

Quadrupole Moments of Wobbling Excitations in ^{163}Lu

Andreas G3rgen et al.

talk by

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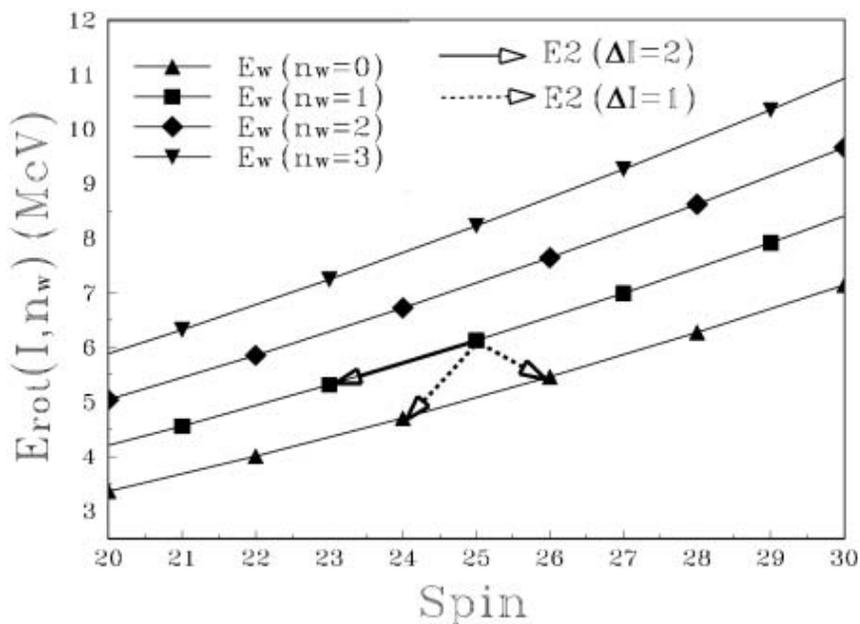
Nuclei at the Limits, Argonne, July 26-30 2004

Wobbling:

Bohr and Mottelson:

triaxiality: $J_x > J_y > J_z$

high spin: $I \approx I_x \gg 1$

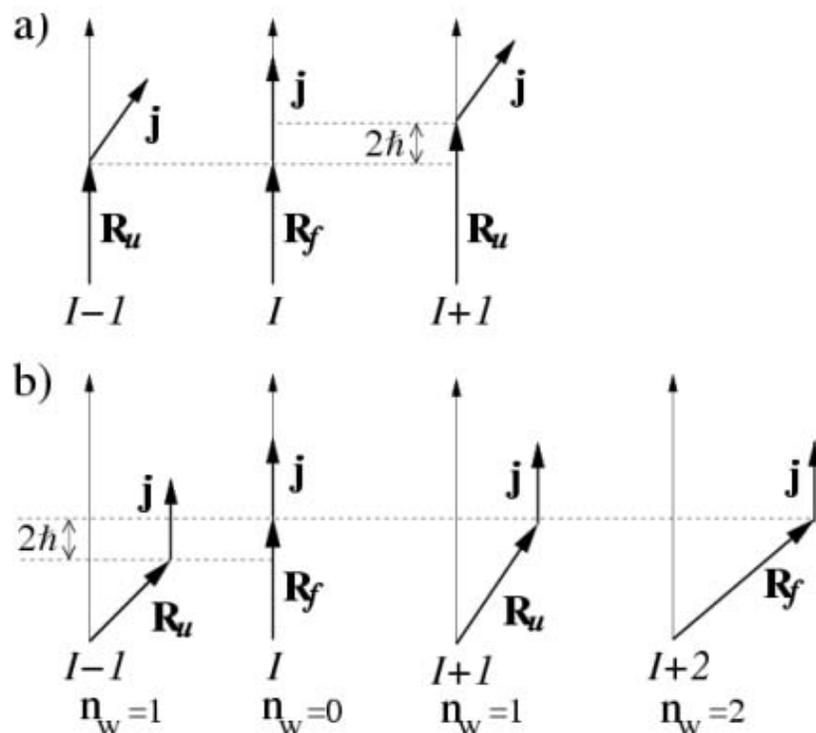


$$E_R(I, n_w) = \frac{I(I+1)}{2J_x} + \hbar\omega_w(n_w + 1/2)$$

$$\hbar\omega_w = \frac{I}{J_x} \sqrt{\frac{(J_x - J_y)(J_x - J_z)}{J_y J_z}}$$

Wobbling with high-j aligned particles:

“cranking scenario”



“wobbling scenario”

I. Hamamoto, Phys. Rev. C **65**, 044305 (2002)

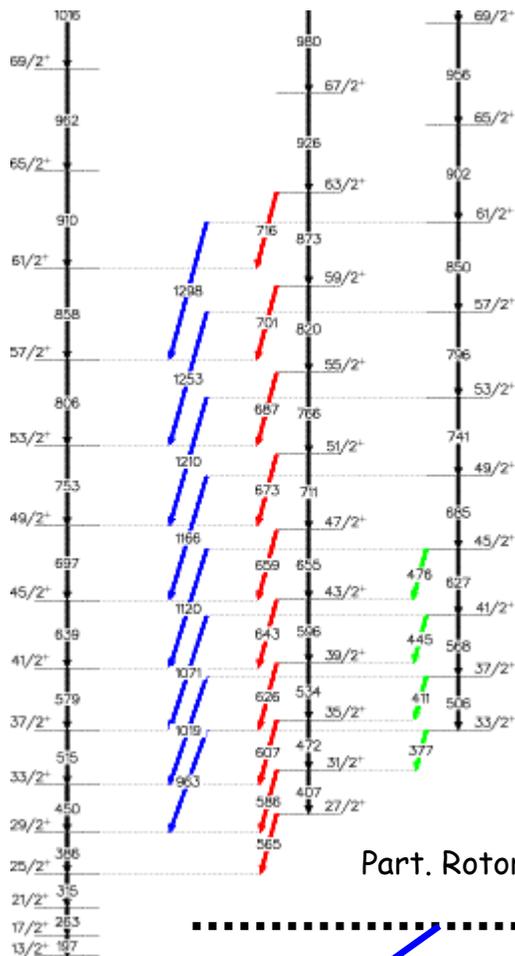
I. Hamamoto and G.B. Hagemann, Phys. Rev. C **67**, 014319 (2003)



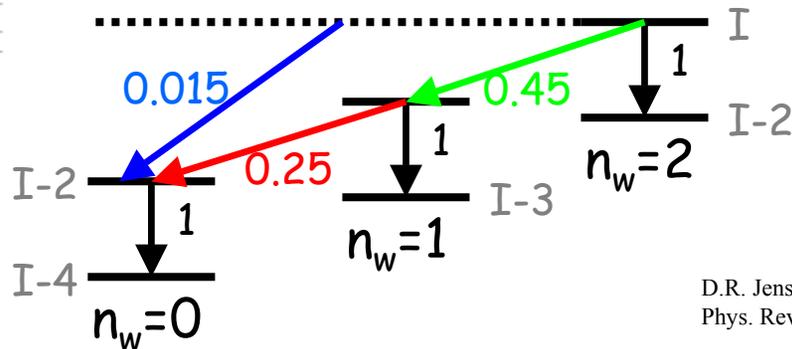
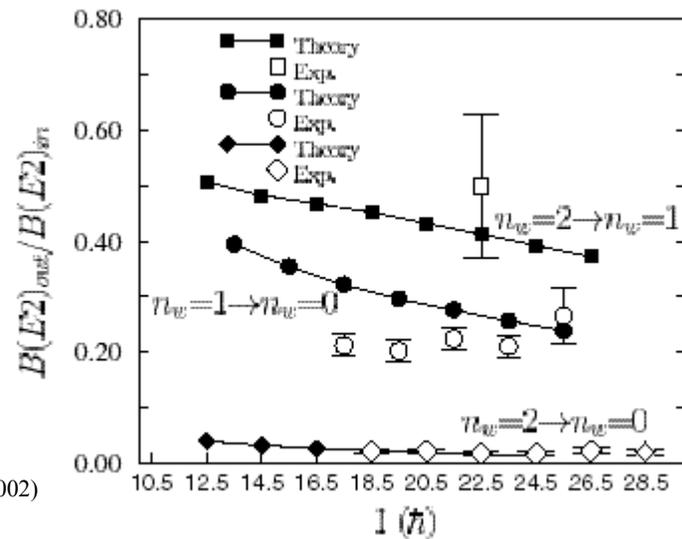
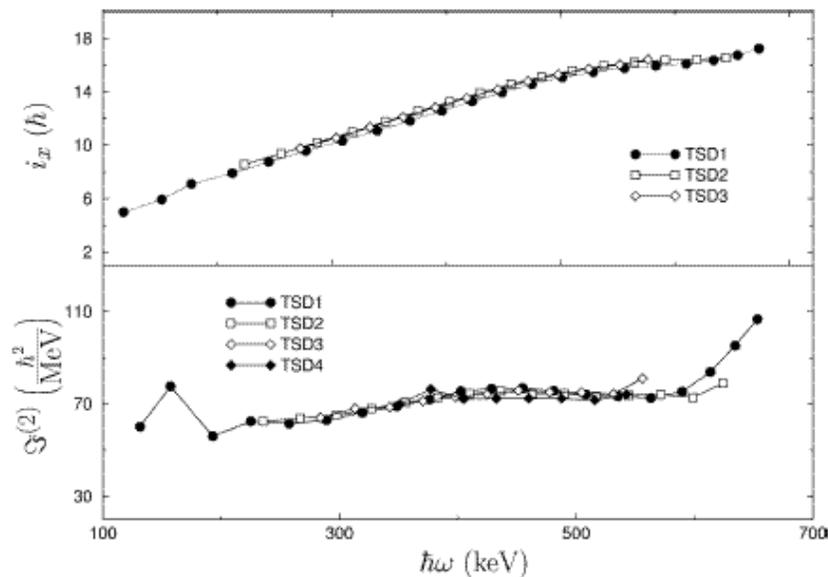
TSD1

TSD2

TSD3



$$B(E2; n_w, I \rightarrow n_w - 1, I - 1) \propto n_w / I$$

Part. Rotor calc., $B(E2)$'sEvidence for the wobbling mode in ^{163}Lu 

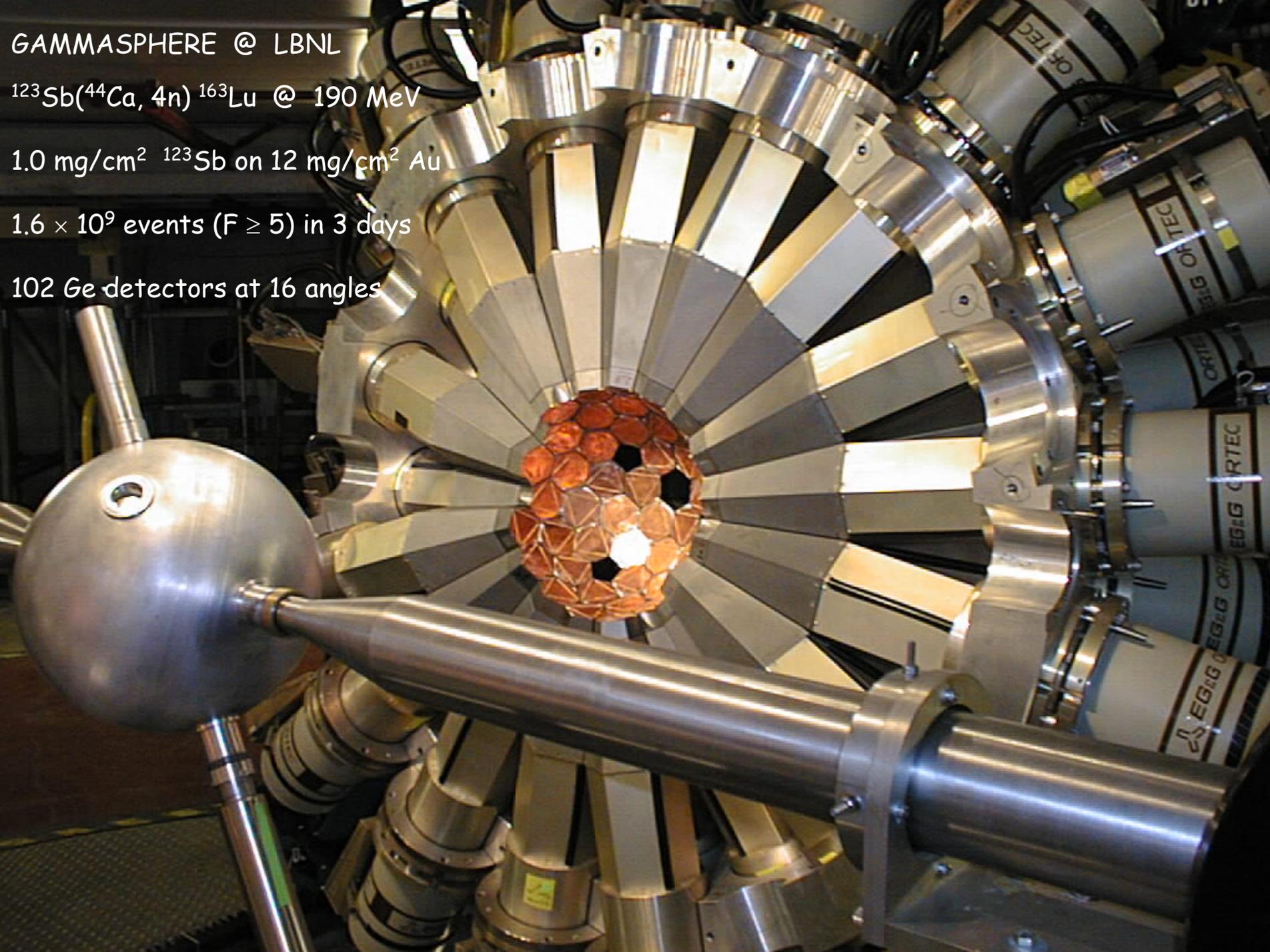
GAMMASPHERE @ LBNL

$^{123}\text{Sb}(^{44}\text{Ca}, 4n)^{163}\text{Lu}$ @ 190 MeV

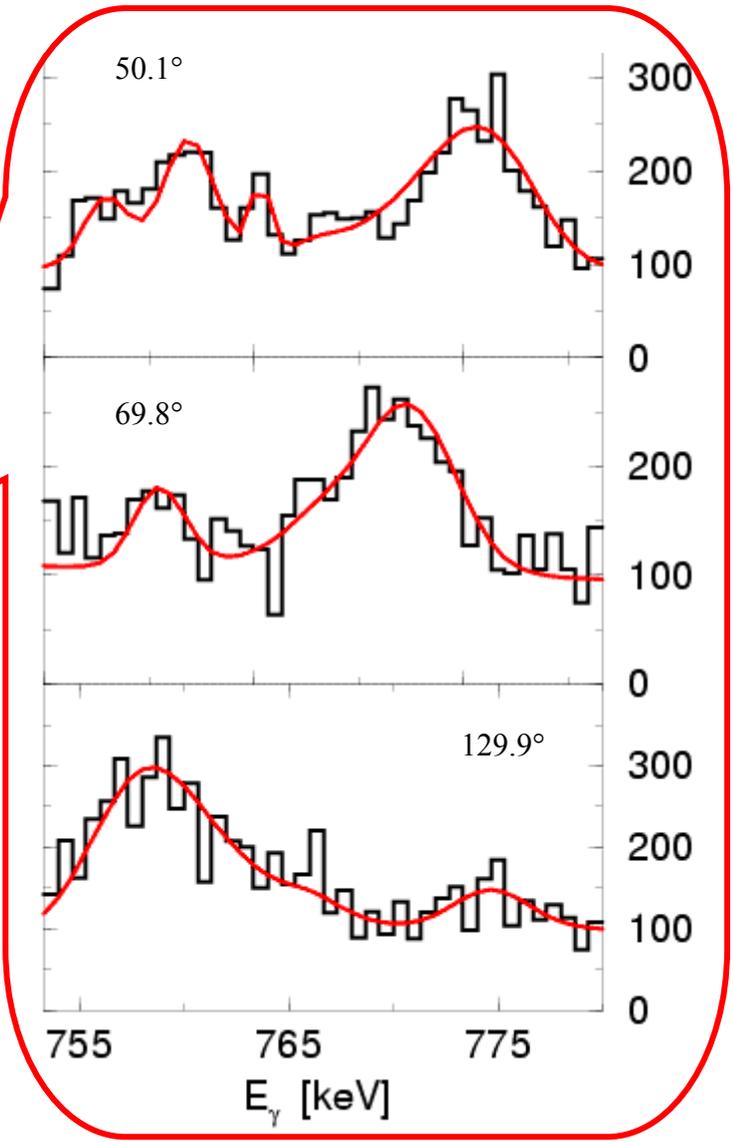
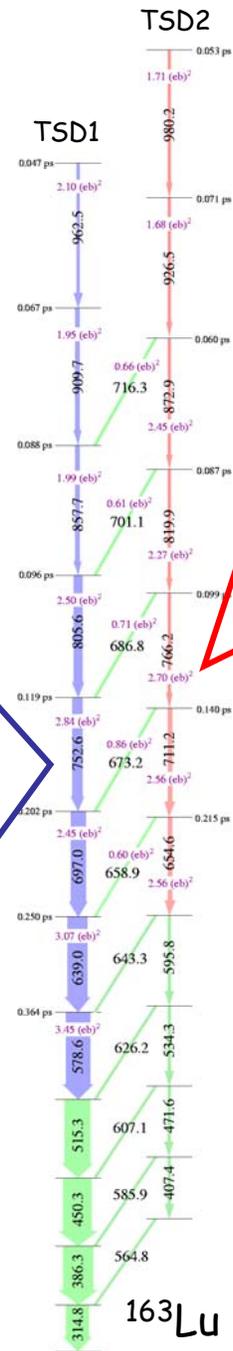
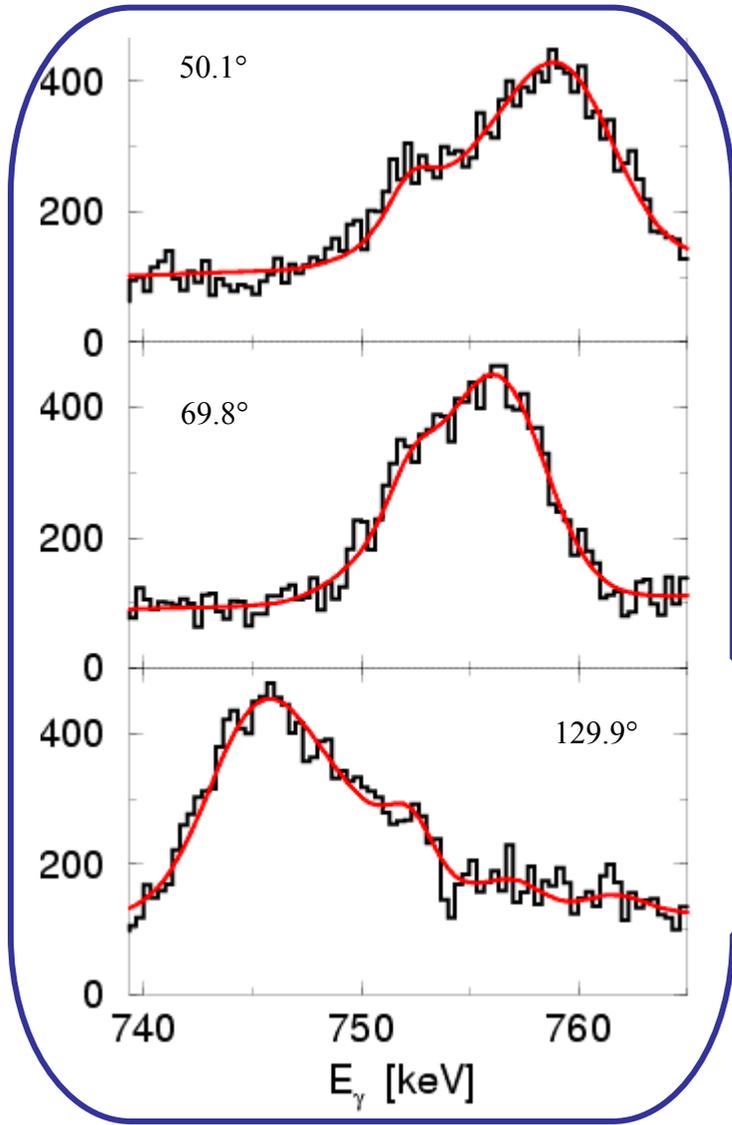
1.0 mg/cm² ^{123}Sb on 12 mg/cm² Au

1.6×10^9 events ($F \geq 5$) in 3 days

102 Ge detectors at 16 angles

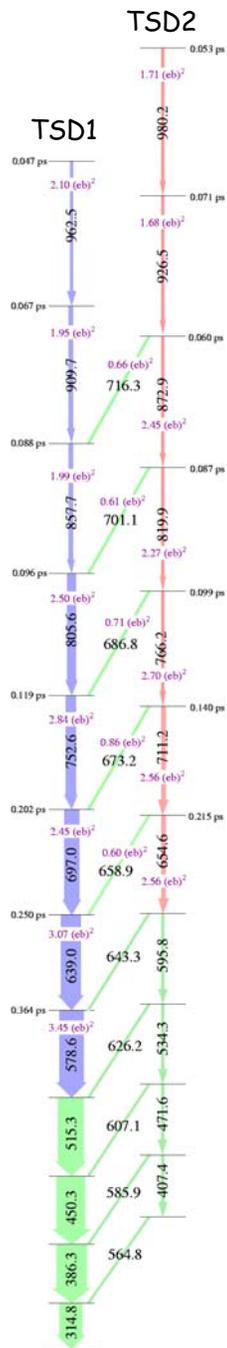


DSAM lineshape analysis



Lifetimes measured for

- 8 states in TSD1
- 7 states in TSD2



We know:

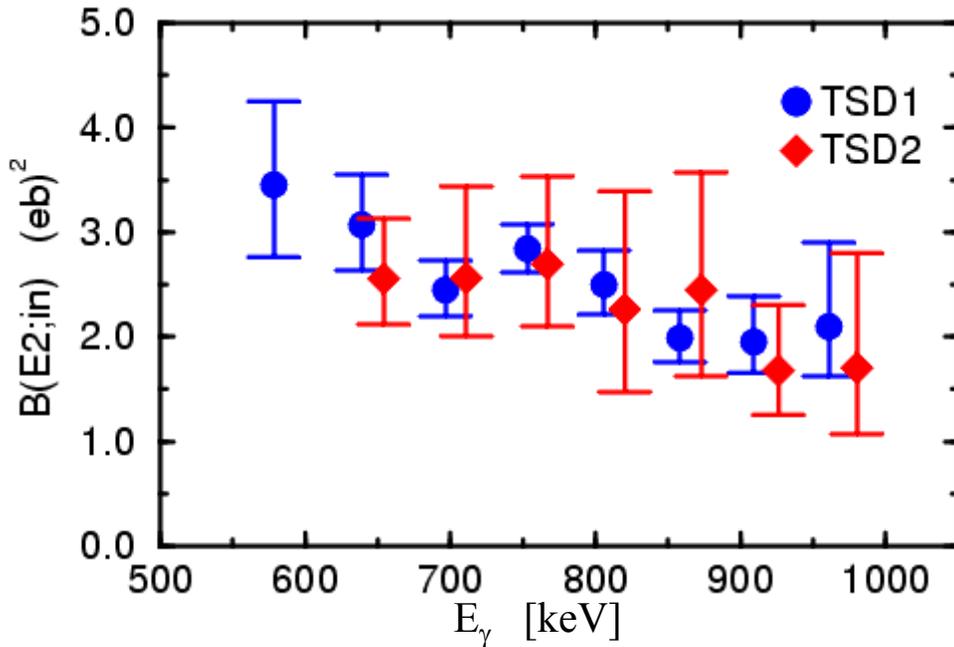
This measurement:

- total transition probabilities $\lambda_{\text{tot}}=1/\tau$
- for TSD1: $B(E2; \text{in}) = B(E2; n_w=0, I \rightarrow n_w=0, I-2)$

Thin target experiments at Euroball:

- branching ratios $\lambda_{\text{out}}/\lambda_{\text{in}}$
 - ⇒ for TSD2: $B(E2; \text{in}) = B(E2; n_w=1, I \rightarrow n_w=1, I-2)$
- mixing ratio δ (from DCO/lin. polarization)
 - $\delta = -3.10^{+0.36}_{-0.44} \Rightarrow 90\% \text{ E2 and } 10\% \text{ M1}$
 - ⇒ $B(E2; \text{out}) = B(E2; n_w=1, I \rightarrow n_w=0, I-1)$
 - ⇒ $B(M1; \text{out}) = B(M1; n_w=1, I \rightarrow n_w=0, I-1)$

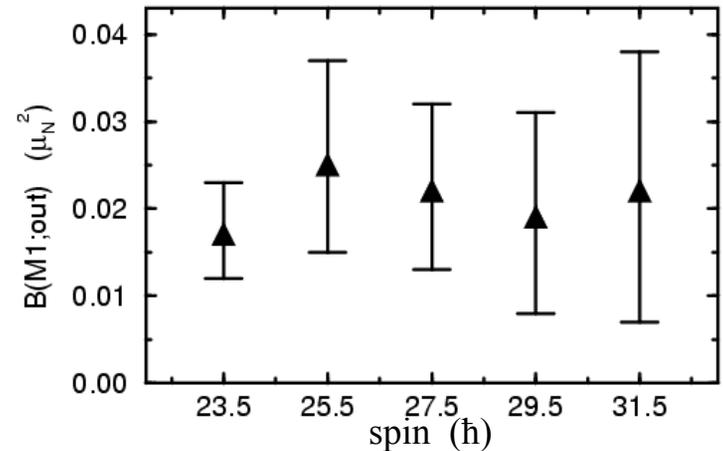
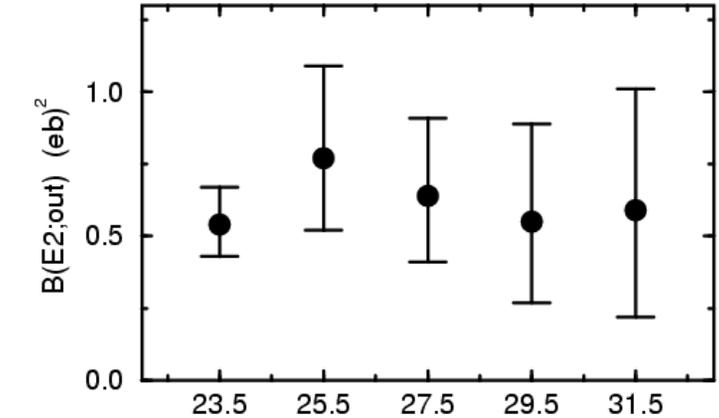
transition strengths



in-band transitions:

- $B(E2)$'s very similar for TSD1 and TSD2
- both show a decrease

What can we learn ?



inter-band transitions:

- $B(E2)$'s and $B(M1)$'s are
~constant (with large errors)

How to extract a quadrupole moment from a B(E2)

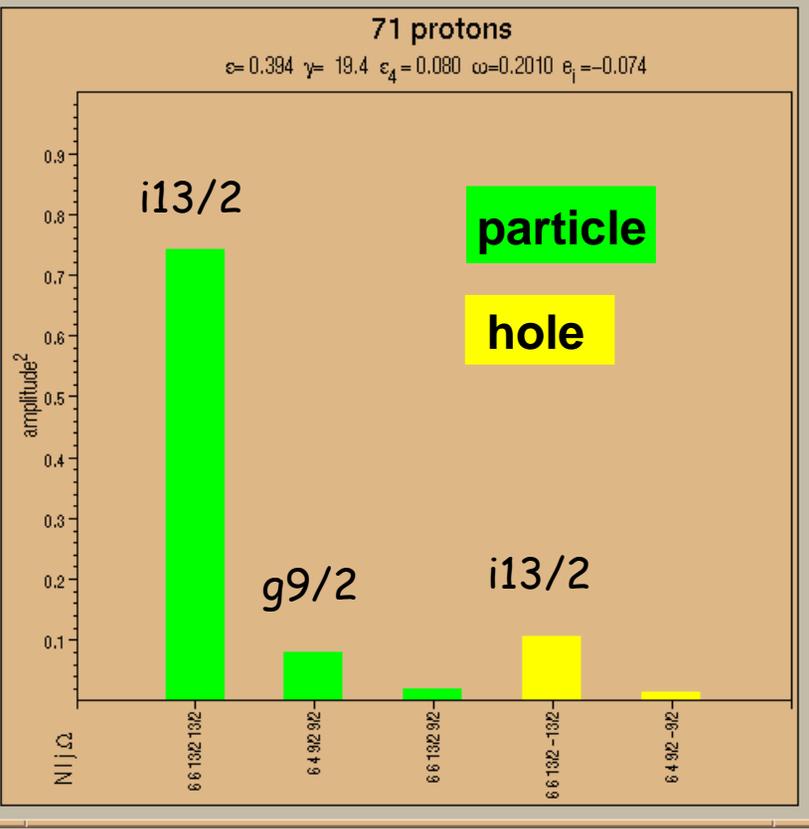
axial symmetric nuclei: $B(E2) = \frac{5}{16\pi} (eQ)^2 \langle I_i K 2 0 | I_f K \rangle^2$

K is not a good quantum number in a triaxial nucleus

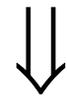
we have to use a distribution of different K values:

$$B(E2) = \frac{5}{16\pi} (eQ)^2 \sum_K C_K^2 \langle I_i K 2 0 | I_f K \rangle^2$$

How do we get the coefficients C_K^2 ?



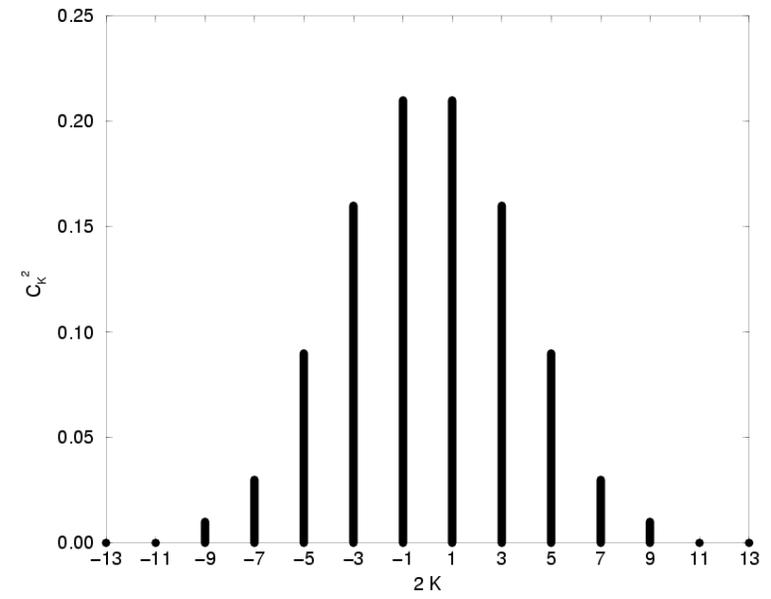
Ultimate Cranker:
 "pure" $i_{13/2}$ configuration
 aligned with rotational axis

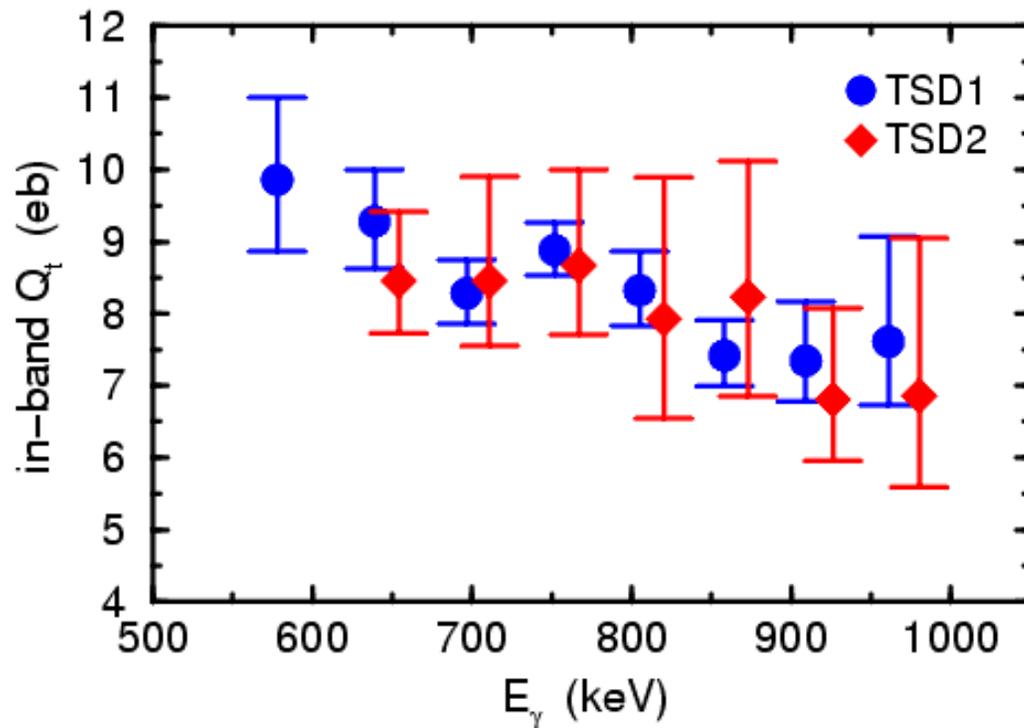
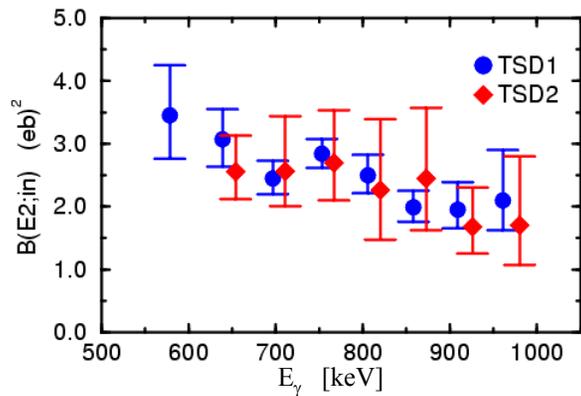


geometric problem

$\Rightarrow C_K^2$ from Wigner's D-functions
 $D_{jK}^j(0, \pi/2, 0)$ with $j = 13/2$

$$B(E2) = \frac{5}{16\pi} (eQ)^2 \sum_K C_K^2 \langle I_i K 2 0 | I_f K \rangle^2$$





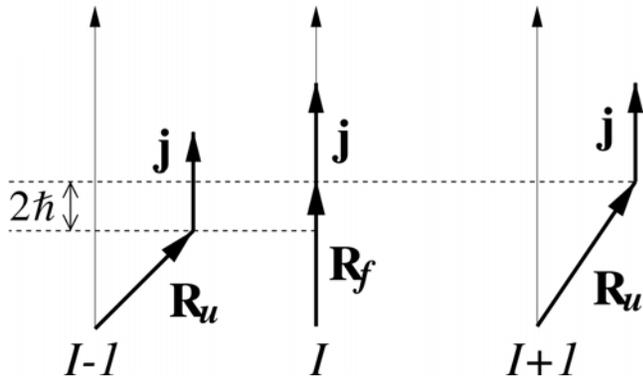
almost identical quadrupole moments

⇒ similar intrinsic structure

we need to understand the decrease

Should the quadrupole moments for TSD1 and TSD2 be the same?

different vector coupling



$$\begin{array}{c} \text{TSD1} \\ \text{K} = \text{I} \end{array}$$

$$\begin{array}{c} \text{TSD2} \\ \text{K} = \text{I}-1 \end{array}$$

(with rotational axis as quantization axis)

relevant Clebsch-Gordan: $\langle \text{I} \text{ K} 2 -2 \mid \text{I}-2 \text{ K}-2 \rangle$

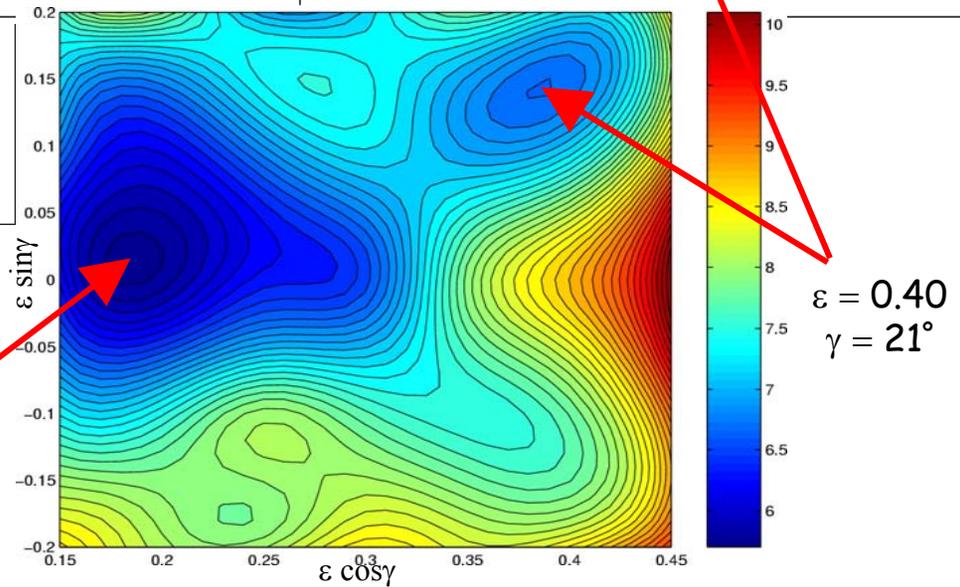
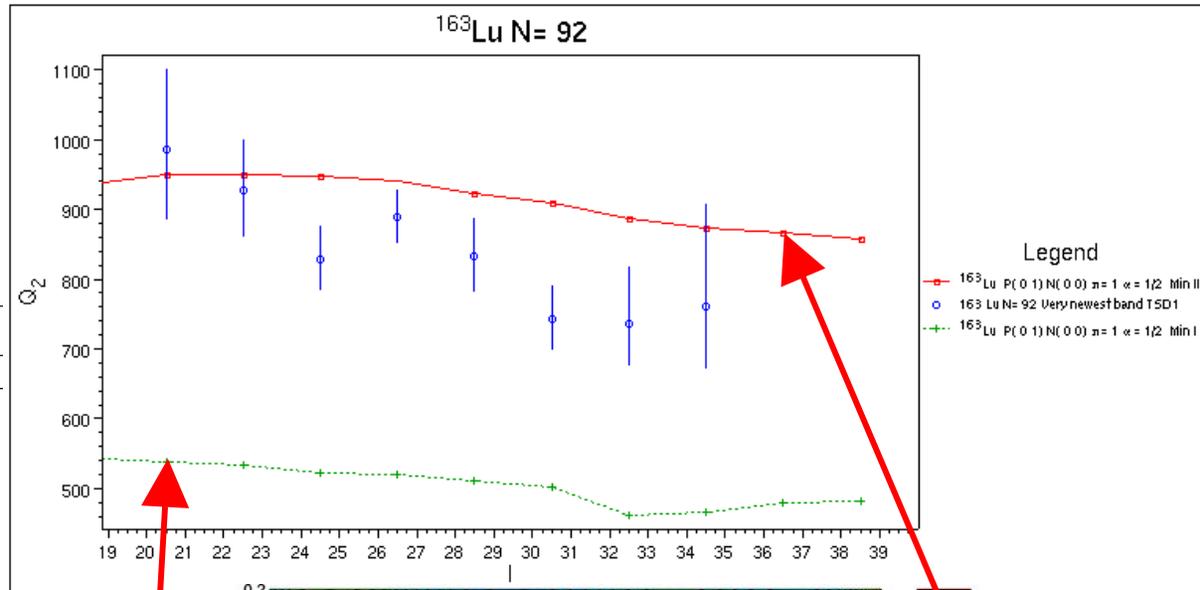
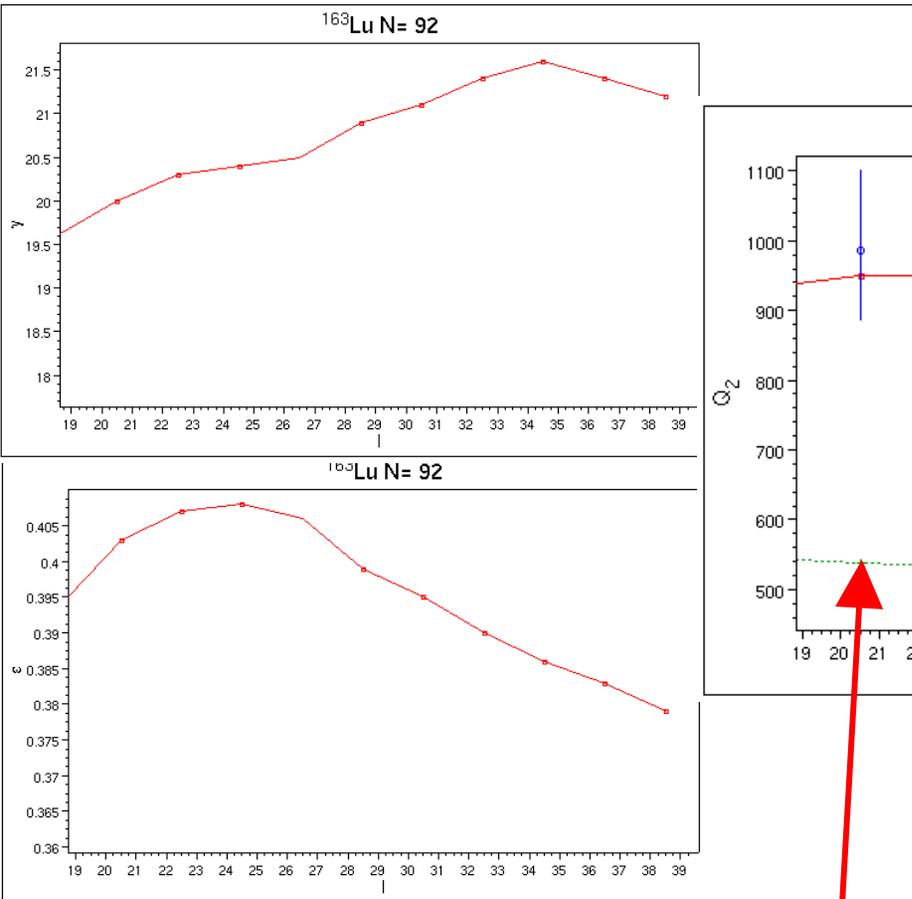
examples:

$$\langle 55/2 \ 55/2 \ 2 \ -2 \mid 51/2 \ 51/2 \rangle = 0.964$$

$$\langle 55/2 \ 53/2 \ 2 \ -2 \mid 51/2 \ 49/2 \rangle = 0.928$$

⇒ ~ 4 % difference between 0-phonon and 1-phonon bands

Cranking calculations for TDS1 (ultimate cranker)

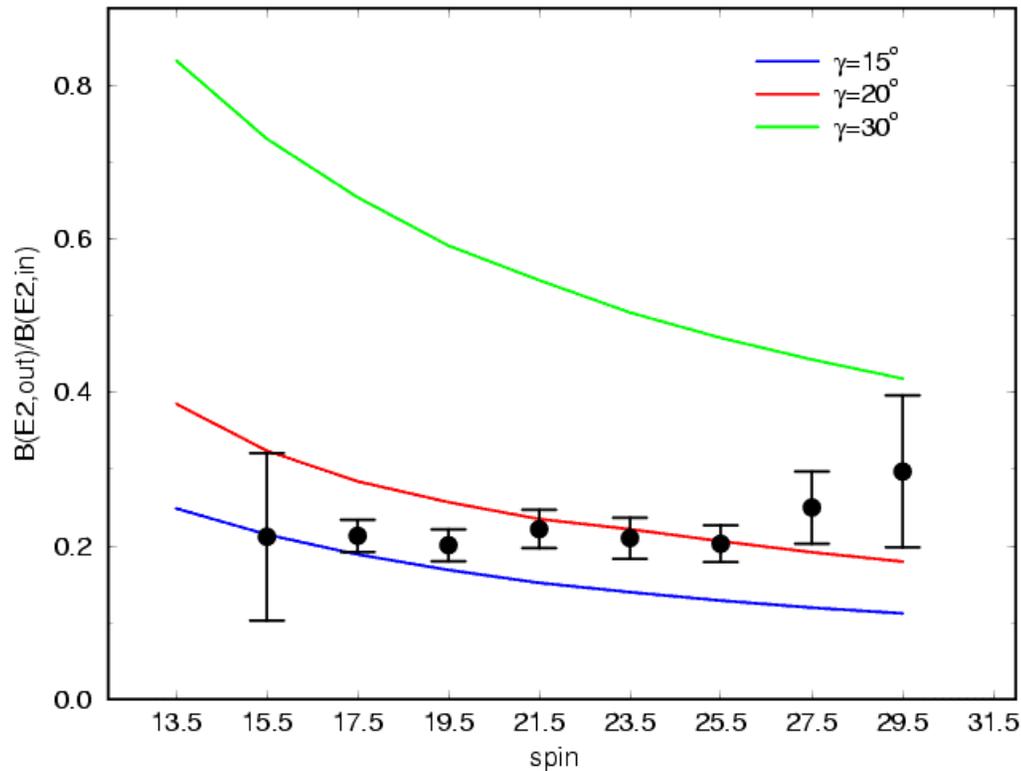


R. Bengtsson, private communication

The ratio $B(E2;out)/B(E2;in)$ is very sensitive to γ ,
but independent of ε

$$B(E2; n_w, I \rightarrow n_w - 1, I - 1) \propto \frac{n_w}{I} \sin^2(\gamma + 30^\circ)$$

$$B(E2; n_w, I \rightarrow n_w, I - 2) \propto \cos^2(\gamma + 30^\circ)$$



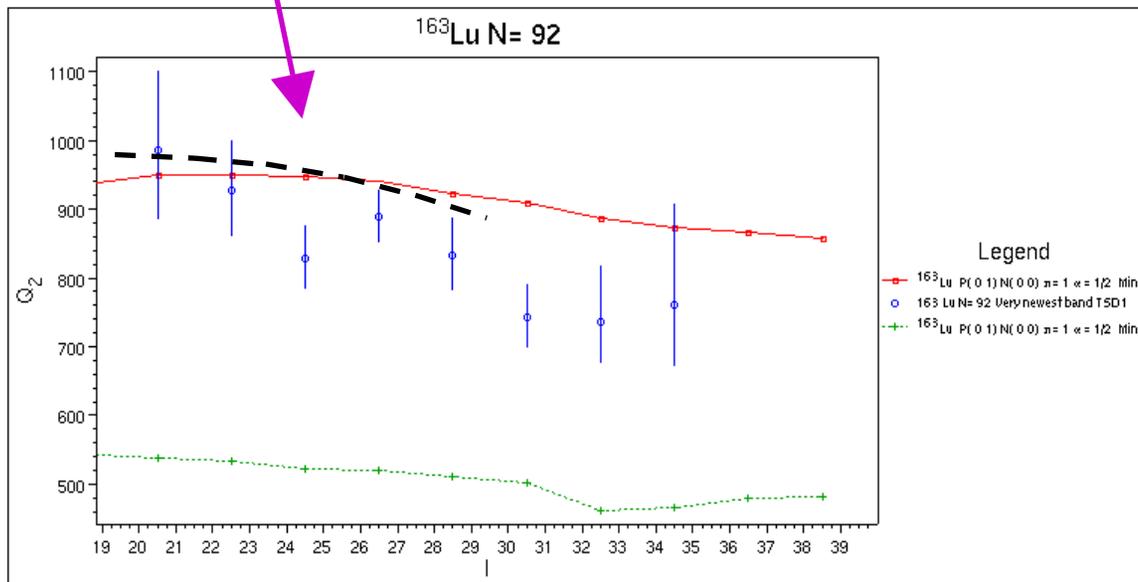
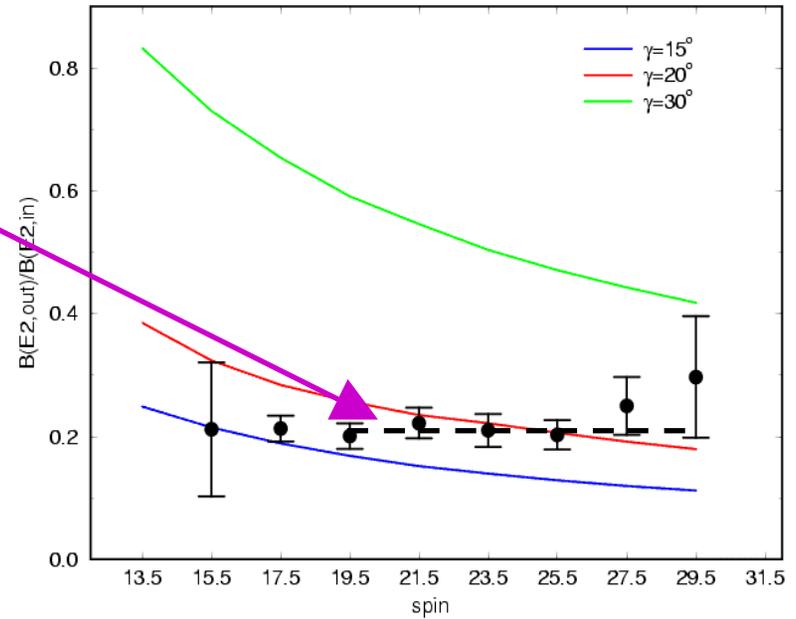
Particle-rotor calculation from:

I. Hamamoto and G.B. Hagemann,
Phys. Rev. C **67**, 014319 (2003)

Constant $B(E2)$ ratio
can be explained by
increase in γ from
 $\sim 16^\circ \dots 22^\circ$

Constant B(E2) ratio can be explained by increase in γ from $\sim 16^\circ \dots 22^\circ$

Explains stronger decrease of Q at the same time.



Summary:

- strong evidence for wobbling phonon excitations in odd-mass Lu isotopes based on characteristic E2 inter-band transitions
 - lifetime measurements in ^{163}Lu find very similar Q_{\dagger} 's for TSD1 and TSD2
⇒ strong indication for similar intrinsic structure of the wobbling bands
 - in-band Q_{\dagger} 's show decrease with angular momentum, trend is reproduced by cranking calculations: increase in γ and decrease in ε
 - constant B(E2)'s of the inter-band transitions can be understood by increase in γ , thus restoring the 1/I dependence inherent in the wobbling mode
- ⇒ new results support the wobbling picture and give a handle on triaxiality for the first time

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