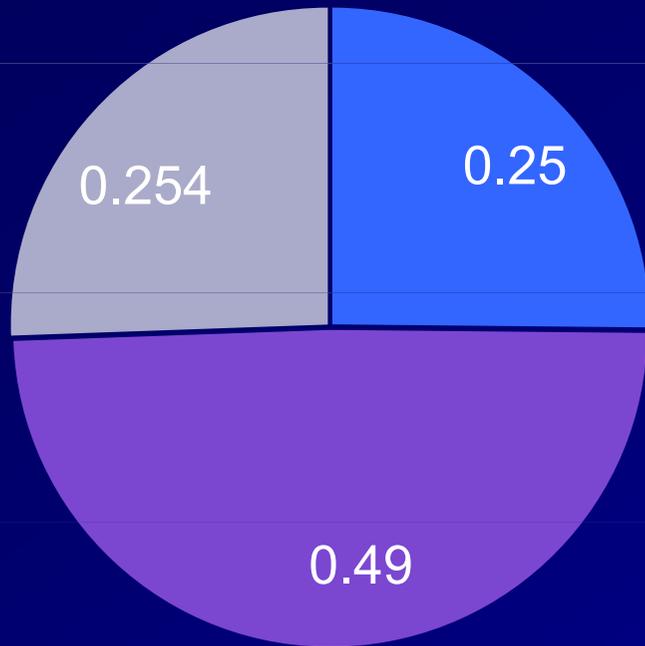
$$\Delta\Sigma = g_A^0 \sim 0.2 - 0.3$$



# Renormalized results:

	CI(u)	CI(d)	CI(u+d)	DI(u/d)	DI(s)	Glue
2J	0.726 (128)	-.072 (82)	0.651 (51)	0.036 (7)	0.023 (7)	0.254 (76)
$g_A$	0.95 (11)	-0.32 (12)	0.65 (8)	-0.12 (1)	-0.12 (1)	
2 L	-0.25 (18)	0.26 (14)	0.00 (10)	0.17 (2)	0.15 (2)	

# Quark Spin, Orbital Angular Momentum, and Glue Angular Momentum



**2 J**

- Quark Spin
- Quark OAM
- Glue AM

**The Sea is Orbiting!**

$$\Delta q \approx 0.25;$$

$$2 L_q \approx 0.49 \text{ (0.0(CI) + 0.49(DI));}$$

$$2 J_g \approx 0.25$$

## Proton Spin Decompositions

$$\mathbf{J}^{\text{Ji}} = \underbrace{\frac{i}{2} q^\dagger (\vec{r} \times \vec{D})^z q}_{L_q} + \underbrace{\frac{1}{2} q^\dagger \sigma^z q}_{\Delta q} + \underbrace{2 \text{Tr} E^j (\vec{r} \times \vec{D})^z A^j}_{L_g} + \underbrace{\text{Tr} (\vec{E} \times \vec{A})^z}_{\Delta g}$$

$$\mathbf{J}^{\text{Jaffe}} = \frac{1}{2} q_+^\dagger (\vec{r} \times i \vec{\nabla})^z q_+ + \frac{1}{2} q_+^\dagger \gamma_5 q_+ + 2 \text{Tr} F^{+j} (\vec{r} \times i \vec{\nabla})^z A^j + \epsilon^{+-ij} \text{Tr} F^{+i} \vec{A}^j$$

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**Ji:** ③ gauge invariant  $\Delta q, L_q, J_g$

✗ **access  $\Delta g$ :** no GI sep<sup>n</sup> of  $\Delta g, L_g$

✓ **measure  $L_q$**  (expt & lattice):  
yes → via GPDs & DVCS

✗ **interpret  $L_q$ :** covariant derivative  
 $D^\mu = \partial^\mu + ig^\mu \leftarrow$  **gluon interac's**

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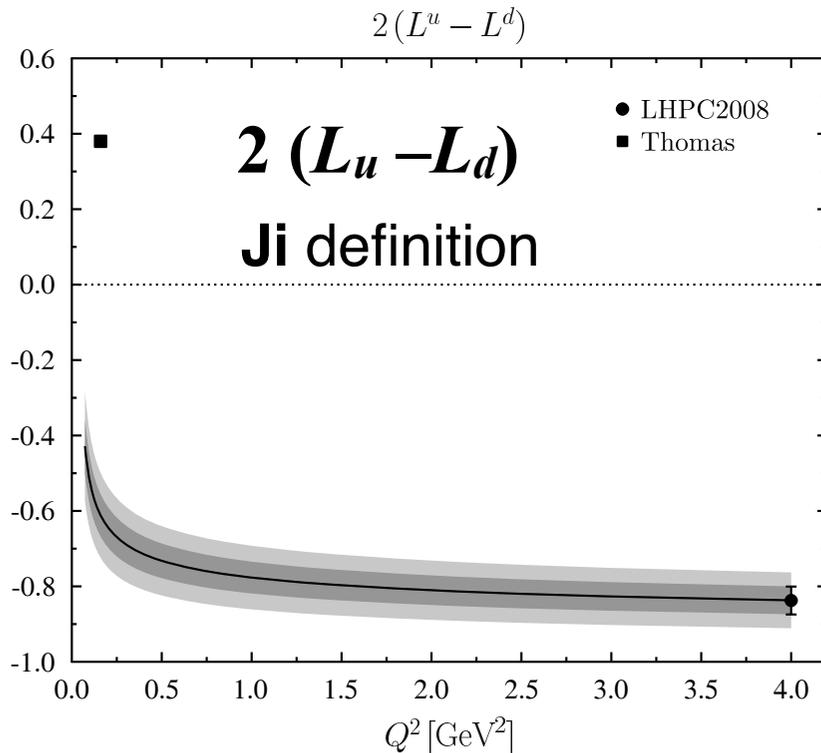
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see ongoing work of **Wakamatsu** PRD 81 (2010), 83 (2011)  
& **Chen et al** PRL 100 (2008), 103 (2009)

Theory: Ji's  $L_{u-d}$  is rock-solid & **negative**



$$2L_q^{\text{Ji}} = \left[ \langle x \rangle_q + E_q^{(2)} \right]_{=J_q} - \Delta q$$

- $\langle x \rangle_{u-d}$ : well known
- $\Delta u - \Delta d = g_A$ : well known
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**and XQSM** agree

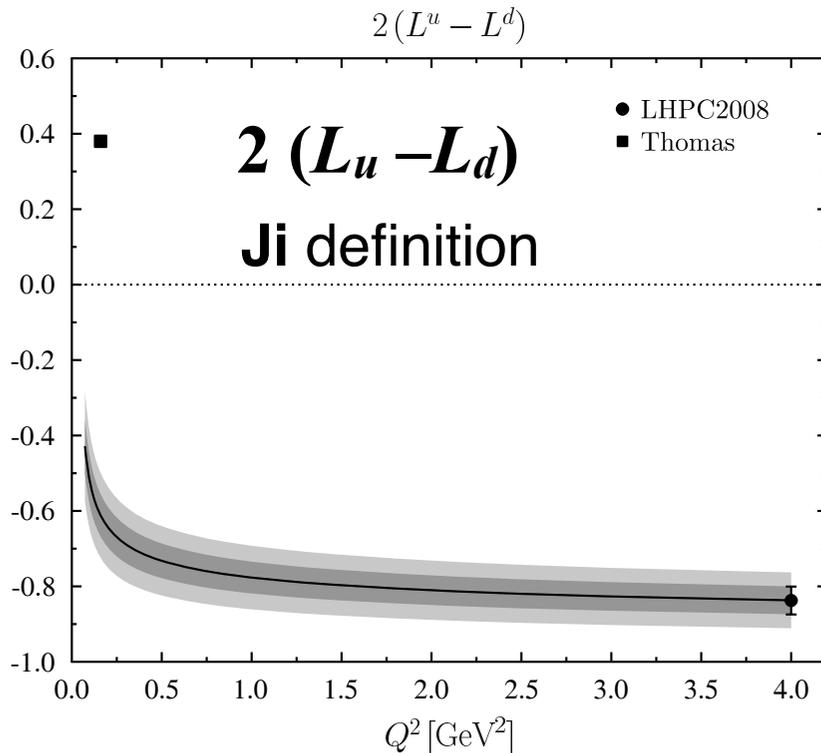
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Wakamatsu, EPJA 44 (2010)

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Compare Jaffe & Ji

calculate explicitly in  $\chi$ QSM;  
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	$L_{u-d}$ Jaffe	$L_{u-d}$ Ji
<b>Valence</b>	+0.147	-0.142
<b>Sea</b>	-0.265	-0.188
<b>Total</b>	-0.115	-0.330

Negative model value  
 dominated by sea quark L !

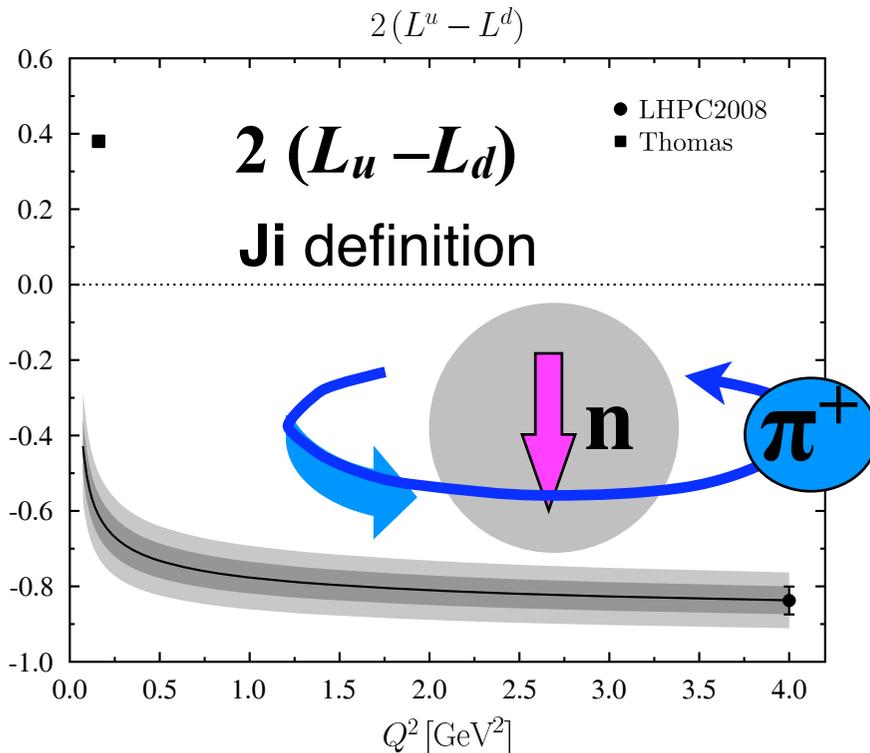
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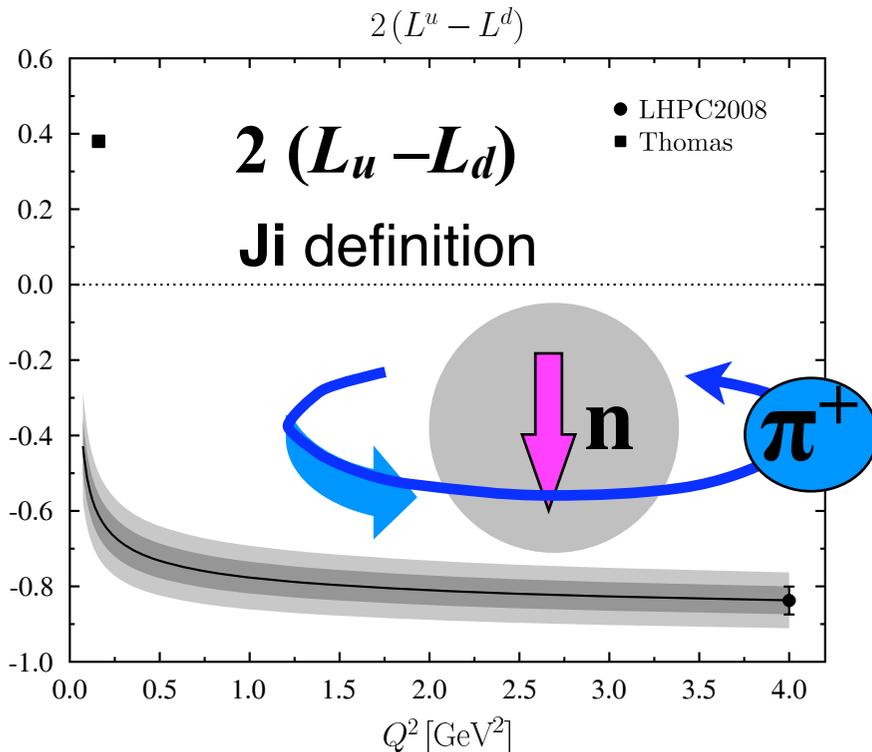
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**Negative model value  
dominated by sea quark L !**

**Need direct measurement of  
Sivers for sea quarks:**

**Spin-dependent Drell-Yan  
with  $p$  or  $\pi^+$  beam & pol'd target**

# TMD Evolution

# Turin standard approach (DGLAP)

- Unpolarized TMDs are factorized in  $x$  and  $k_{\perp}$ . Only the collinear part evolves with DGLAP evolution equation. No evolution in the transverse momenta:

$$\hat{f}_{q/p}(x, k_{\perp}; Q) = f_{q/p}(x; Q) \frac{e^{-k_{\perp}^2 / \langle k_{\perp}^2 \rangle}}{\pi \langle k_{\perp}^2 \rangle}$$

Collinear PDF (DGLAP evolution)

Normalized Gaussian: no evolution

# TMD evolution formalism\*

- \* *J.C. Collins, Foundation of Perturbative QCD, Cambridge Monographs on Particle Physics, Nuclear Physics and Cosmology, No. 32, Cambridge University Press, 2011.*
- *S. M. Aybat and T. C. Rogers, Phys. Rev. D83, 114042 (2011), arXiv:1101.5057 [hep-ph]*
- *S. M. Aybat, J. C. Collins, J.-W. Qiu and T.C. Rogers, arXiv:1110.6428 [hep-ph]*



**Re<sup>e</sup>volution!**

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# TMD evolution formalism

➤ At LO the evolution equation can be summarized by the following expression:


$$\tilde{F}(x, \mathbf{b}_T; Q) = \tilde{F}(x, \mathbf{b}_T; Q_0) \tilde{R}(Q, Q_0, b_T) \exp \left\{ -g_K(b_T) \ln \frac{Q}{Q_0} \right\}$$

**Output** function at the scale  $Q$   
in the impact parameter space

**Input** function at the scale  $Q_0$   
in the impact parameter space

Evolution kernel

# TMD evolution formalism

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- **Non Perturbative** (scale independent) part of the evolution kernel that needs to be empirically modeled

$$g_K(b_T) = \frac{1}{2} g_2 b_T^2$$

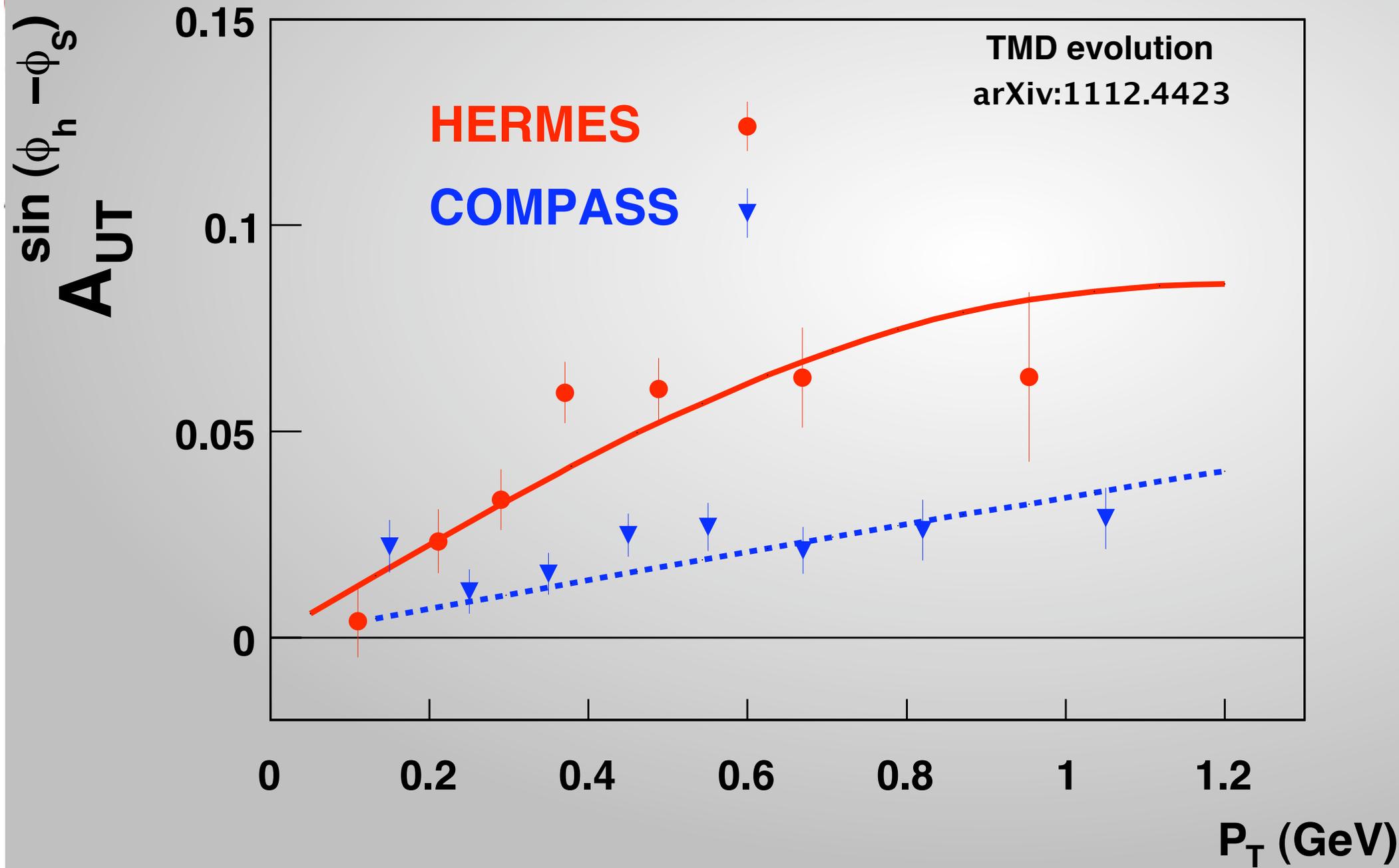
$$g_2 = 0.68 \text{ GeV}^2$$

Common choice used in the unpolarized DY data analyses in the CSS formalism

Landry et al. Phys Rev D67, 073016

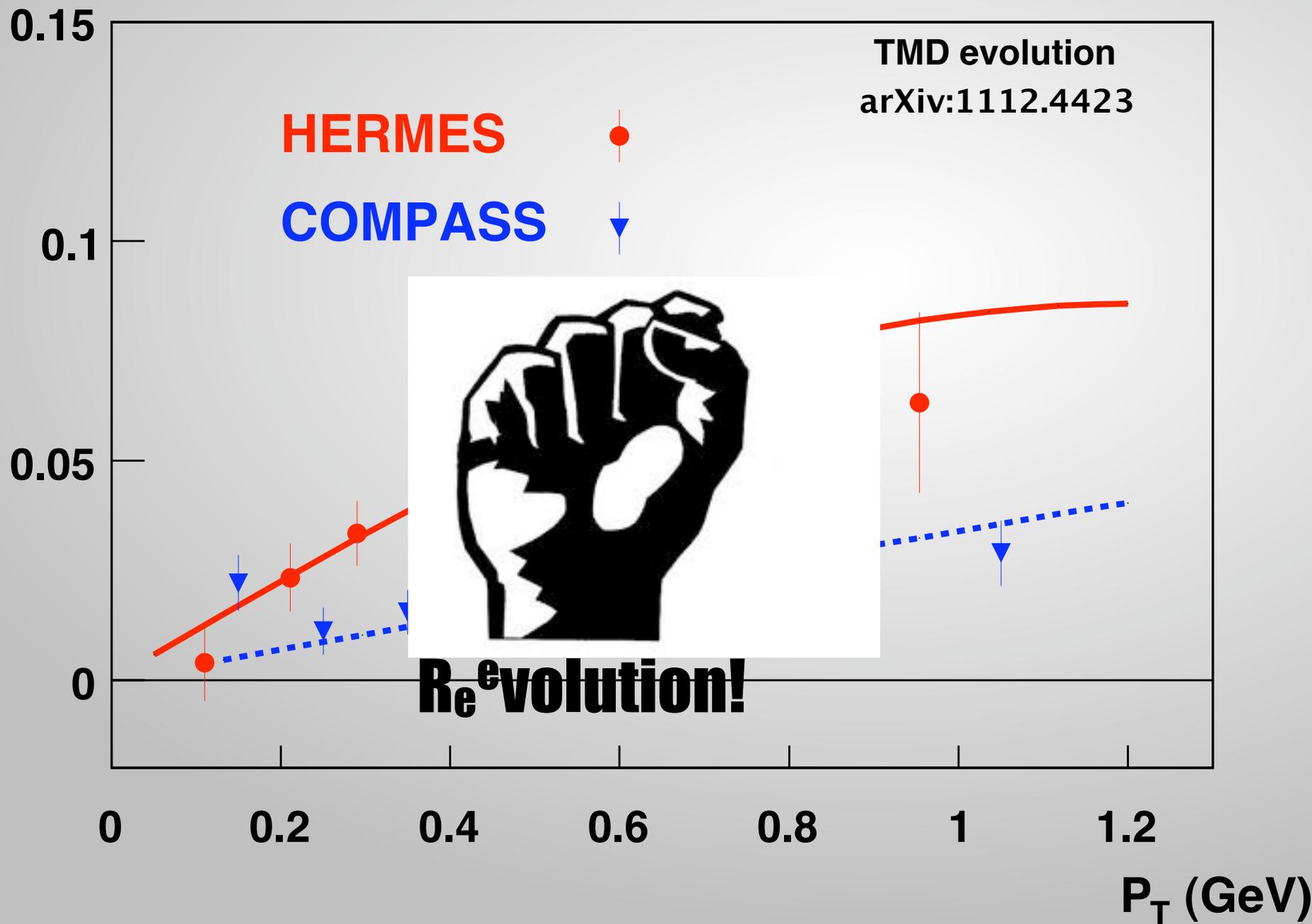
**Crucial parameters:  
 $g_2$  and  $b_{T\text{max}}$**

# Great Success in bridging HERMES-COMPASS!



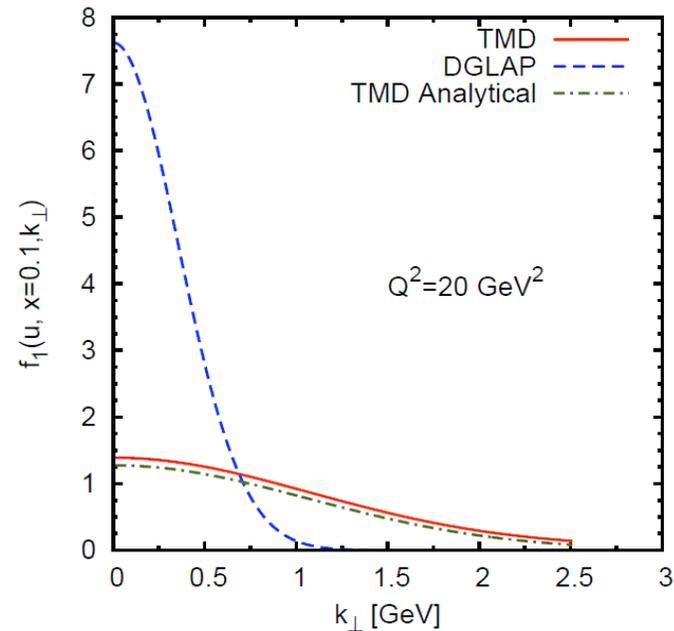
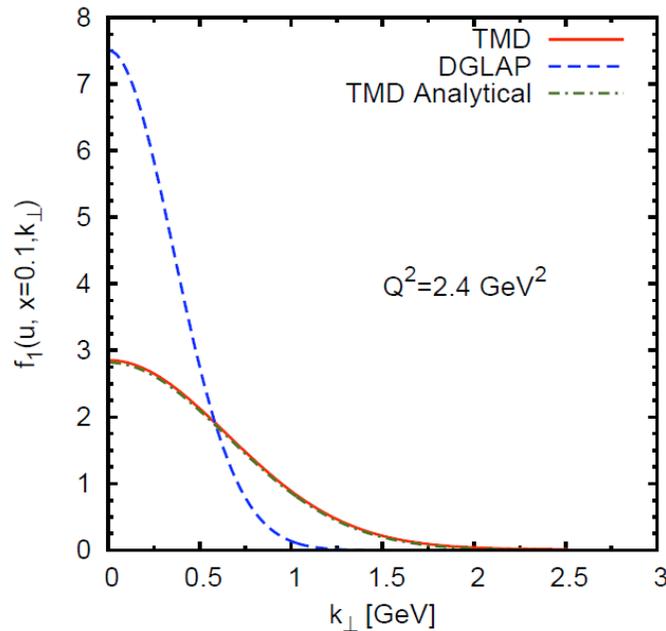
# Great Success in bridging HERMES-COMPASS!

$\sin(\phi_h - \phi_s)$   
 $A_{UT}$



# Comparative analysis of TMD evolution equations

**f1: huge effect!**



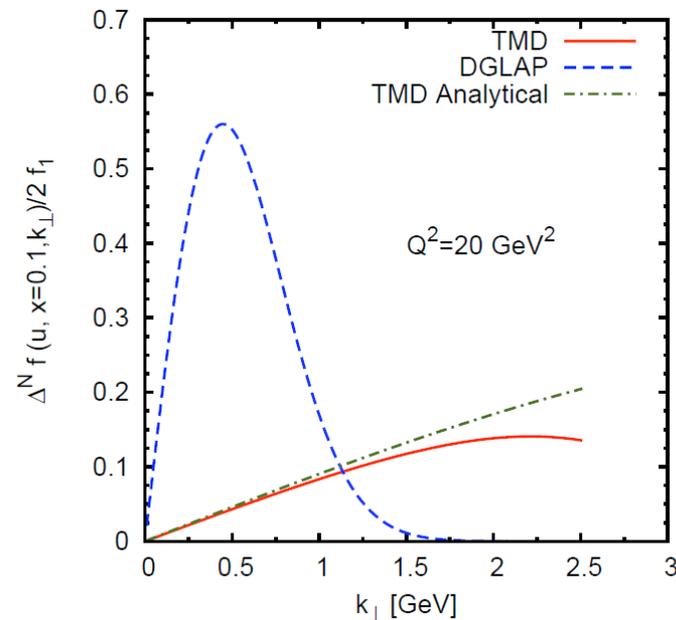
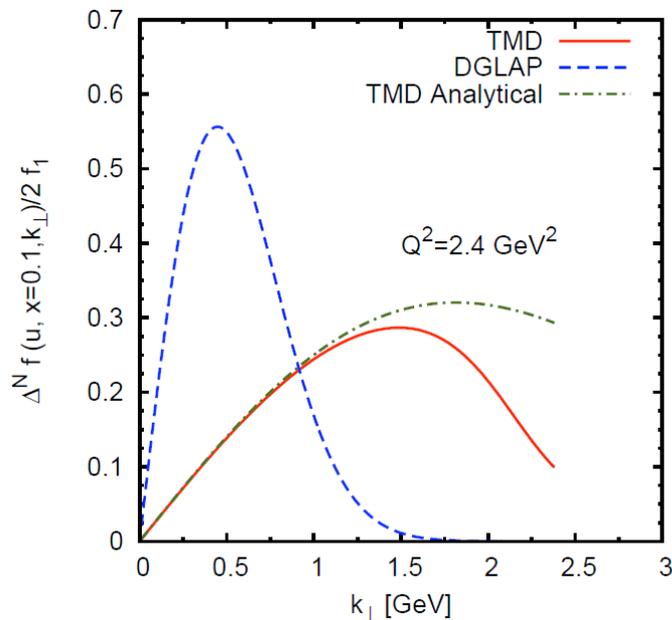
Starting scale  $Q_0 = 1 \text{ GeV}$   
Same function at  $Q_0$

DGLAP evolution is slow at moderate  $x$  and in this range of  $Q^2$

For the unpolarized PDF, the analytical approximation holds up to large  $k_\perp$

# Comparative analysis of TMD evolution equations

**Sivers:  
huge effect!**



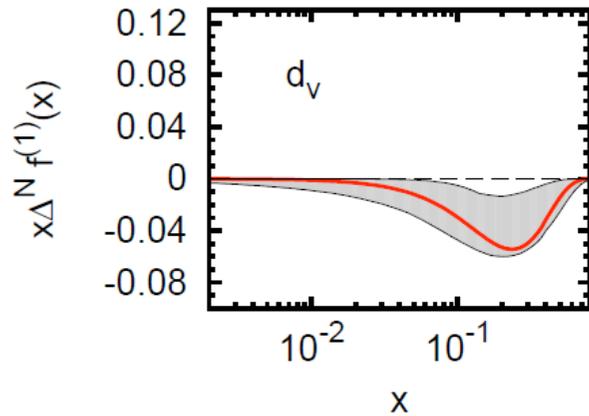
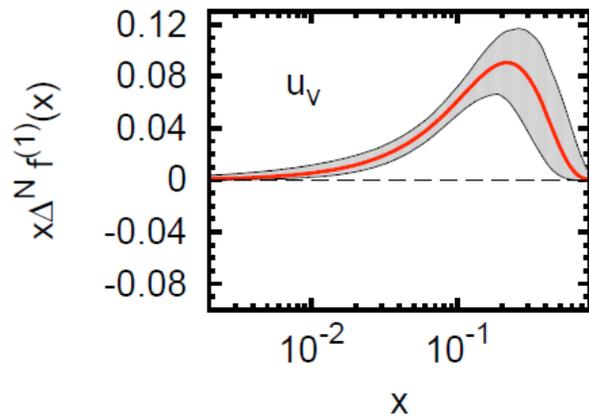
Starting scale  $Q_0 = 1 \text{ GeV}$   
Same function at  $Q_0$

For the Sivers function,  
the analytical approximation  
breaks down at large  $k_\perp$  values

# BIGGER Siverson function from SIDIS!

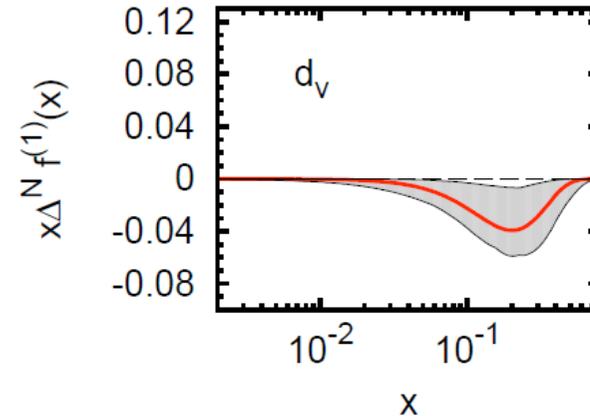
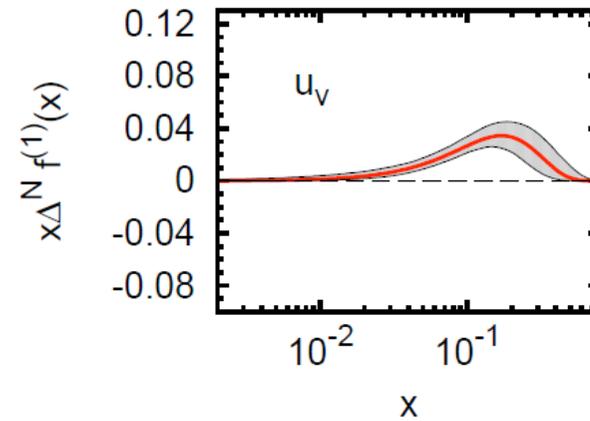
## Fit of HERMES and COMPASS SIDIS data

### TMD Evolution



$Q_0 = 1 \text{ GeV}$

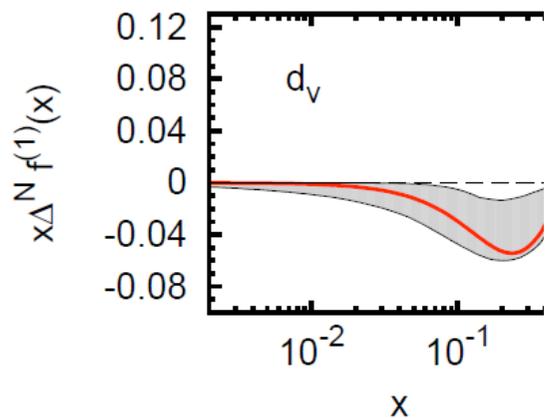
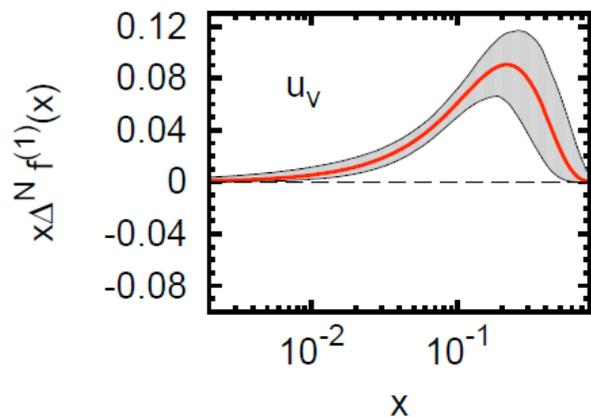
### DGLAP Evolution



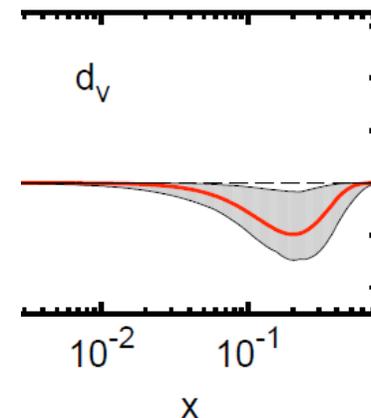
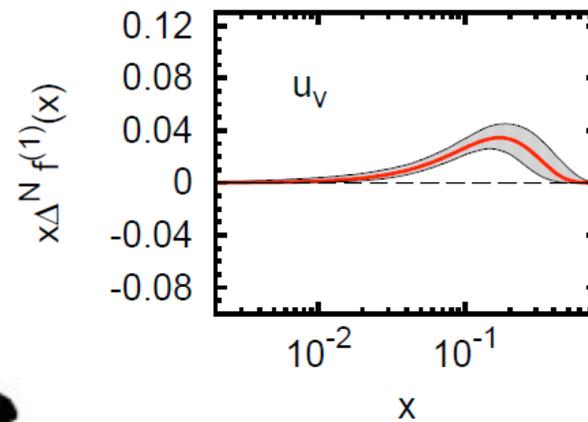
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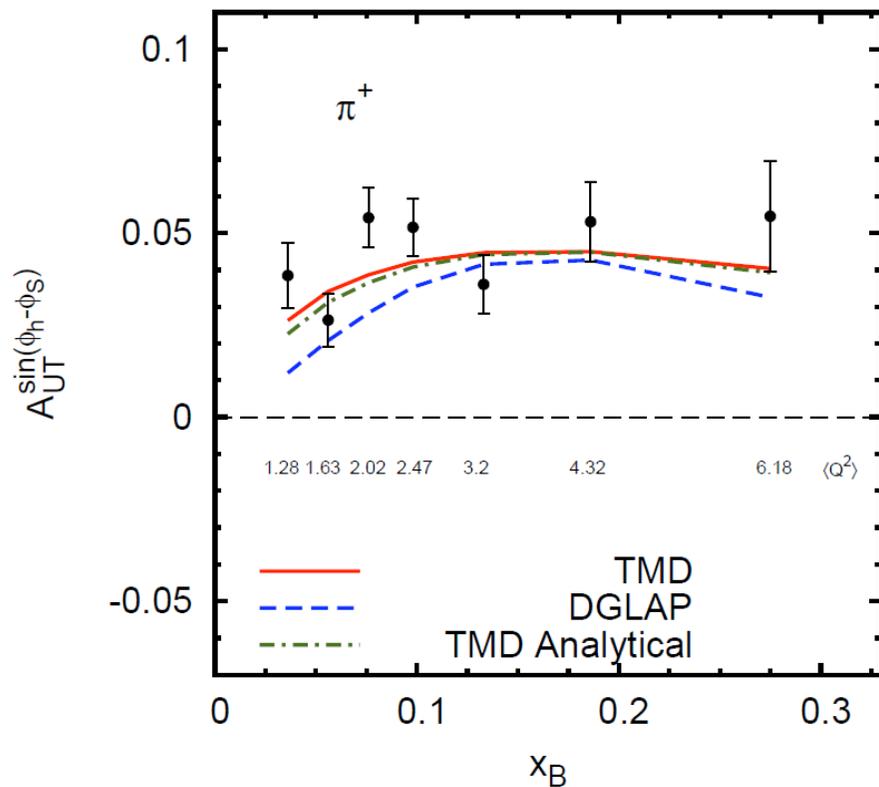


**Re<sup>e</sup>volution!**

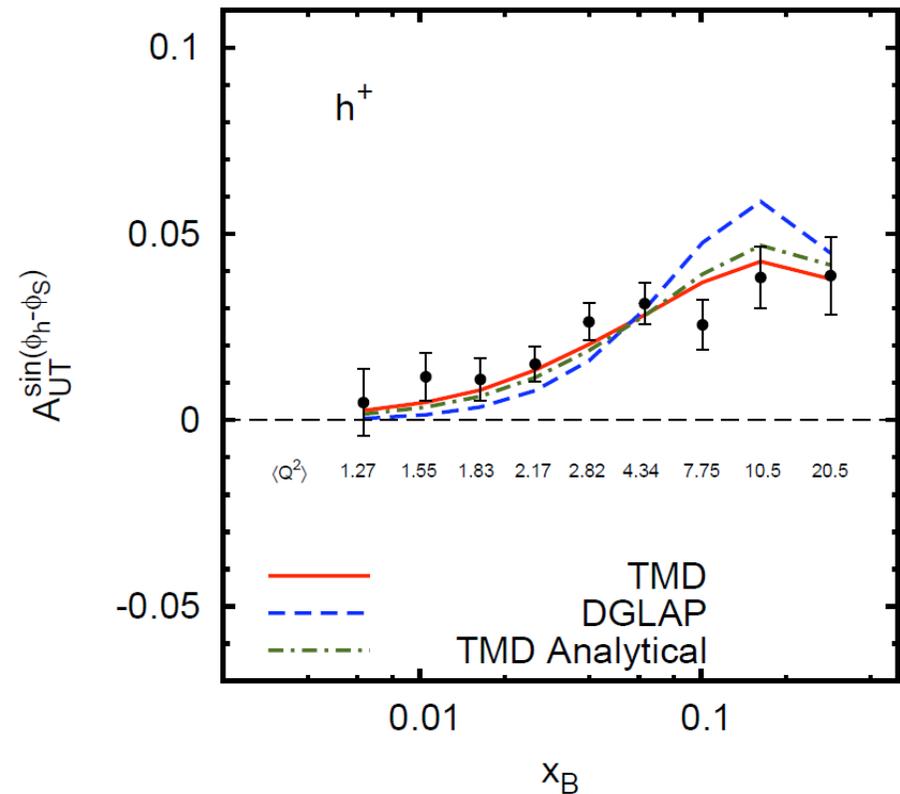
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**Sivers Asym: barely change in SIDIS**

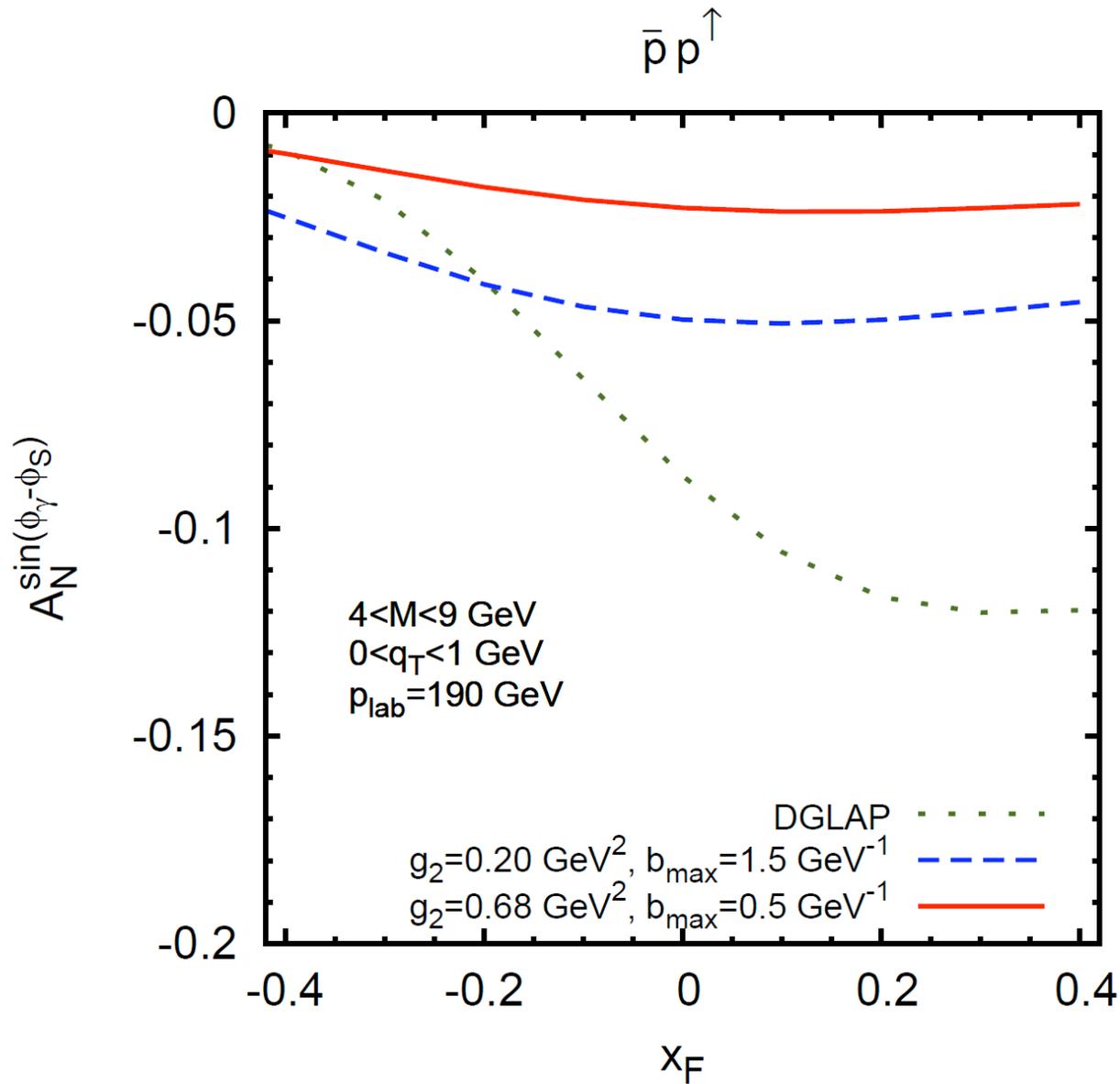
HERMES PROTON



COMPASS PROTON

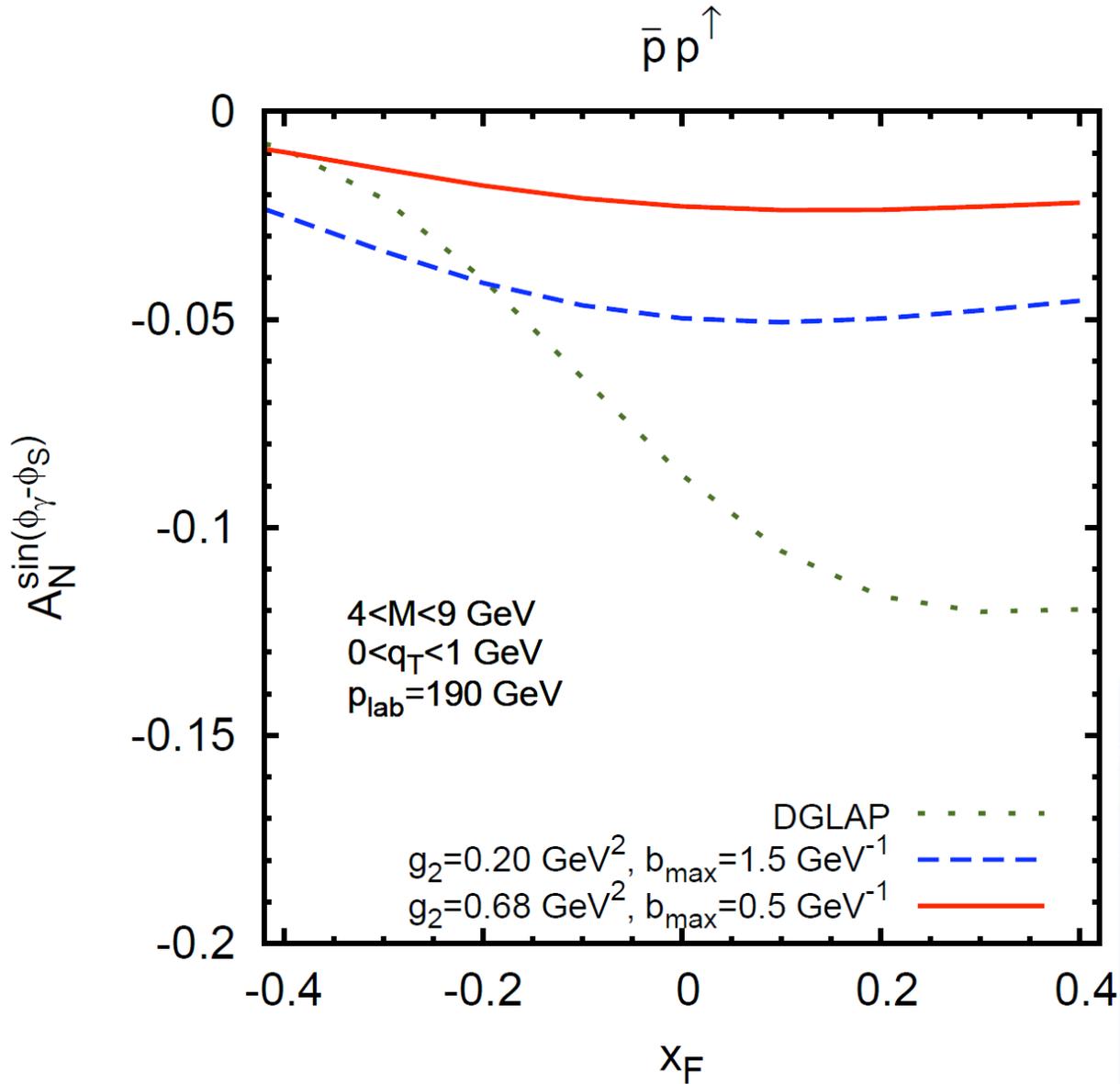


# Sivers Asym: huge change in DY!



- smaller
- very dependent on  $g_2, b_{\text{max}}$

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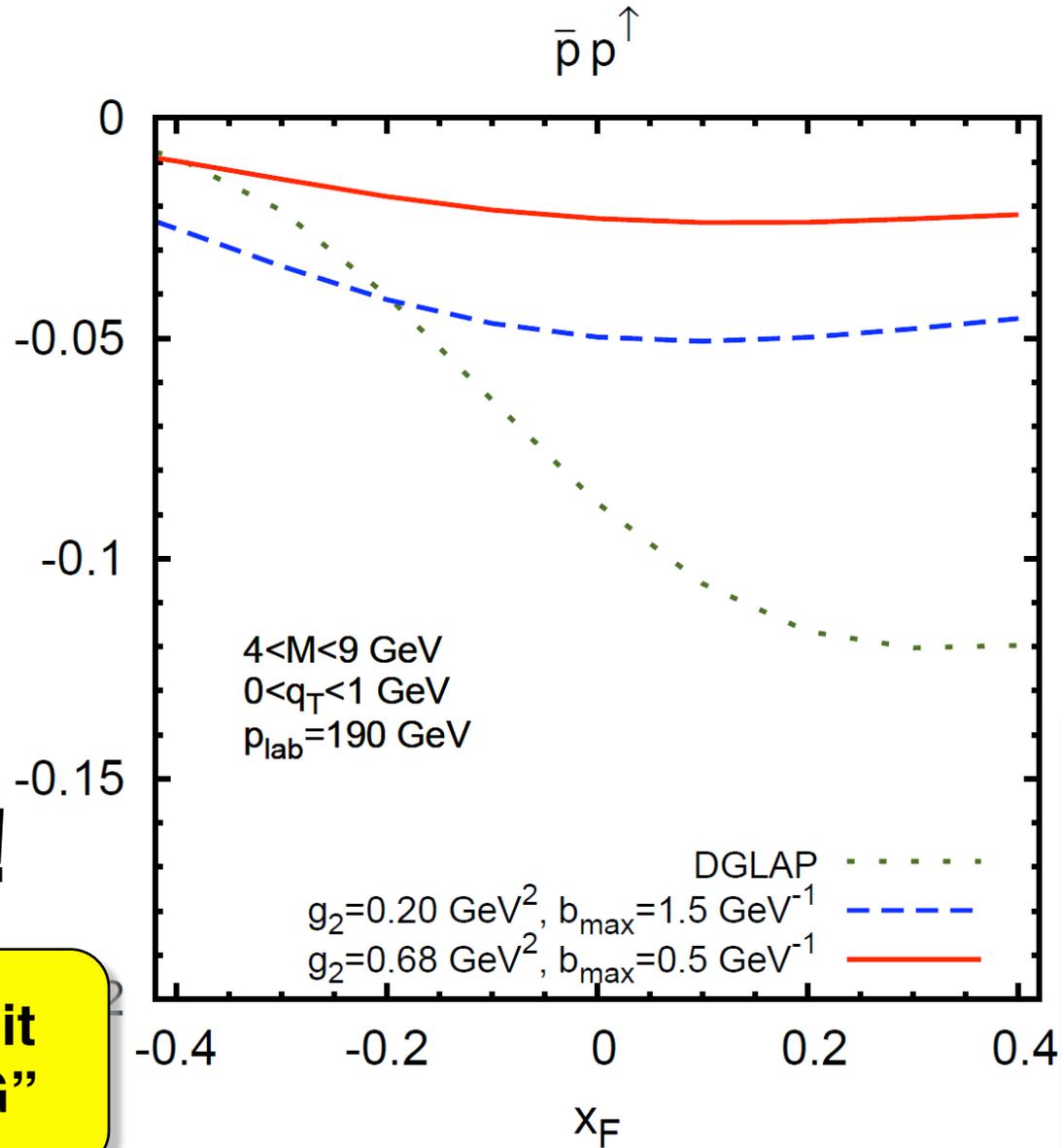


- smaller
- very dependent on  $g_2, b_{\text{max}}$



**Grumpy Bear**

# Sivers Asym: huge change in DY!



- smaller
- very dependent on  $g_2, b_{\text{max}}$



**Re<sup>e</sup>volution!**

**Melis: "Revisit EVERYTHING"**



**Grumpy Bear**

$\langle k^2_{\perp} \rangle$ , Multi-D, and  
Global Fits

# Boer-Mulders function in DY from fits

- ▶ Can we safely assume that the average transverse momentum is the same in SIDIS and in DY?

 Gaussian smearing for unpolarized PDFs

- $f_{q/p}(x, k_{\perp}) = f_q(x) \frac{1}{\pi \langle k_{\perp}^2 \rangle} e^{-k_{\perp}^2 / \langle k_{\perp}^2 \rangle}$

From SIDIS:  $\langle k_{\perp}^2 \rangle = 0.25 \text{ (GeV/c)}^2$

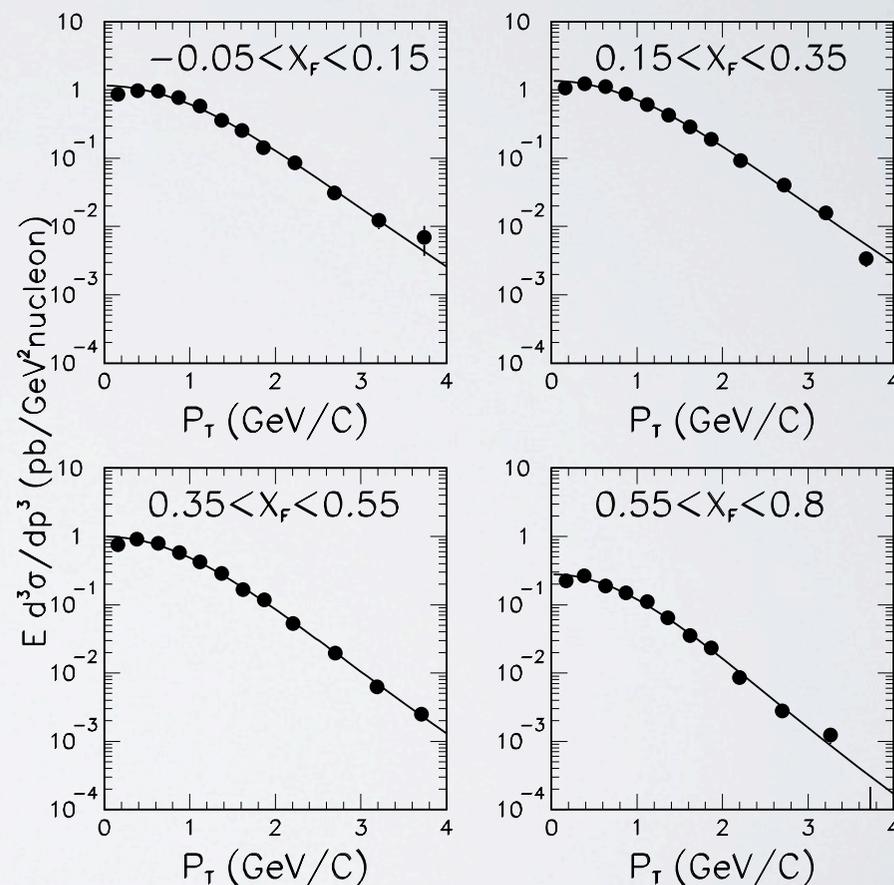
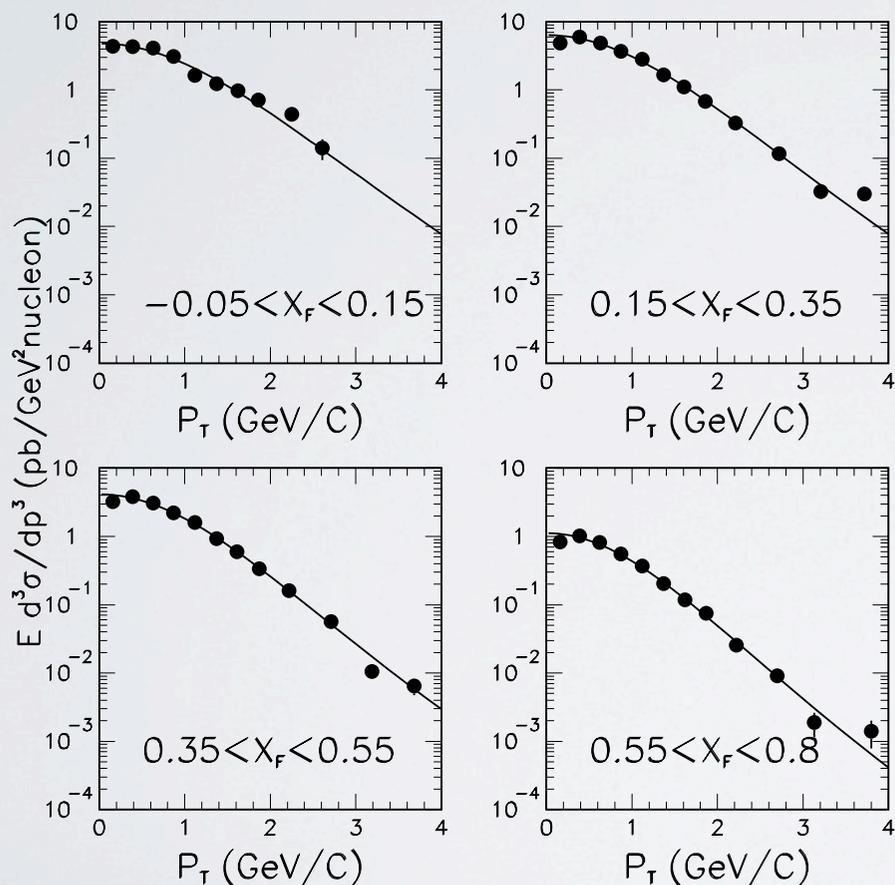
Typical DY :  $\langle k_{\perp}^2 \rangle \simeq 0.5 - 1 \text{ (GeV/c)}^2$

→ Let us try to change this value

# Multidim. studies are needed

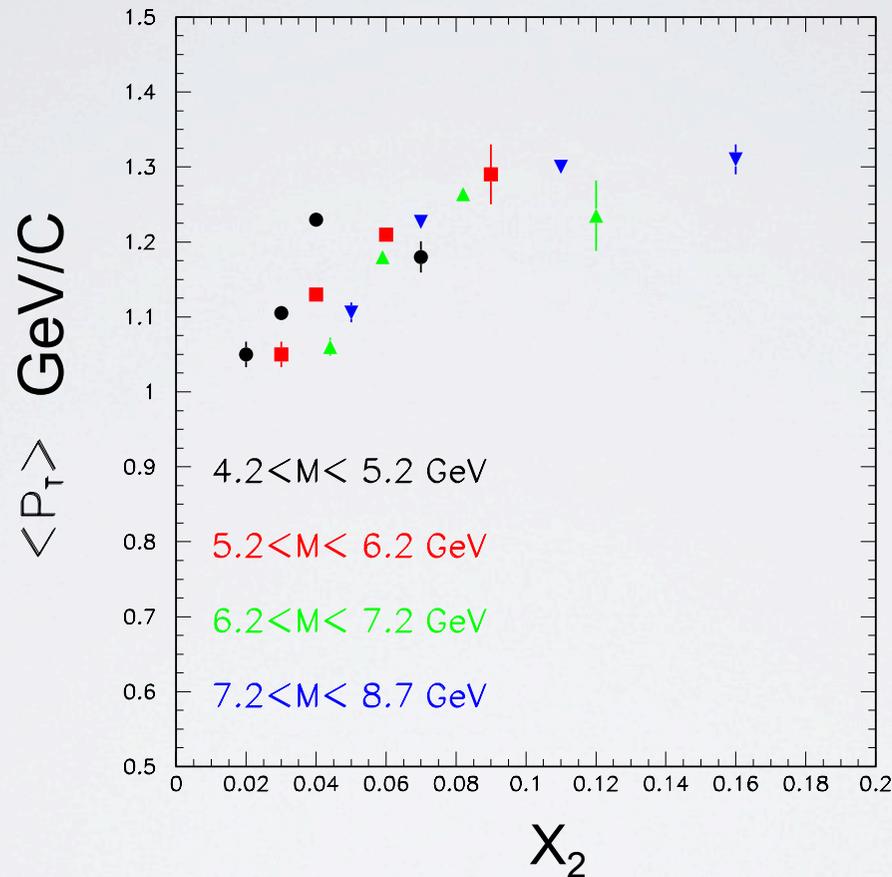
5.2 < M < 6.2 GeV

7.2 < M < 8.7 GeV



*E866/NuSea preliminary,  
talk by J.-C. Peng at DY@BNL workshop*

# New E866 data



Behavior opposite to BLNY fit

*E866/NuSea preliminary,  
talk by J.-C. Peng at DY@BNL workshop*

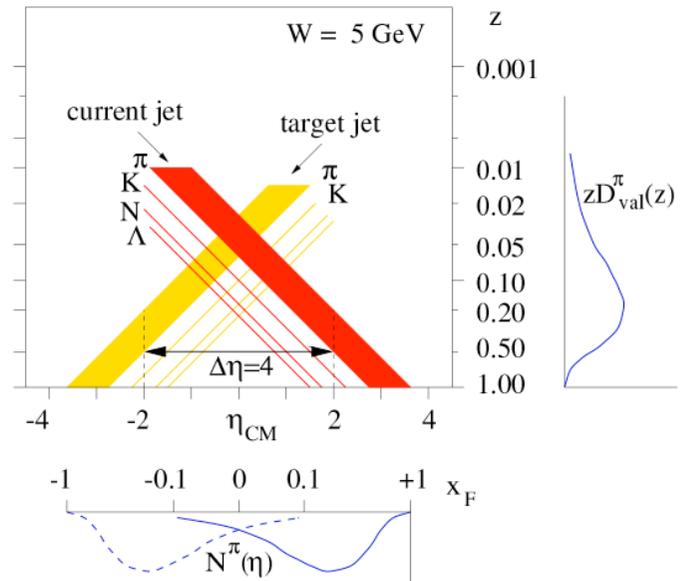
# multi-D SIDIS Multiplicities Coming $\rightarrow \langle k_{\perp} \rangle$ and $\langle p_{\perp} \rangle$

- How well do the **favored / disfavored** symmetries & **x-z factoriz<sup>n</sup>** hold?  
 ... assumed in  $\approx$  all FF global fits & PDF extractions  
 ... not exact at HERMES energies, acc to Lund MC

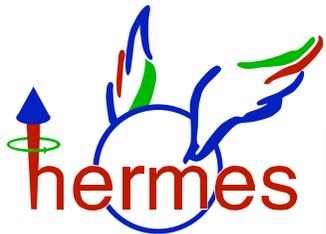
$$D_{\text{fav}} \equiv D_u^{\pi^+} = D_d^{\pi^-} = \dots$$

$$D_{\text{disfav}} \equiv D_u^{\pi^-} = D_d^{\pi^+} = \dots$$

- Are there **any** such FF symmetries for ***Kaons***?
- Does **intrinsic quark  $\langle k_T \rangle$**  vary with **x**?  
 ... with ***flavor***? (holy grail!)

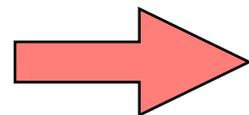


- Can the **Lund model** describe fragmentation at different **energies** / different **processes** (SIDIS vs e+e-) ***without retuning***?



*paper permanently  
in progress*

$d\sigma(x, z, p_T)$   
 $d\sigma(Q^2, z, p_T)$  for  $\pi^{\pm}, \pi^0, K^{\pm}, p, pbar$



***compare***

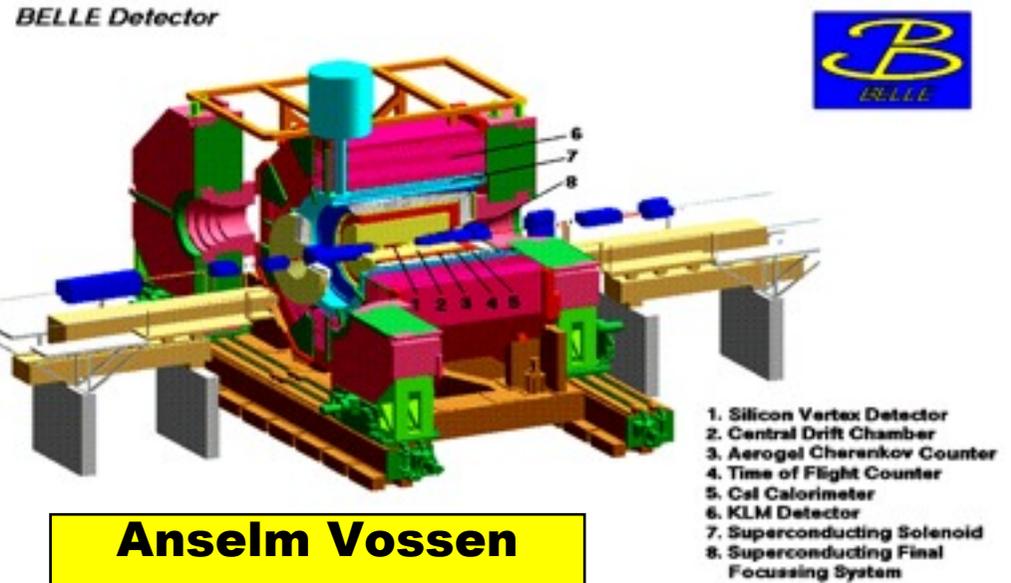
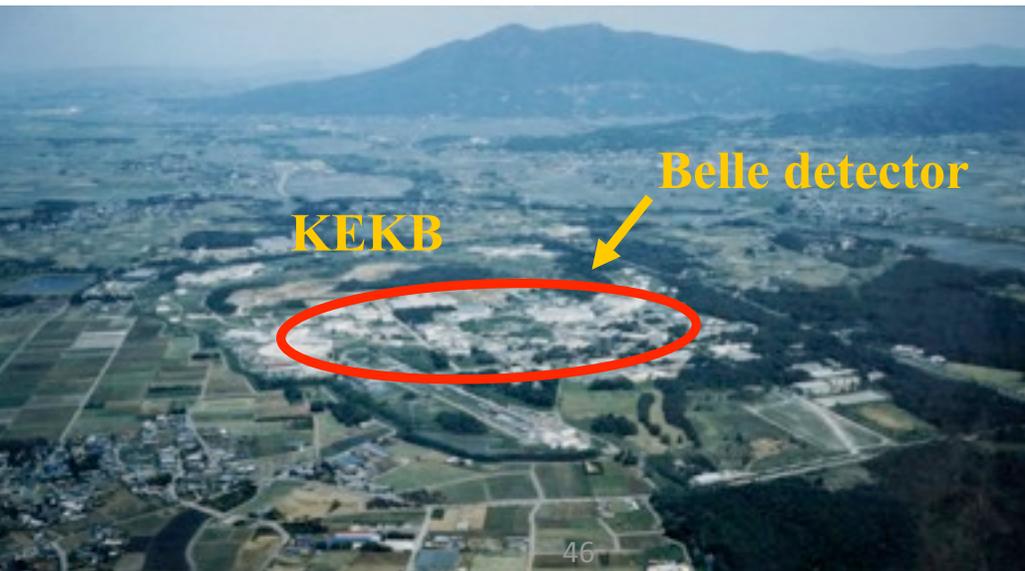
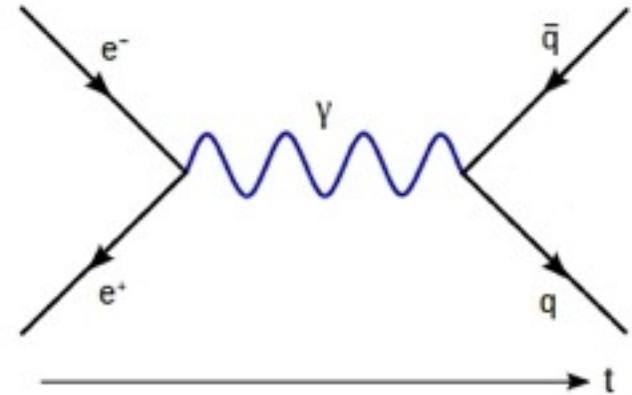
## COMPASS-II $\mu^{\pm}p$

- pure LH2 target
- higher energy
- RICH upgrade
- full 4D binning



# Measurements of Fragmentation Functions in $e^+e^-$ at Belle

- KEK-B: asymmetric  $e^+$  (3.5 GeV)  $e^-$  (8 GeV) collider:
  - $\sqrt{s} = 10.58$  GeV,  $e^+e^- \rightarrow Y(4S) \rightarrow B \bar{B}$
  - $\sqrt{s} = 10.52$  GeV,  $e^+e^- \rightarrow qq\bar{q}$  (u,d,s,c) 'continuum'
- ideal detector for high precision measurements:
  - tracking acceptance  $\theta$  [ $17^\circ$ ;  $150^\circ$ ]: Azimuthally symmetric
  - particle identification (PID):  $dE/dx$ , Cherenkov, ToF, EMcal, MuID
- Available data:
  - $\sim 1.8 \cdot 10^9$  events at 10.58 GeV,
  - $\sim 220 \cdot 10^6$  events at 10.52 GeV

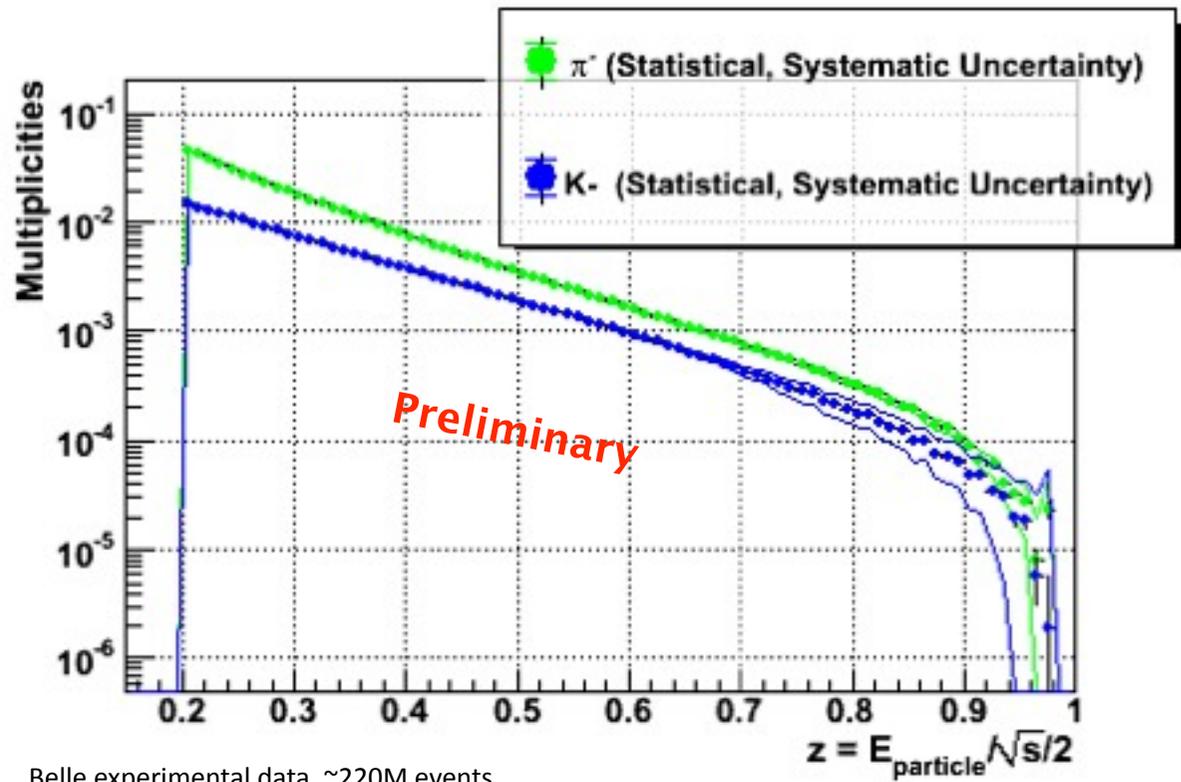




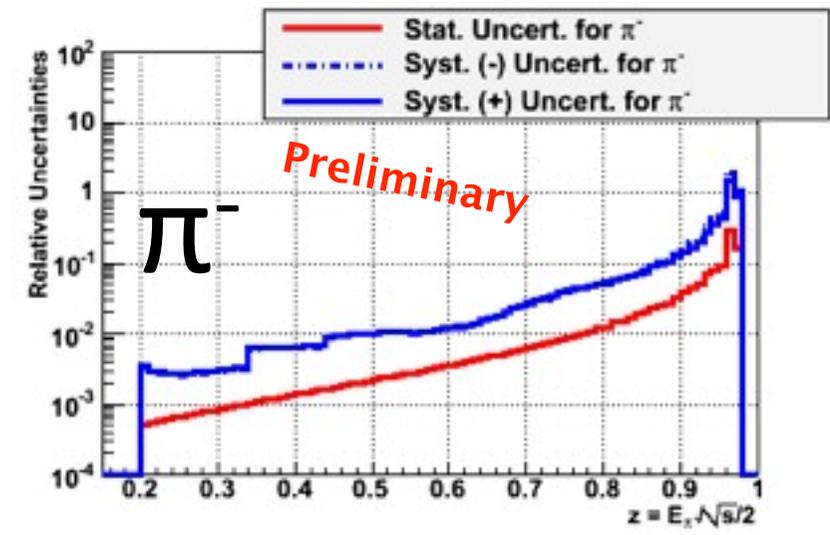
# Pion and Kaon Multiplicities

## Preliminary Results

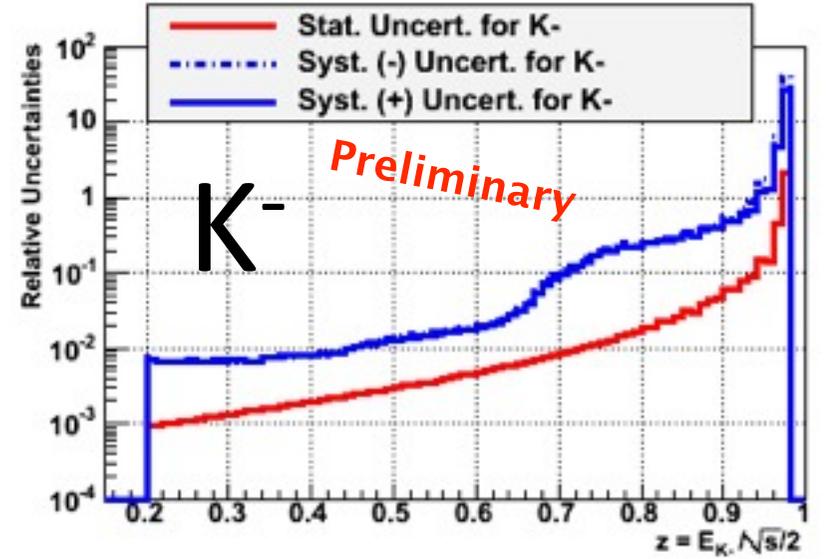
- Binning in  $z$ : width = 0.01; yields normalized to hadronic cross section
- Systematic uncertainties:  $z \sim 0.6$ : 1% (2%) for  $\pi$  (K);  
 $z \sim 0.9$ : 14% (50%) for  $\pi$  (K)



Belle experimental data, ~220M events



Additional normalization uncertainty of 1.4% not shown.



## Can we disentangle $\langle k^2_{\perp} \rangle$ , $\langle p^2_{\perp} \rangle$ ?

- Global fit coming for new unpol **multiD-binned** xsecs from
  - SIDIS: HERMES, COMPASS, JLab
  - $e^+e^-$ : BELLE (impossibly large statistics)
  - DY: COMPASS-II, SeaQuest
- Can relations like  $\mathbf{p}^h_{\perp} = \mathbf{k}_{\perp} - \mathbf{z} \mathbf{p}_{\perp}$  help? Too sloppy?
- $\langle k^2_{\perp} \rangle$  and  $\langle p^2_{\perp} \rangle$  likely depend on **flavor, PDF/FF, scale ...**

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**Is this Doable? At least for pions / light-quarks?**



*work together,  
bears*



Conventions:

Trento II

# Trans. mom. notations

Amsterdam	Torino	INT	Description
$p$	$k$	$k$	momentum of parton in distribution function
$p_T$	$k_{\perp}$	$k_{\perp}$	parton transverse momentum in distribution function
$k$	$p$	$p$	momentum of fragmenting parton
$k_T$		$p_{\perp}$	trans. momentum of fragmenting parton w.r.t. final hadron
$K_T$	$p_{\perp}$	$P_{\perp}$	trans. momentum of final hadron w.r.t. fragmenting parton
$P_{h\perp}$	$P_T$	$P_{hT}$	transverse momentum of final hadron w.r.t. virtual photon
$q_T$		$q_T$	transverse momentum of final photon w.r.t. hadron-hadron axis

## Ingredients of symbology

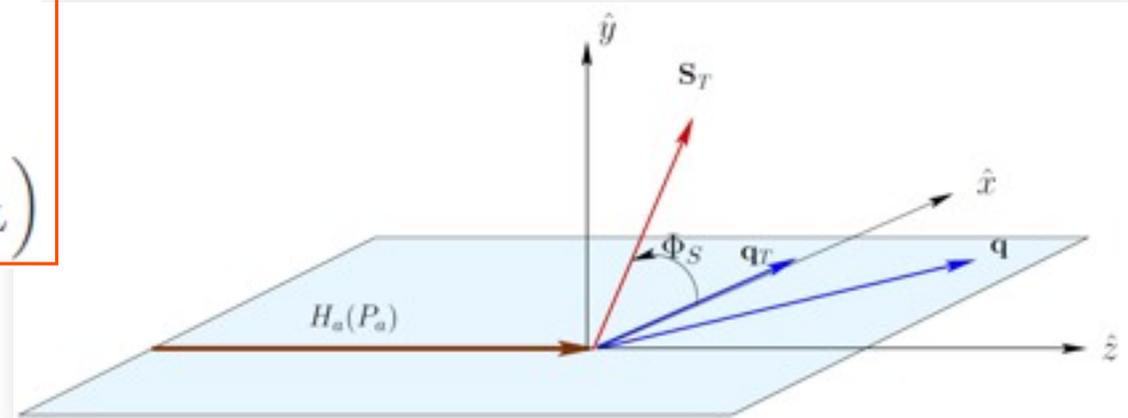
- $k$  vs  $p$  parton from distribution vs fragmentation function
- $\perp$  vs  $T$  internal vs measurable transverse component
- $a$  vs  $A$  quark vs hadron variable

**Source: Alessandro Bacchetta,  
Marco Radici**

# Angular variables

figures: Aram Kotzinian

$$\begin{aligned}
 P_{a,TF}^\mu &= (E, 0, 0, P_{a,TF}^3), \\
 P_{b,TF}^\mu &= (M_b, 0, 0, 0), \\
 q_{TF}^\mu &= (q_{0,TF}, q_T, 0, q_{L,TF}), \\
 S_{TF}^\mu &= (0, |\vec{S}_T| \cos \phi_S, |\vec{S}_T| \sin \phi_S, S_L)
 \end{aligned}$$



Target rest frame (TF)

1. Define  $\phi_{\text{spin}}$  as shown

$$\begin{aligned}
 l_{CS}^\mu &= \frac{q}{2} (1, \sin \theta \cos \phi, \sin \theta \sin \phi, \cos \theta), \\
 l_{CS}^{\prime\mu} &= \frac{q}{2} (1, -\sin \theta \cos \phi, -\sin \theta \sin \phi, -\cos \theta)
 \end{aligned}$$

Collins-Soper frame (CS)

1. Boost along beam until  $q_L = 0$
2. Boost along  $q$  until  $q_T = 0$
3. lepton  $\theta, \phi$  defined with respect to lepton ( $\mu^-, e^-$ ) not anti-lepton

