

# Congratulations!

## John

On Fifty Remarkable Years of  
Research at Argonne





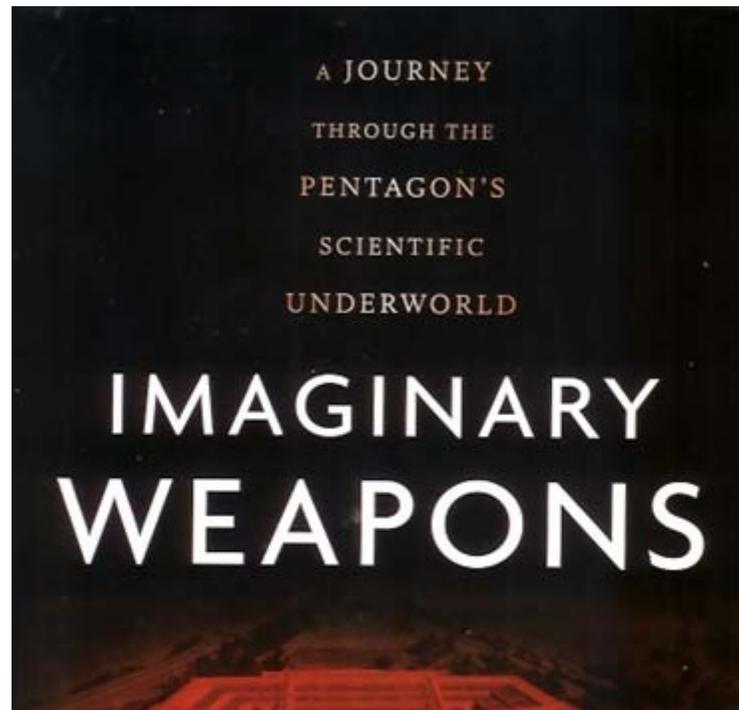
me

Me again  
John P. Schiffer

art

BUT WAIT,  
THESE  
ARE THE  
BEST TO FILL

CAN STRAIGHTEN



their students. Among colleagues, Schiffer had a reputation as a world-class scientist, albeit a tad blunt, with a habit of calling people fools to their face. But Schiffer was not the

Me

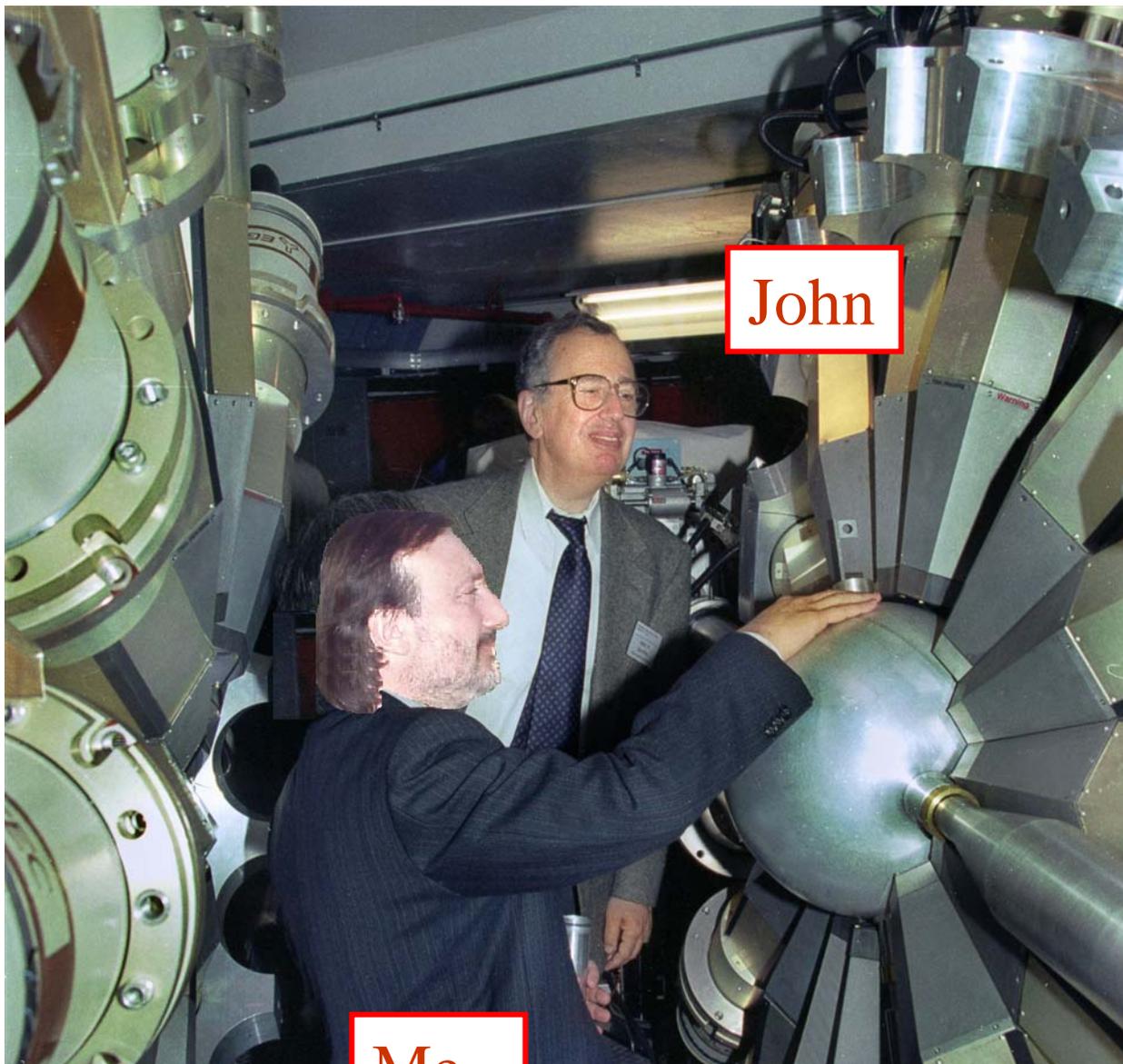
Me

John



GORDON RESEARCH CONFERENCE  
Colby Sawyer College (1)  
NUCLEAR CHEMISTRY

# Checking a setup in GAMMASPHERE !

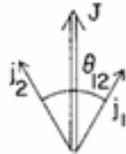
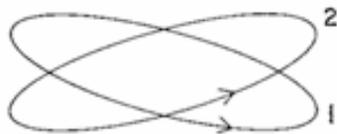


John

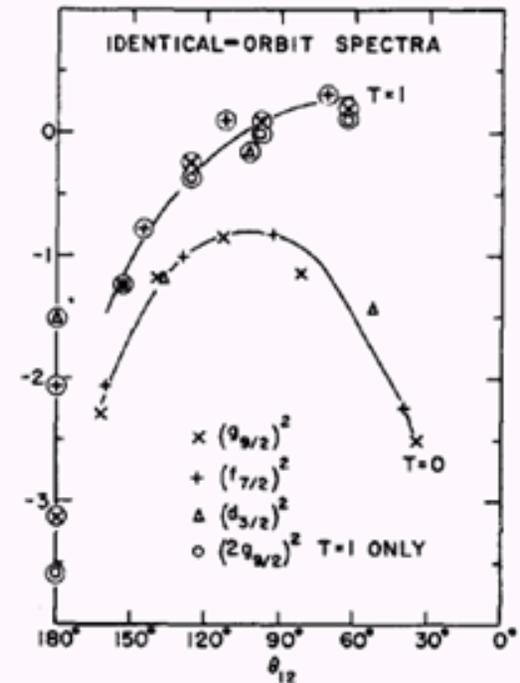
Lee

# The effective interaction between nucleons deduced from nuclear spectra

John P. Schiffer, William W. True



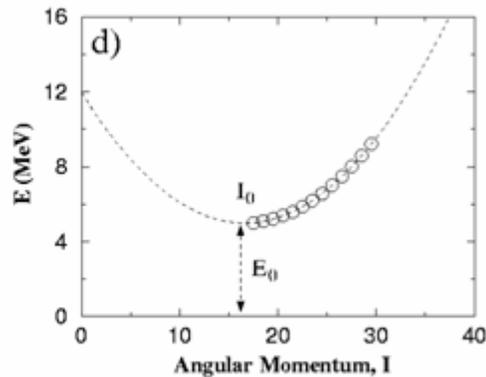
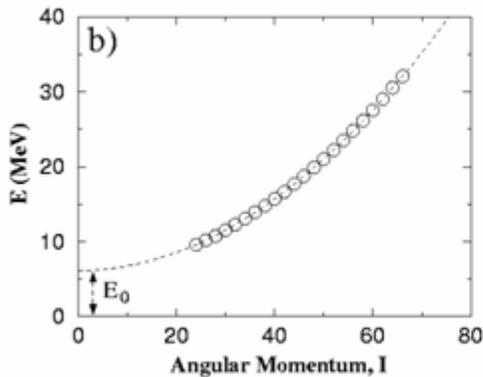
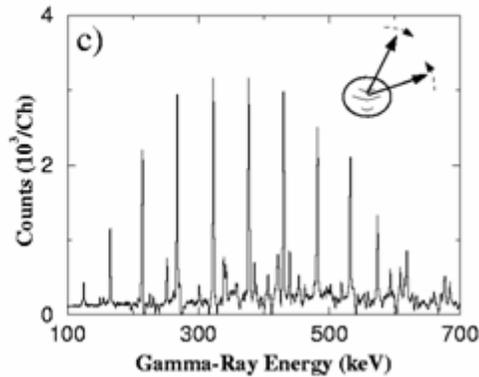
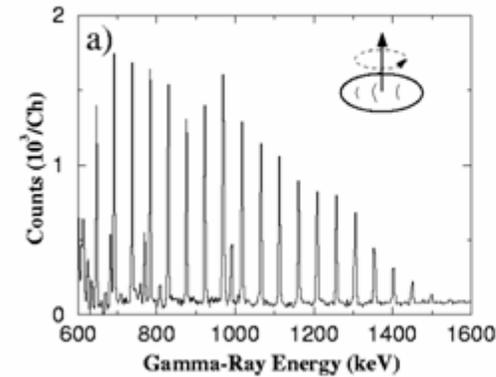
$$\theta_{12} = \cos^{-1} \frac{j_1(j_1+1) + j_2(j_2+1) - J(J+1)}{[2j_1j_2(j_1+1)(j_2+1)]^{1/2}}$$



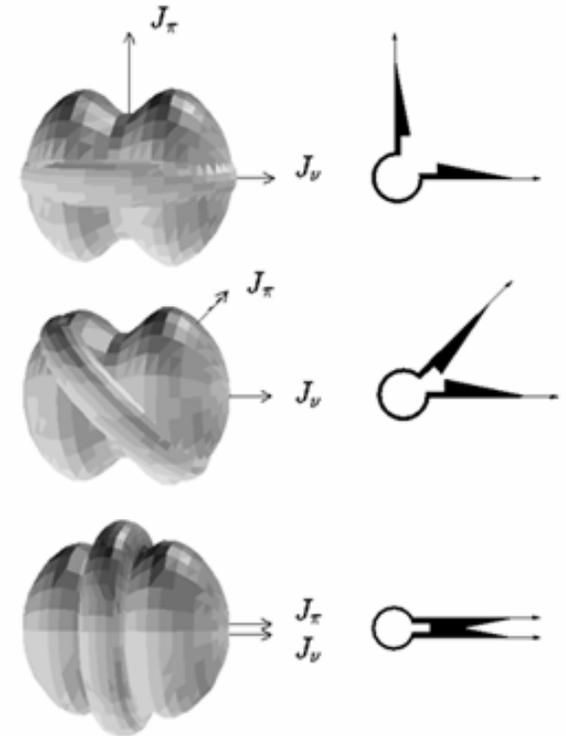
# Rotational-like (Shears) bands in near-spherical nuclei

## SD Band in $^{152}\text{Dy}$

## M1 Band in $^{199}\text{Pb}$



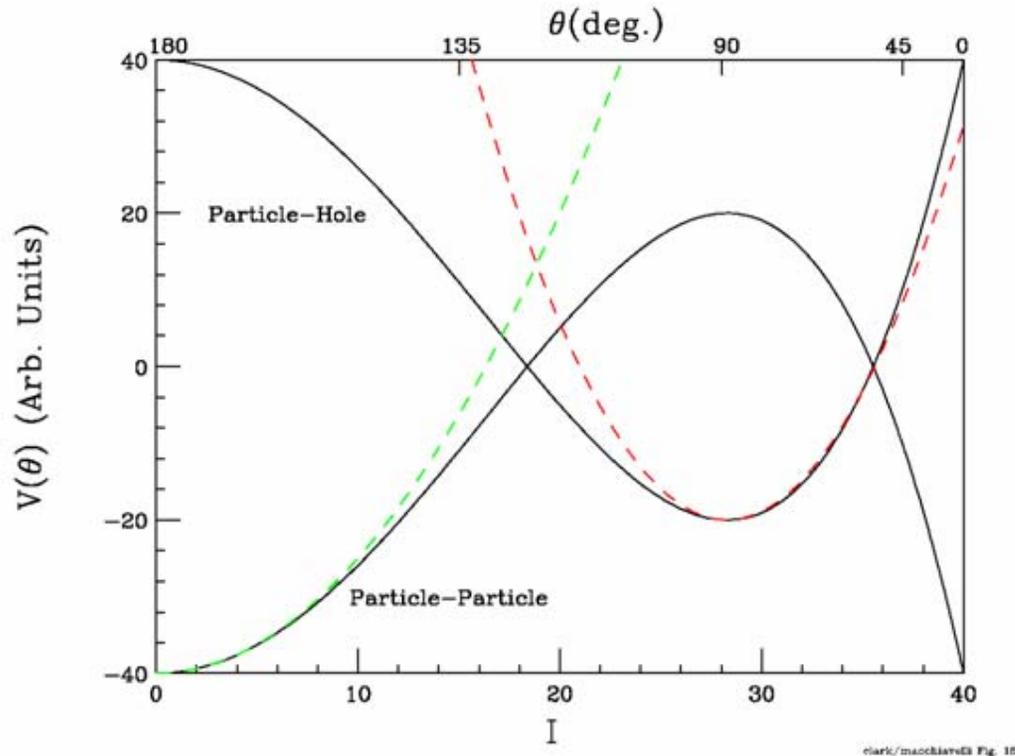
darkmacchiwell Fig. 1



Angular momentum generated by the recoupling of the protons and neutrons components (shears)

$B(M1)$ , proportional to  $\mu_{\perp}^2 \Rightarrow$  should decrease with  $I$

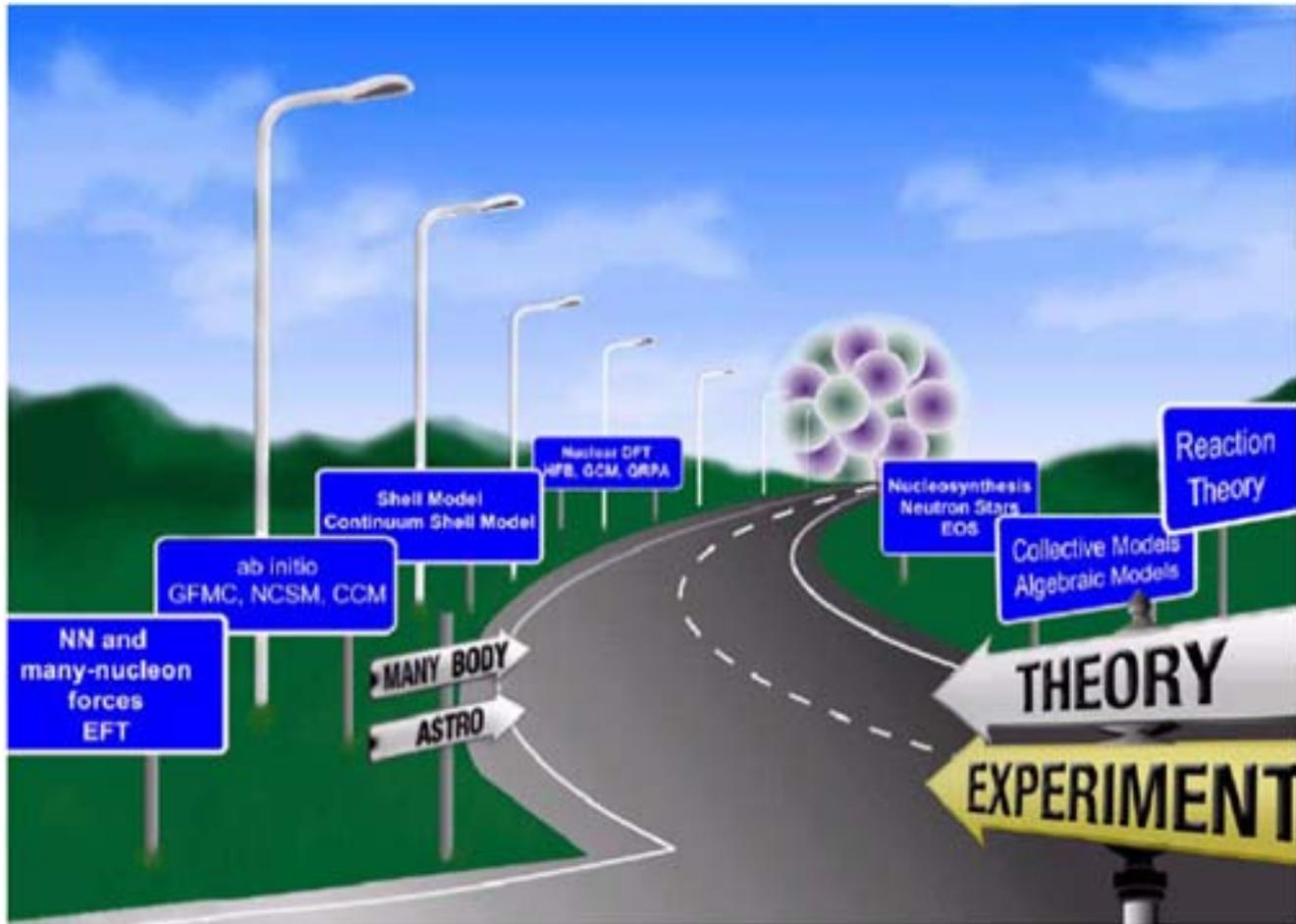
## Effective $P_2$ Interaction ?



**Effective  $P_2$  interaction can give rise to the rotational-like spectrum**  
**Particle-particle case will correspond to Anti-magnetic rotation**

**$\Rightarrow$  Although the main ingredients and the general properties of the shears bands are understood, the regularity of these bands is still an important theoretical question from a microscopic point of view.**

Inspired by the RIA Theory Blue Book and the recent RIA brochure ..



# Transfer Reactions and Nuclear Structure: The Road Ahead

*A.O. Macchiavelli*  
*Lawrence Berkeley National Laboratory*

Many Thanks to Rod Clark, Paul Fallon, I. Yang Lee,  
David Radford and Alan Wuosmaa

**Nuclear Physics, The Core of Matter, the Fuel of Stars**  
**Argonne National Laboratory, IL, September 21-22, 2006**

# Outline

- Short introduction
- One-nucleon transfer reactions
  - Single-particle degrees of freedom
  - Spin-orbit splitting
  - Ab initio calculations
- Two-nucleon transfer reactions
  - Correlations
    - NP pairing
    - Pairing phase transition
- Heavy-ion transfer reactions
- Summary and conclusions

A central theme of study in nuclear structure has been the understanding of the elementary modes of excitation of the atomic nucleus, and their evolution with  $A$  (size),  $I$  (rotational frequency),  $(N-Z)$  (Isospin) and  $E^*$  (Temperature)

Transfer reactions have played a major role in this endeavor, particularly in the characterization of the single particle degrees of freedom and their correlations.

With the development of exotic beams and new instruments to handle the “reverse kinematics” nature of these reactions, there has been a *renaissance* of transfer reaction studies.

On Angular Distributions from  $(d, p)$  and  $(d, n)$   
Nuclear Reactions

S. T. BUTLER\*

Department of Mathematical Physics, University of Birmingham,  
Birmingham, England

October 30, 1950

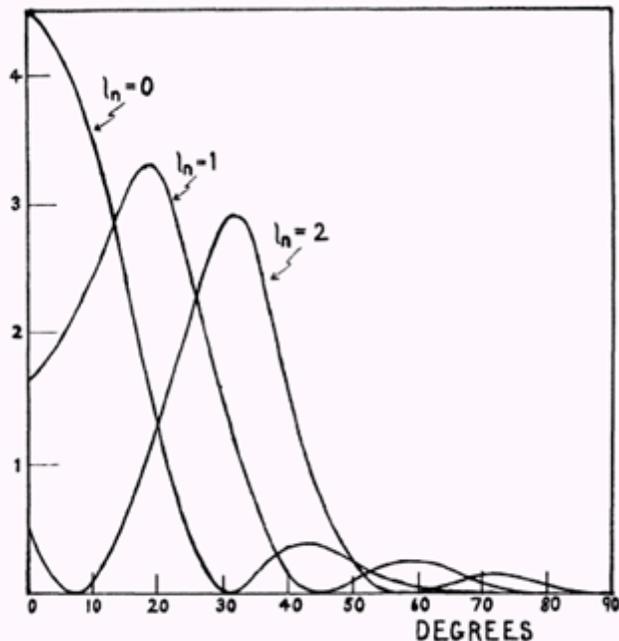


FIG. 1. Theoretical angular distributions for  $(d, p)$  and  $(d, n)$  reactions for different angular momentum transfers to the initial nucleus.

Angular distribution of the outgoing particles reflect the transferred angular momentum  $\Rightarrow$  l-value

Spectroscopic factors. Overlap between initial and final state  $\Rightarrow$  Test wave functions

Ideal tool to study the single-particle degree of freedom. Complementary to other processes that probe collective aspects.

**One of the important goals in the development of exotic beams is to study the evolution of the nuclear shell structure with neutron excess.**

**In fact, it is already known that in light nuclei the magic numbers are not as “robust” as originally conceived.**

**Of particular interest here is the spin-orbit force.**

## Is the Nuclear Spin-Orbit Interaction Changing with Neutron Excess?

J. P. Schiffer,<sup>1</sup> S. J. Freeman,<sup>1,2</sup> J. A. Caggiano,<sup>3</sup> C. Deibel,<sup>3</sup> A. Heinz,<sup>3</sup> C.-L. Jiang,<sup>1</sup> R. Lewis,<sup>3</sup> A. Parikh,<sup>3</sup> P. D. Parker,<sup>3</sup>  
K. E. Rehm,<sup>1</sup> S. Sinha,<sup>1</sup> and J. S. Thomas<sup>4</sup>

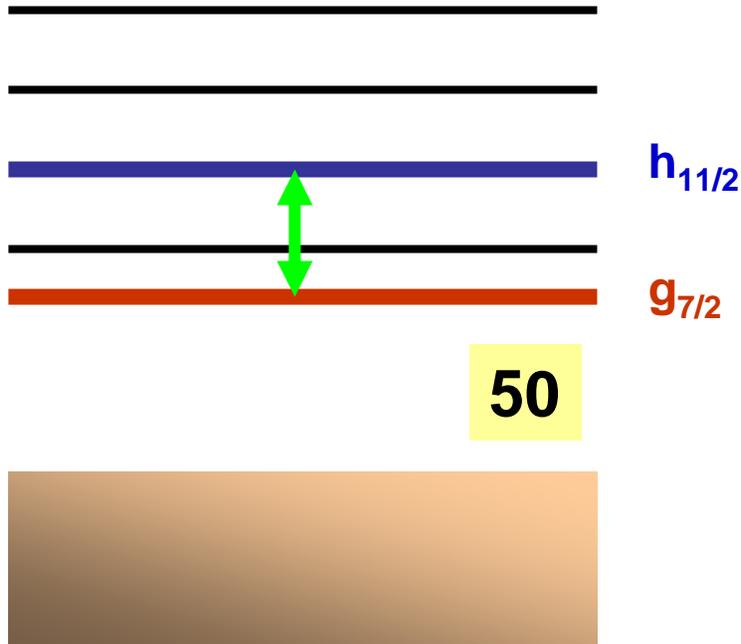
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<sup>2</sup>University of Manchester, Manchester M13 9PL, United Kingdom

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(Received 17 December 2003; published 20 April 2004)



$$\delta E \approx \hbar \omega_0 - 4.5V_{ls} / A^{2/3}$$

## Use $(\alpha,t)$ reaction to populate high- $l$ orbitals

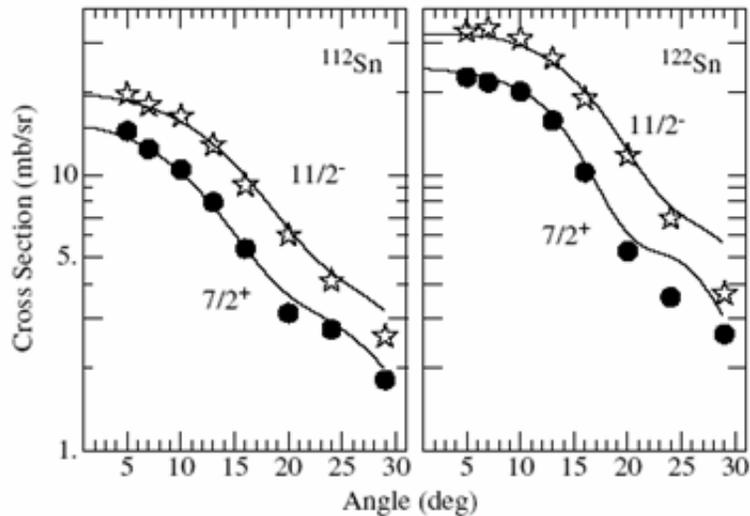
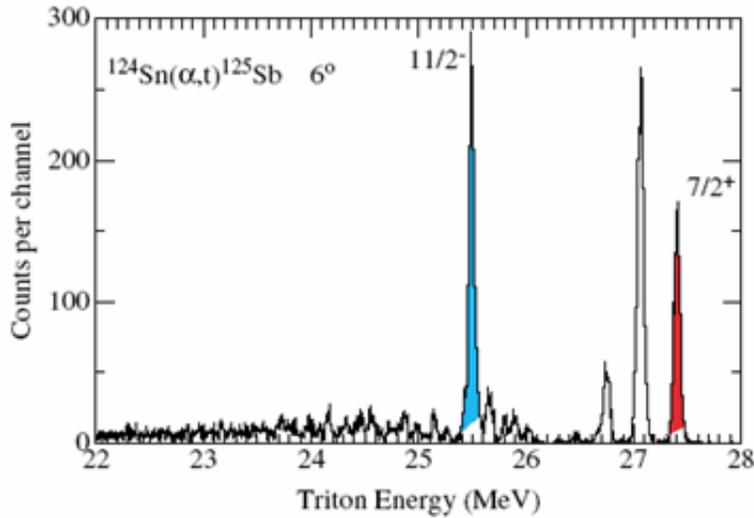
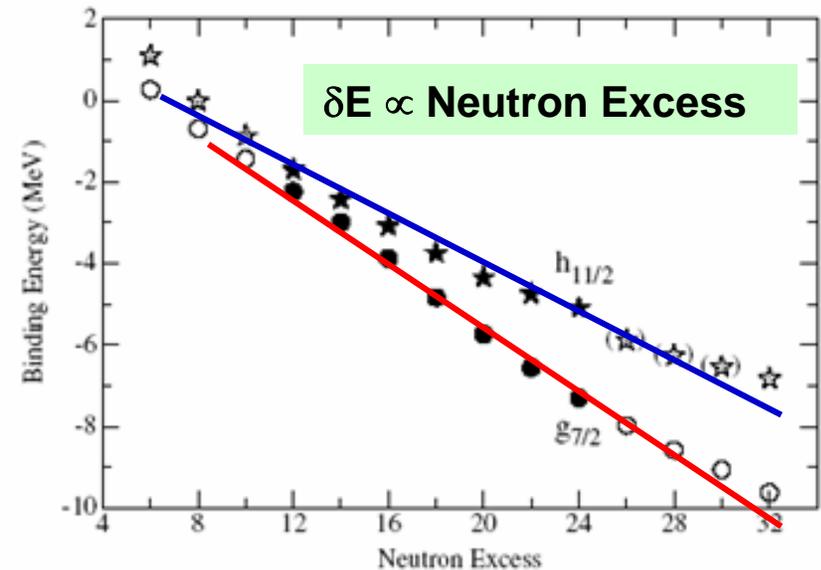


TABLE I. Cross sections (mb/sr) at  $6^\circ$  for the lowest  $7/2^+$  and  $11/2^-$  states, their ratios, and spectroscopic factors. The uncertainties in the cross sections are estimated at 10% and those in the ratio, at about 5%. The accuracy of the relative spectroscopic factors are estimated at 5%.

Target	$7/2^+$	$11/2^-$	Ratio	$C^2S_{7/2}$	$C^2S_{11/2}$
$^{112}\text{Sn}$	14.6	21.4	1.47	0.99	0.84
$^{114}\text{Sn}$	19.6	27.3	1.39	1.10	0.93
$^{116}\text{Sn}$	19.7	30.9	1.57	0.95	0.97
$^{118}\text{Sn}$	20.4	33.5	1.64	0.88	0.99
$^{120}\text{Sn}$	27.9	39.4	1.41	1.13	1.12
$^{122}\text{Sn}$	24.6	35.5	1.45	0.98	1.00
$^{124}\text{Sn}$	24.7	39.2	1.59	1.00	1.12



# QUENCHING OF THE $2p_{1/2}-2p_{3/2}$ PROTON SPIN-ORBIT SPLITTING IN THE Sr-Zr REGION

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*Department of Physics and Astronomy, San Francisco State University, San Francisco, CA 94132, USA*

S. PITTEL <sup>2</sup>

*Bartol Research Foundation of the Franklin Institute, University of Delaware, Newark, DE 19716, USA*

and

A. ETCHEGOYEN

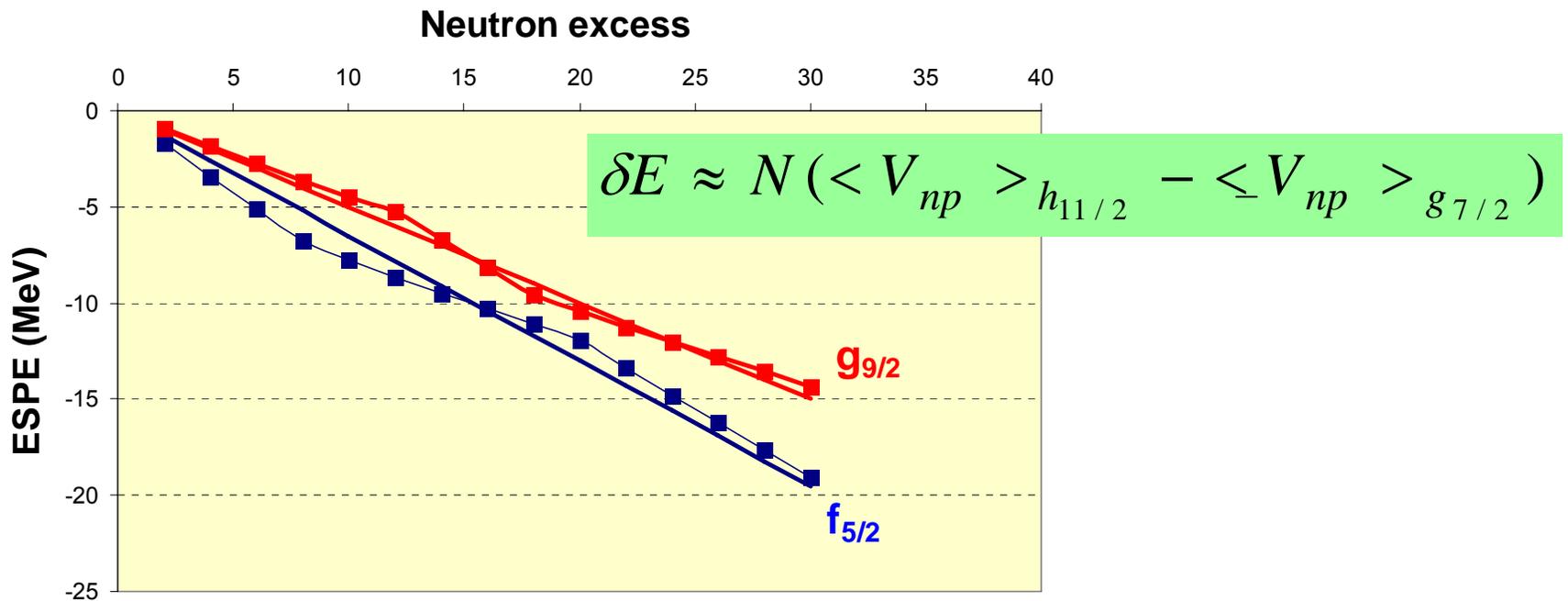
*Departamento de Física, Comisión Nacional de Energía Atómica, Av. del Libertador 8250, 1429 Buenos Aires, Argentina*

<sup>95</sup>Y

$$\Delta\epsilon(N = 56) - \Delta\epsilon(N = 50)$$

<sup>89</sup>Y

$$= 6 \left( \overline{V}_{p_{1/2}-d_{5/2}}^{n-p} - \overline{V}_{p_{3/2}-p_{5/2}}^{n-p} \right),$$



**Reduction of the Spin-Orbit Splittings at the  $N = 28$  Shell Closure**

L. Gaudefroy,<sup>1,2</sup> O. Sorlin,<sup>2,1</sup> D. Beaumel,<sup>1</sup> Y. Blumenfeld,<sup>1</sup> Z. Dombrádi,<sup>3</sup> S. Fortier,<sup>1</sup> S. Franchoo,<sup>1</sup> M. Gélín,<sup>2</sup> J. Gibelin,<sup>1</sup> S. Grévy,<sup>2</sup> F. Hammache,<sup>1</sup> F. Ibrahim,<sup>1</sup> K. W. Kemper,<sup>4</sup> K.-L. Kratz,<sup>5,6</sup> S. M. Lukyanov,<sup>7</sup> C. Monrozeau,<sup>1</sup> L. Nalpas,<sup>8</sup> F. Nowacki,<sup>9</sup> A. N. Ostrowski,<sup>5,6</sup> T. Otsuka,<sup>10</sup> Yu.-E. Penionzhkevich,<sup>7</sup> J. Piekarewicz,<sup>4</sup> E. C. Pollacco,<sup>8</sup> P. Roussel-Chomaz,<sup>2</sup> E. Rich,<sup>1</sup> J. A. Scarpaci,<sup>1</sup> M. G. St. Laurent,<sup>2</sup> D. Sohler,<sup>11</sup> M. Stanoiu,<sup>12</sup> T. Suzuki,<sup>13</sup> E. Tryggestad,<sup>1</sup> and D. Verney<sup>1</sup>

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**$^{46}\text{Ar}(d,p)^{47}\text{Ar}$  reaction at SPIRAL and MUST detector array**

**In recent years substantial progress has been made in the development of *Ab Initio* theories of light nuclei.**

**The Quantum Monte Carlo method has been very successful in reproducing BE and excitation spectra of light nuclei starting with the basic NN force and 3-body force.**

Neutron Spectroscopic Factors in  ${}^9\text{Li}$  from  ${}^2\text{H}({}^8\text{Li}, p){}^9\text{Li}$ 

A. H. Wuosmaa,<sup>1</sup> K. E. Rehm,<sup>2</sup> J. P. Greene,<sup>2</sup> D. J. Henderson,<sup>2</sup> R. V. F. Janssens,<sup>2</sup> C. L. Jiang,<sup>2</sup> L. Jisonna,<sup>3</sup> E. F. Moore,<sup>2</sup>  
 R. C. Pardo,<sup>2</sup> M. Paul,<sup>4</sup> D. Peterson,<sup>2</sup> Steven C. Pieper,<sup>2</sup> G. Savard,<sup>2</sup> J. P. Schiffer,<sup>2</sup> R. E. Segel,<sup>3</sup>  
 S. Sinha,<sup>2</sup> X. Tang,<sup>2</sup> and R. B. Wiringa<sup>2</sup>

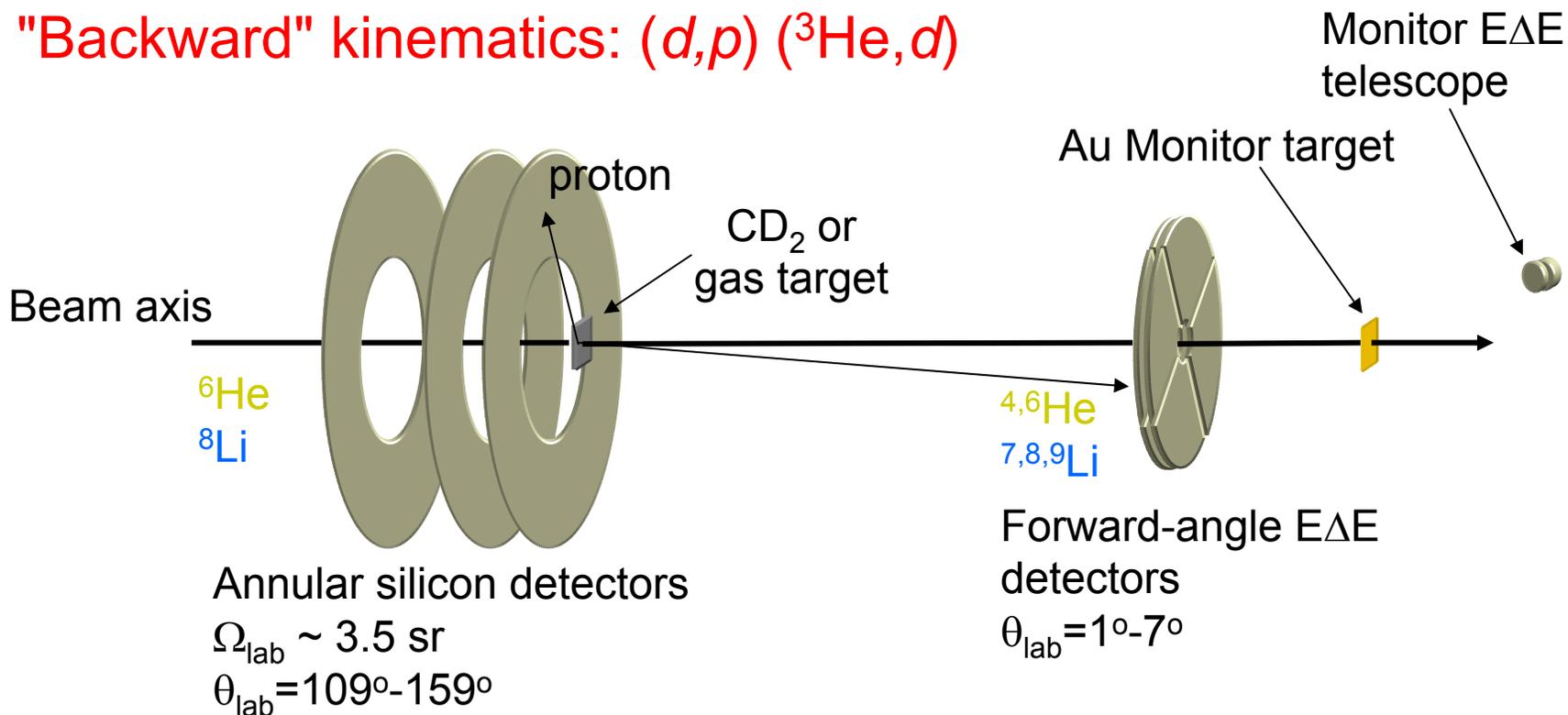
<sup>1</sup>Physics Department, Western Michigan University, Kalamazoo, Michigan 49008-5252, USA

<sup>2</sup>Physics Division, Argonne National Laboratory, 9700 S. Cass Ave, Argonne Illinois 60439, USA

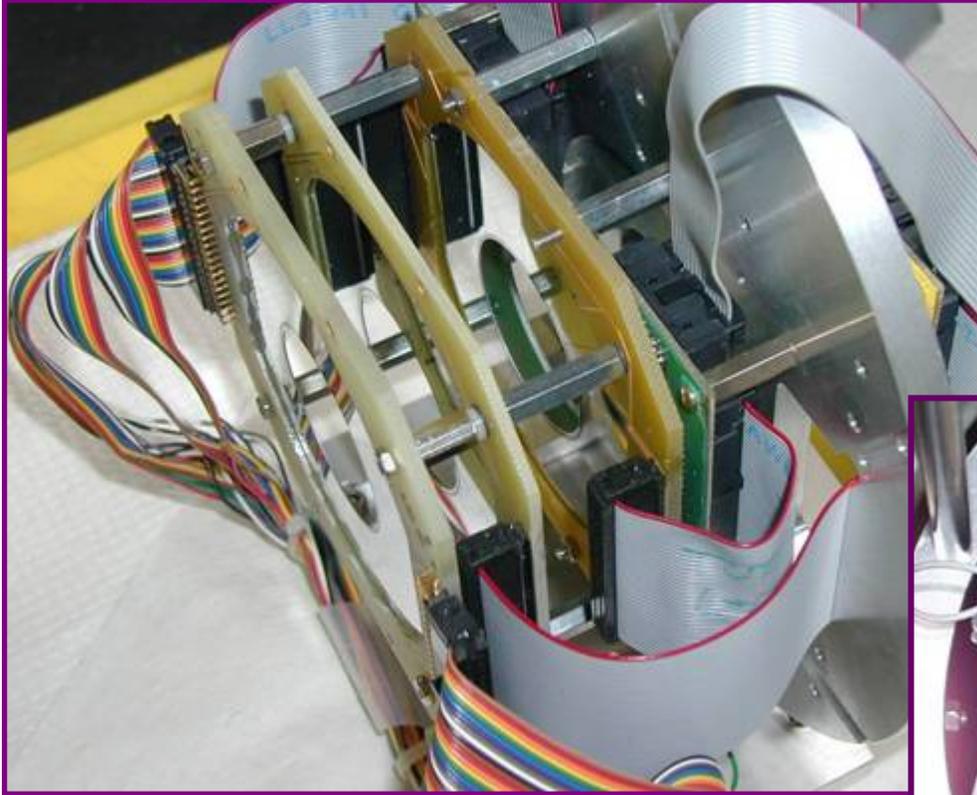
<sup>3</sup>Physics Department, Northwestern University, Evanston, Illinois 60208, USA

<sup>4</sup>Hebrew University, Jerusalem, Israel 91904

## "Backward" kinematics: $(d,p)$ ( ${}^3\text{He}, d$ )

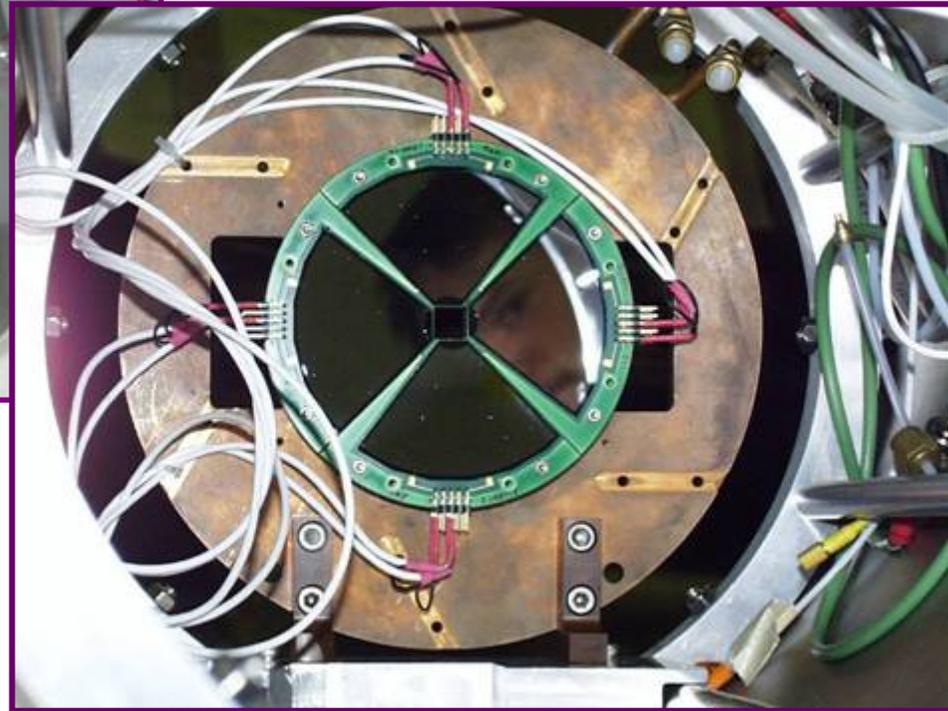
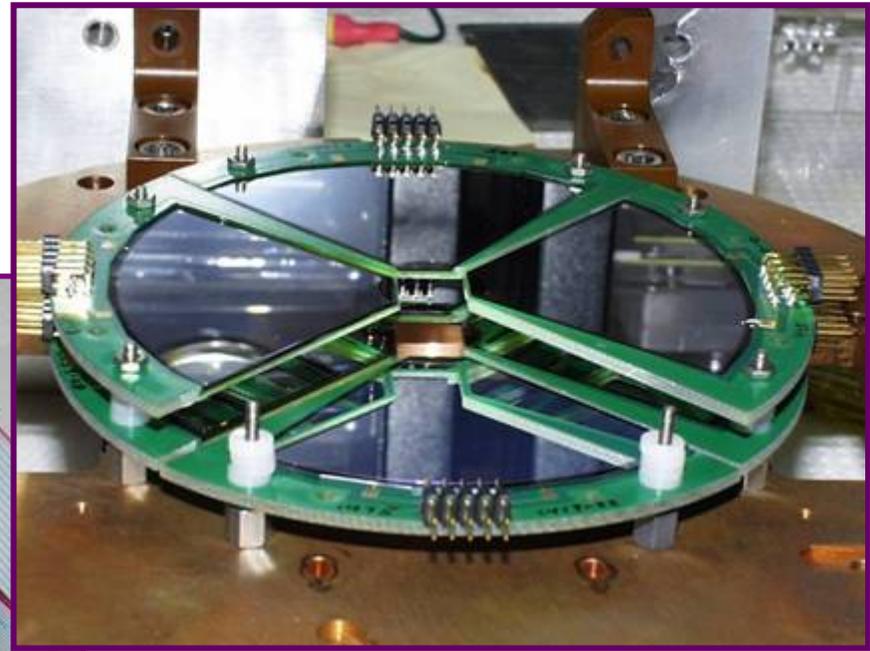


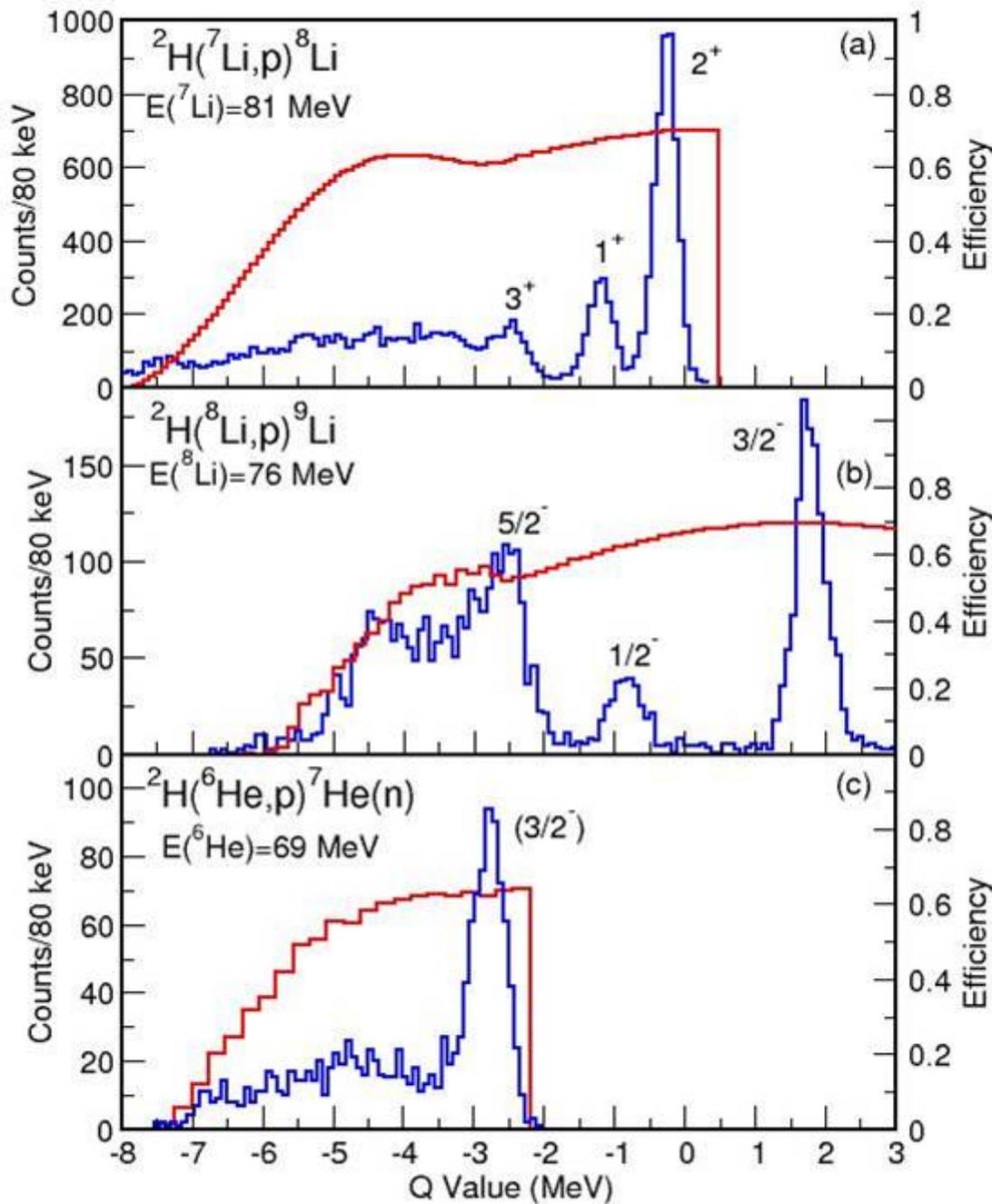
## Detectors



Segmented proton detectors

500 $\mu\text{m}$ /1000 $\mu\text{m}$  silicon E $\Delta$ E telescope



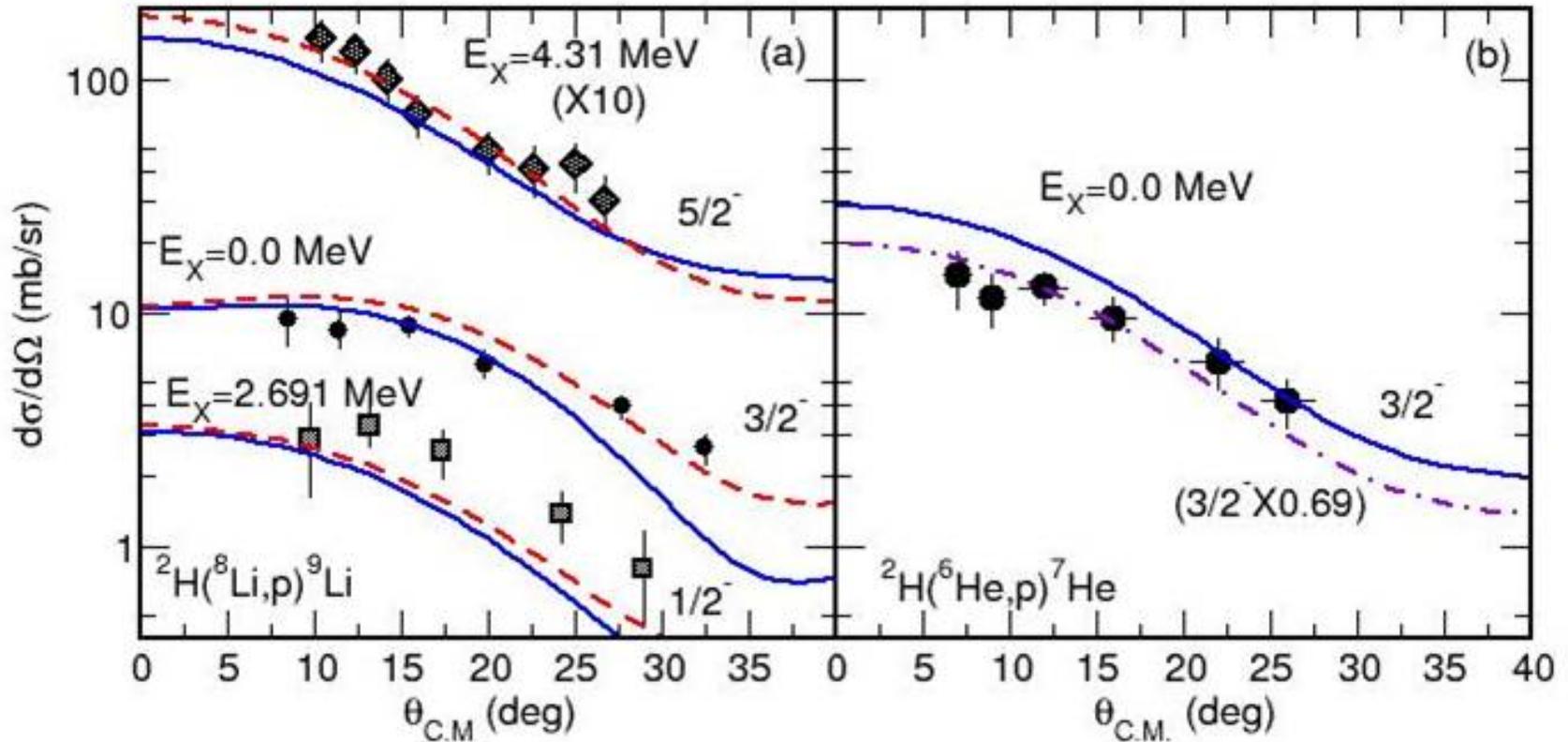


Q-value spectra from (d,p)

$p$ -Li or  $p$ - ${}^6\text{He}$  coincidences  
 Efficiency from Monte Carlo simulations

PRL **94**, 082502 (2005),  
 PRC **72**, 061301(R) (2005)

## $(d,p)$ Angular Distributions - narrow states



${}^2\text{H}({}^8\text{Li},p){}^9\text{Li}$  DWBA calculations: **Red**,  
**blue** curves: QMC predictions with  
different OMP, no extra normalization

${}^2\text{H}({}^6\text{He},p){}^7\text{He}_{g.s.}$  DWBA calculations  
QMC calculations. **Blue**: no normalization  
**violet**- QMC X 0.69.

Optical-model parameters from Schiffer et al, PRC **164**

**Shell structure near  $^{132}\text{Sn}$ :**

**Neutron levels in  $^{133}\text{Sn}$  via the  $^{132}\text{Sn}(d,p)$**

# Successful “proof of principle” experiment with $^{124}\text{Sn}$

PHYSICAL REVIEW C 70, 067602 (2004)

## Study of the $^{124}\text{Sn}(d,p)$ reaction in inverse kinematics close to the Coulomb barrier

K. L. Jones,<sup>1</sup> R. L. Kozub,<sup>2</sup> C. Baktash,<sup>3</sup> D. W. Bardayan,<sup>3</sup> J. C. Blackmon,<sup>3</sup> W. N. Catford,<sup>4</sup> J. A. Cizewski,<sup>1</sup> R. P. Fitzgerald,<sup>5</sup> M. S. Johnson,<sup>6</sup> R. J. Livesay,<sup>7</sup> Z. Ma,<sup>8</sup> C. D. Nesaraja,<sup>2</sup> D. Shapira,<sup>3</sup> M. S. Smith,<sup>3</sup> J. S. Thomas,<sup>1</sup> and D. W. Visser<sup>5</sup>

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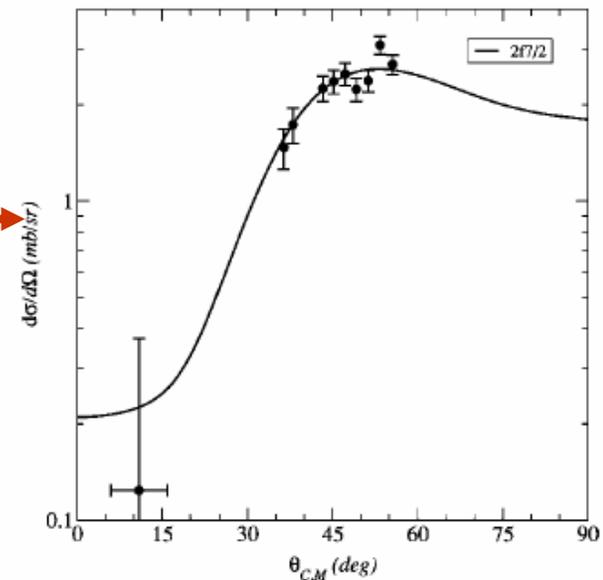
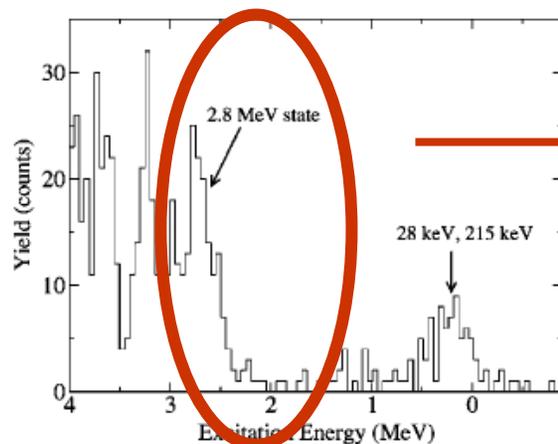
<sup>4</sup>Physics Department, University of Surrey, Guildford, Surrey, GU2 7XH, United Kingdom

<sup>5</sup>Department of Physics and Astronomy, University of North Carolina, Chapel Hill, North Carolina 27599, USA

<sup>6</sup>Oak Ridge Associated Universities, Oak Ridge, Tennessee 37831, USA

<sup>7</sup>Colorado School of Mines, Golden, Colorado 80401, USA

<sup>8</sup>Department of Physics and Astronomy, University of Tennessee, Knoxville, Tennessee 37996, USA



**THE** experiment with an improved detector array should be running, as we speak, at HRIBF

## **PAIRING IN NUCLEI**

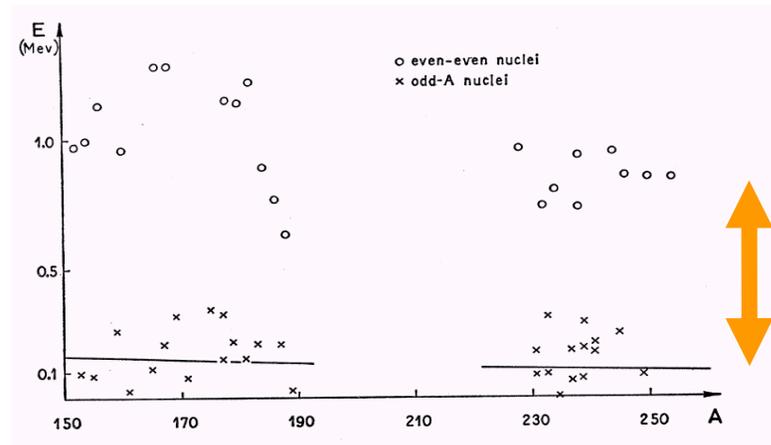
## Possible Analogy between the Excitation Spectra of Nuclei and Those of the Superconducting Metallic State

A. BOHR, B. R. MOTTELSON, AND D. PINES\*

*Institute for Theoretical Physics, University of Copenhagen, Copenhagen, Denmark, and Nordisk Institut for Teoretisk Atomfysik, Copenhagen, Denmark*

(Received January 7, 1958)

The evidence for an energy gap in the intrinsic excitation spectrum of nuclei is reviewed. A possible analogy between this effect and the energy gap observed in the electronic excitation of a superconducting metal is suggested.

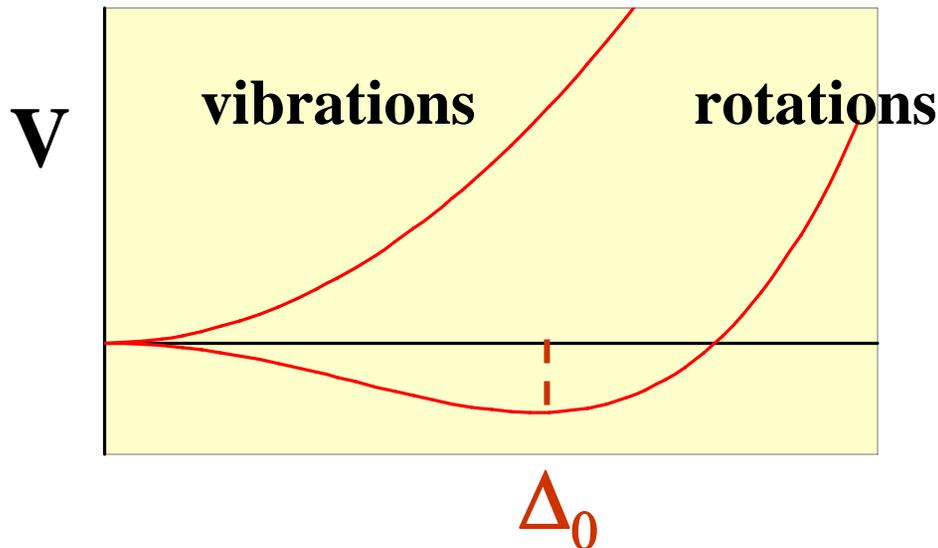


- Gap in the excitation spectra of even- $A$  nuclei
- Odd-even mass differences
- Rotational moments of inertia

Relevance to other finite Fermion systems such as  $^3\text{He}$  clusters, Fermi-gas condensates, quantum dots, metal clusters, ....

## Two-nucleon Transfer Reactions

Generalized densities  $a+a^+$ ,  $aa$  represent the pair field and in close analogy to the collective excitations corresponding to the ordinary density, they can give rise to collective modes.



$$\Delta = G \langle \sum a_v^+ a_{\bar{v}}^+ \rangle$$

Two particle transfer reactions like (t,p) or (p,t), where 2 neutrons are deposited or picked up at the same point in space provide an specific tool to probe the amplitude of this collective motion. The transition operator  $\langle f | a+a^+ | i \rangle$  will be proportional to the pair density of the nucleus.

**NOTE ON THE TWO-NUCLEON STRIPPING REACTION**

SHIRO YOSHIDA †

*Radiation Laboratory, University of Pittsburgh, Pittsburgh, Pennsylvania ††*

Received 9 February 1962

**Abstract:** The magnitude of the two-nucleon s interaction model. The calculation also is a types of reaction a collective enhancement o

\* *Proceedings Int. Symp. on nuclear structure*<sup>20</sup>  
(IAEA, Vienna) 1968  
p. 179

**PAIR CORRELATIONS AND  
DOUBLE TRANSFER REACTIONS**

A. BOHR  
THE NIELS BOHR INSTITUTE,  
UNIVERSITY OF COPENHAGEN,  
COPENHAGEN, DENMARK

PHYSICS REPORT

**ISOVECTOR PAIRING VIBRATIONS**

D.R. BES

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State University of New York at Stony Brook, Physics Department, Stony Brook, New York 11794, USA and  
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*The Niels Bohr Institute, University of Copenhagen, DK-2100 Copenhagen Ø, Denmark and  
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and

Ole HANSEN and O. NATHAN

*The Niels Bohr Institute, University of Copenhagen, DK-2100 Copenhagen Ø, Denmark*

## **Neutron Proton Pairing**

# NP Pairing

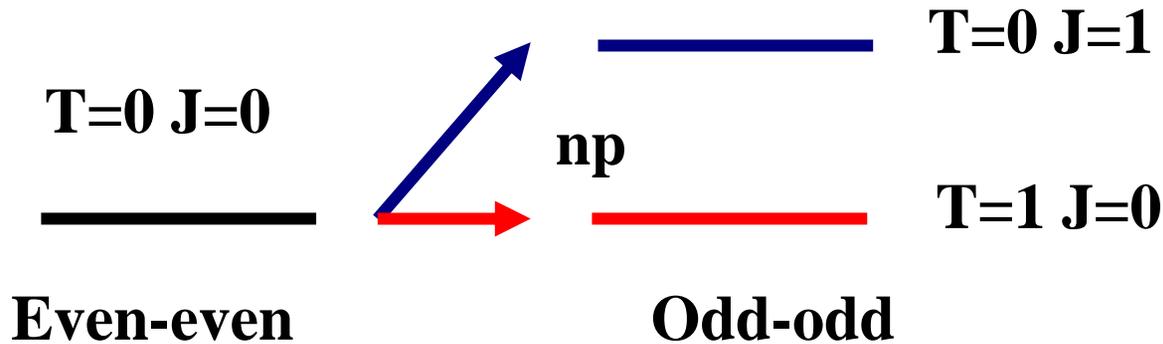
- ★ N=Z nuclei, unique systems to study  $np$  correlations
  - As you move out of N=Z  $nn$  and  $pp$  pairs are favored
- ★ Role of isoscalar ( $T=0$ ) and isovector ( $T=1$ ) pairing
  - Large spatial overlap of  $n$  and  $p$
  - Pairing vibrations (normal system)
  - Pairing rotations (superfluid system)
- ★ Does isoscalar pairing give rise to collective modes?
- ★ What is (are) the “smoking-gun(s)”?
  - Binding energy differences
  - Ground states of odd-odd self-conjugate nuclei
  - Rotational properties: moments of inertia, alignments
  - Two-particle transfer cross-sections



# (3He,p) Transfer Reactions

$$\sigma \propto < 1 | T | 0 >^2 \sim \Omega \quad \text{or} \quad \Omega^2$$

$\sigma ?$



$(^3\text{He}, p)$

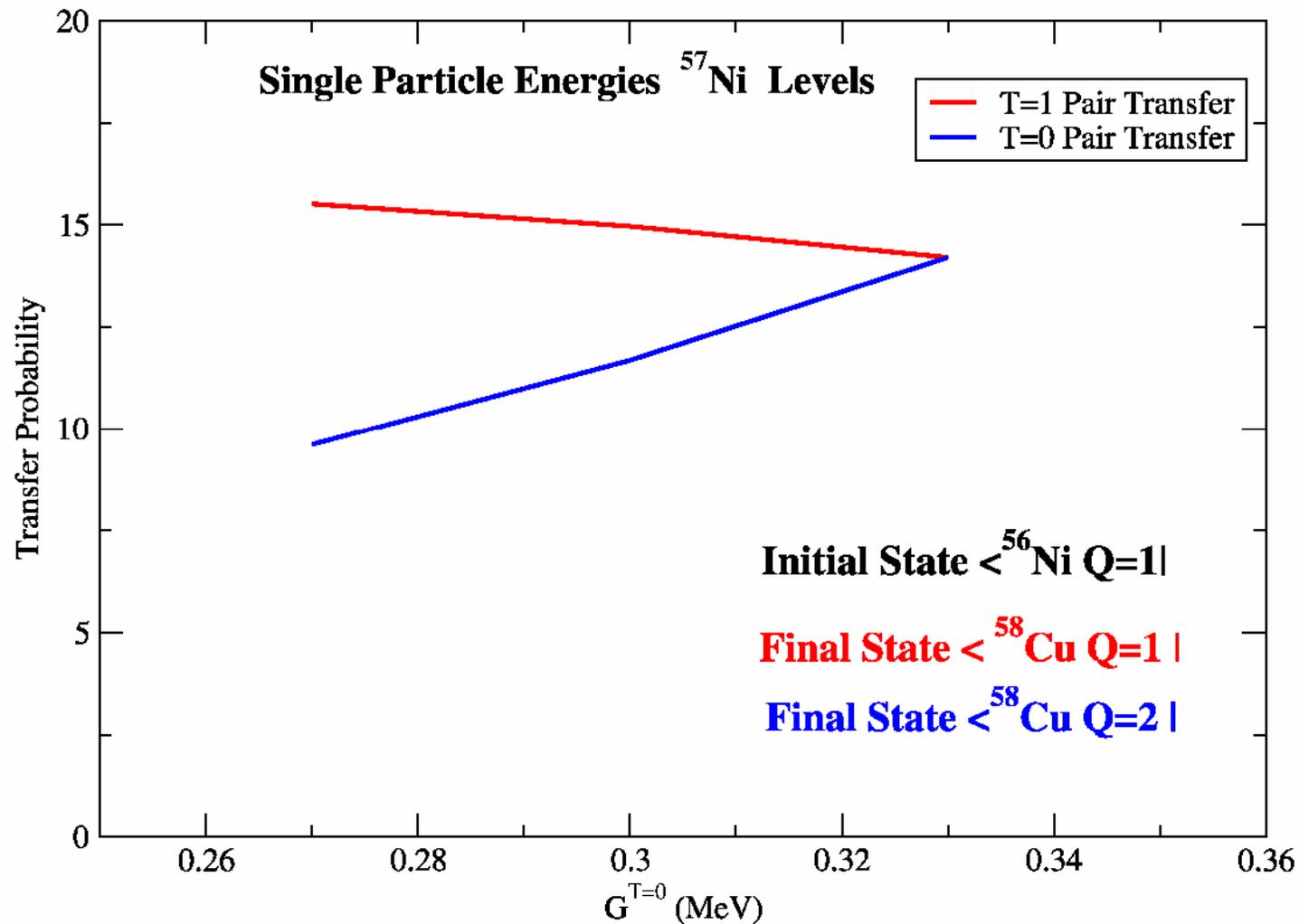
$L=0$  transfer – forward peaked

Measure the  $np$  transfer cross section to  $T=1$  and  $T=0$  states

Both absolute  $\sigma(T=0)$  and  $\sigma(T=1)$  and relative  $\sigma(T=0) / \sigma(T=1)$  tell us about the character and strength of the correlations

# n-p Pair Transfer Probability

$$G^{T=1} = 0.33 \text{ MeV}$$



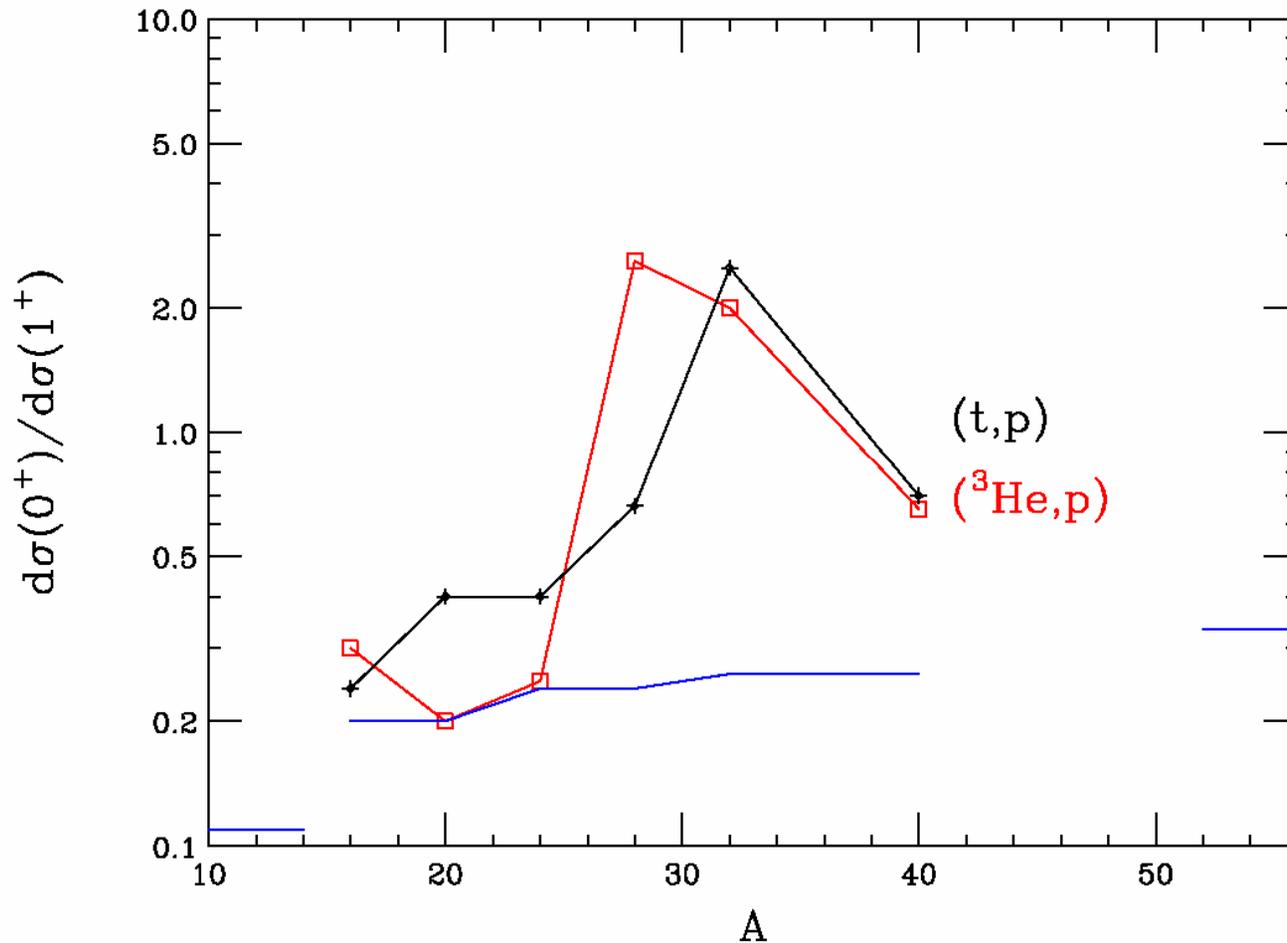
Deuteron Transfer in  $N = Z$  NucleiP. Van Isacker,<sup>1</sup> D.D. Warner,<sup>2</sup> and A. Frank<sup>3</sup><sup>1</sup>*Grand Accélérateur National d'Ions Lourds, B.P. 55027, F-14076 Caen Cedex 5, France*<sup>2</sup>*CCLRC Daresbury Laboratory, Daresbury, Warrington WA4 4AD, United Kingdom*<sup>3</sup>*Instituto de Ciencias Nucleares, UNAM, Apdo. Postal 70-543, 04510 México, D.F. Mexico*

(Received 14 September 2004; published 29 April 2005)

TABLE I. Predicted deuteron-transfer intensities  $C_T^2$  between even-even (EE) and odd-odd (OO)  $N = Z$  nuclei in the  $SU(4)$  ( $b/a = 0$ ) and  $U_T(3) \otimes U_S(3)$  ( $|b/a| \gg 1$ ) limits.

Limit	Reaction	$C_{T=0}^2$	$C_{T=1}^2$
$b/a = 0$	$EE \rightarrow OO_{T=0}$	$\frac{1}{2}(N_b + 6)$	0
	$EE \rightarrow OO_{T=1}$	0	$\frac{1}{2}(N_b + 6)$
$b/a \ll -1$	$EE \rightarrow OO_{T=0}$	$N_b + 3$	0
	$EE \rightarrow OO_{T=1}$	0	3
$b/a \gg +1$	$EE \rightarrow OO_{T=0}$	3	0
	$EE \rightarrow OO_{T=1}$	0	$N_b + 3$

# Systematic of ( $^3\text{He},p$ ) and $(t,p)$ reactions in stable $N=Z$ nuclei



Single-particle estimate  $\sim (\text{spin}) \times (^3\text{He}) \times (\text{LS} \rightarrow \text{jj})$

## Study of the $^{56}\text{Ni}(d,p)^{57}\text{Ni}$ Reaction and the Astrophysical $^{56}\text{Ni}(p,\gamma)^{57}\text{Cu}$ Reaction Rate

K. E. Rehm,<sup>1</sup> F. Borasi,<sup>1</sup> C. L. Jiang,<sup>1</sup> D. Ackermann,<sup>1</sup> I. Ahmad,<sup>1</sup> B. A. Brown,<sup>2</sup> F. Brumwell,<sup>1</sup> C. N. Davids,<sup>1</sup>  
 P. Decroock,<sup>1</sup> S. M. Fischer,<sup>1</sup> J. Görres,<sup>3</sup> J. Greene,<sup>1</sup> G. Hackmann,<sup>1</sup> B. Harss,<sup>1</sup> D. Henderson,<sup>1</sup> W. Henning,<sup>1</sup>  
 R. V. F. Janssens,<sup>1</sup> G. McMichael,<sup>1</sup> V. Nanal,<sup>1</sup> D. Nisius,<sup>1</sup> J. Nolen,<sup>1</sup> R. C. Pardo,<sup>1</sup> M. Paul,<sup>4</sup> P. Reiter,<sup>1</sup> J. P. Schiffer,<sup>1</sup>  
 D. Seweryniak,<sup>1</sup> R. E. Segel,<sup>5</sup> M. Wiescher,<sup>3</sup> and A. H. Wuosmaa<sup>1</sup>

<sup>1</sup>Argonne National Laboratory, Argonne, Illinois 60439

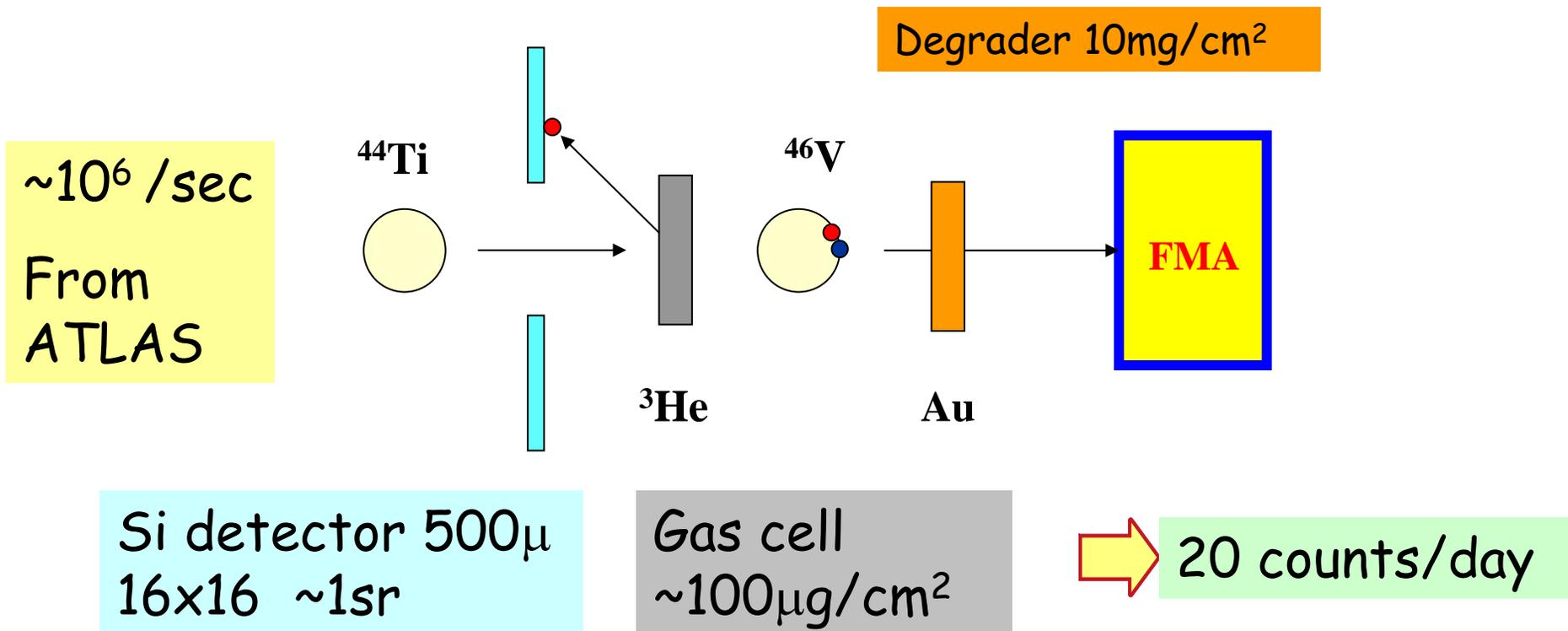
<sup>2</sup>Michigan State University, East Lansing, Michigan 48824

<sup>3</sup>University of Notre Dame, South Bend, Indiana 46556

<sup>4</sup>Hebrew University, Jerusalem, Israel

<sup>5</sup>Northwestern University, Evanston, Illinois 60208

(Received 29 August 1997)



# Proof of principle

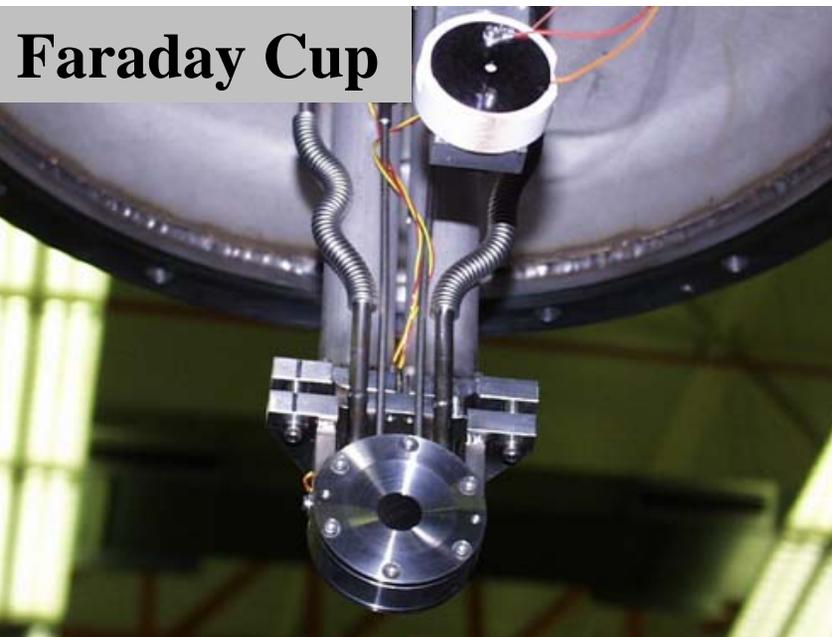
A.O.Macchiavelli<sup>1</sup>, E.Rehm<sup>2</sup>, P.Fallon<sup>1</sup>, M.Cromaz<sup>1</sup>, I.Ahmad<sup>1</sup>, C.N.Davis<sup>2</sup>, J.Greene<sup>2</sup>, R.V.F.Janssens<sup>2</sup>, C.L.Jiang<sup>2</sup>, E.F.Moore<sup>2</sup>, R.Pardo<sup>2</sup>, D.Seweryniak<sup>2</sup>, J.P.Schiffer<sup>2</sup>, A.Wuosmaa<sup>3</sup>, J.Cizewski<sup>3</sup>,

<sup>1</sup> Lawrence Berkeley National Laboratory

<sup>2</sup> Argonne National Laboratory

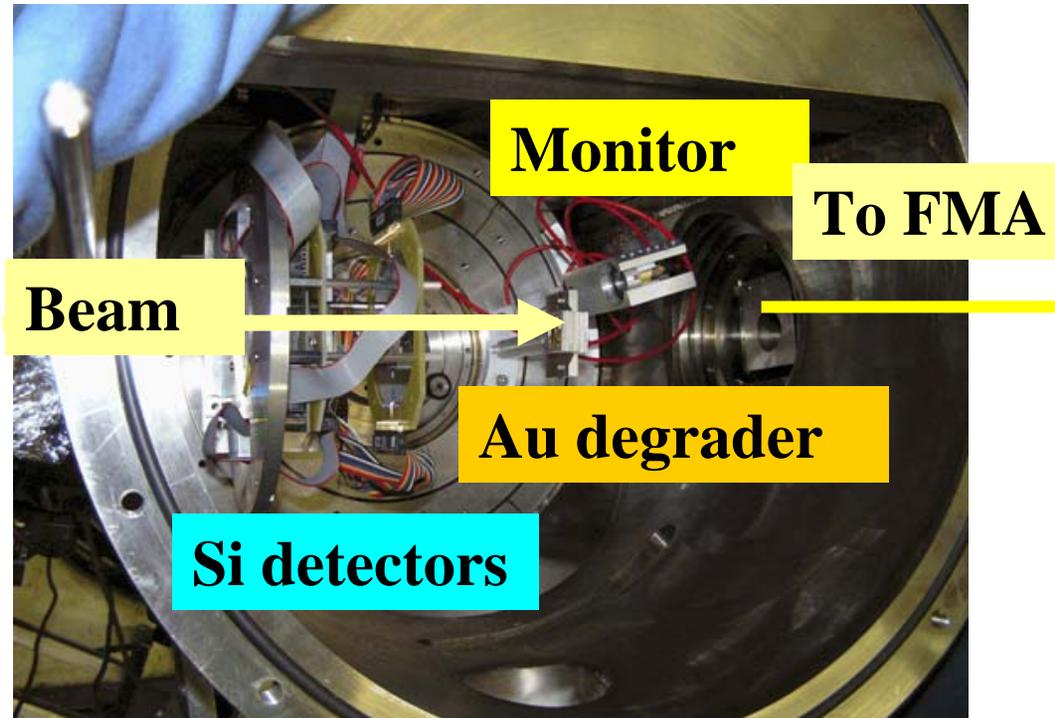
<sup>3</sup> Western Michigan University

<sup>3</sup> Rutgers University



**Faraday Cup**

**Gas Cell**



**Beam**

**Monitor**

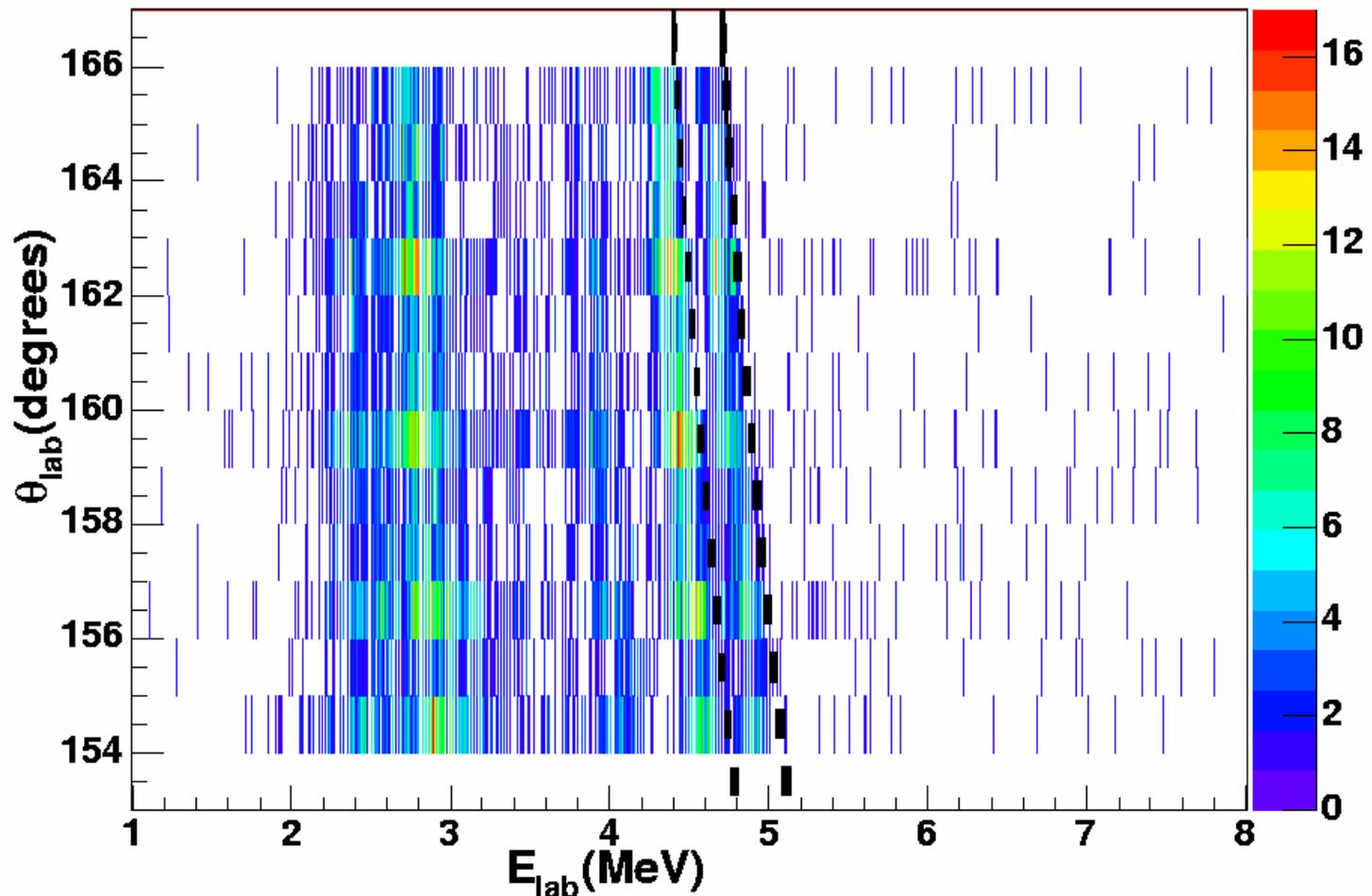
**To FMA**

**Au degrader**

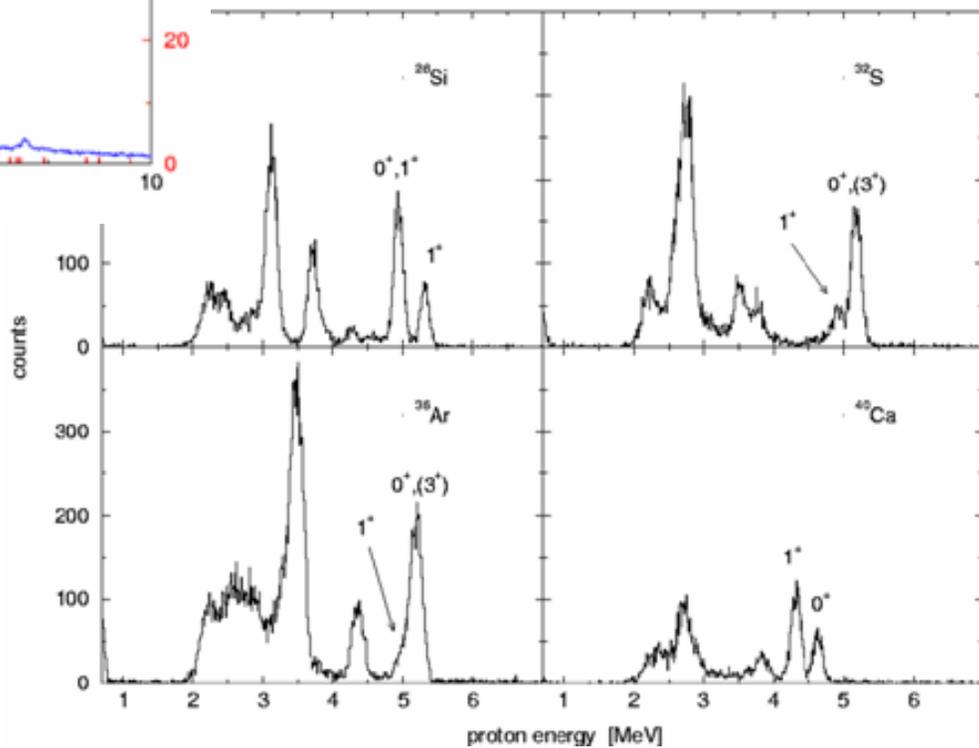
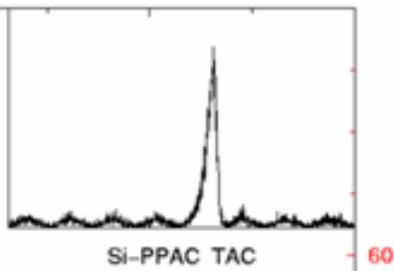
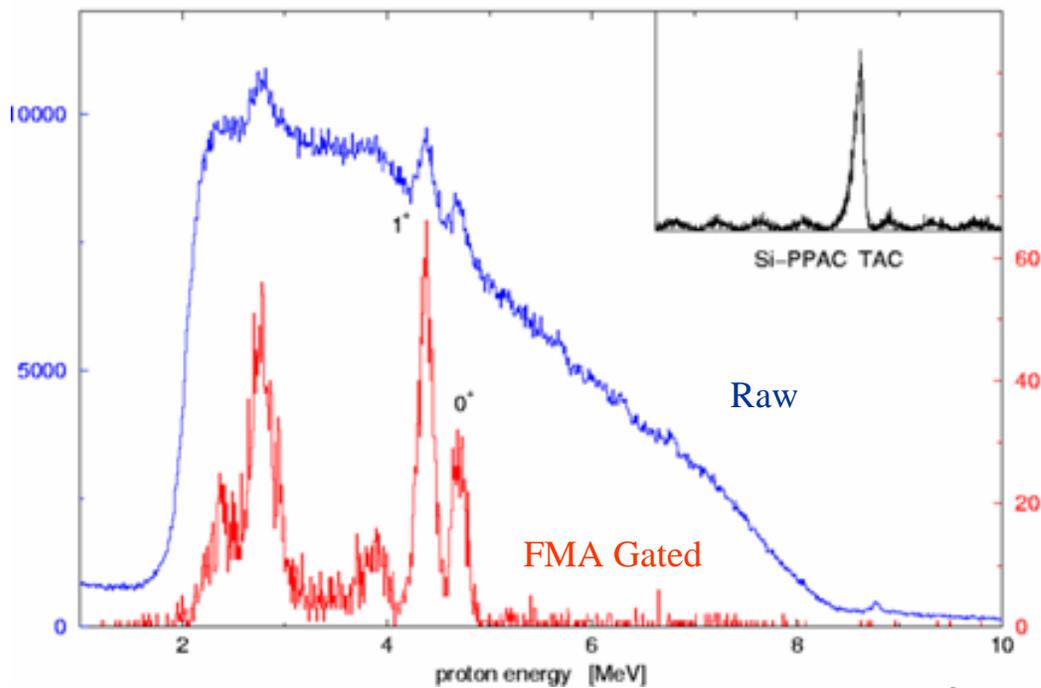
**Si detectors**

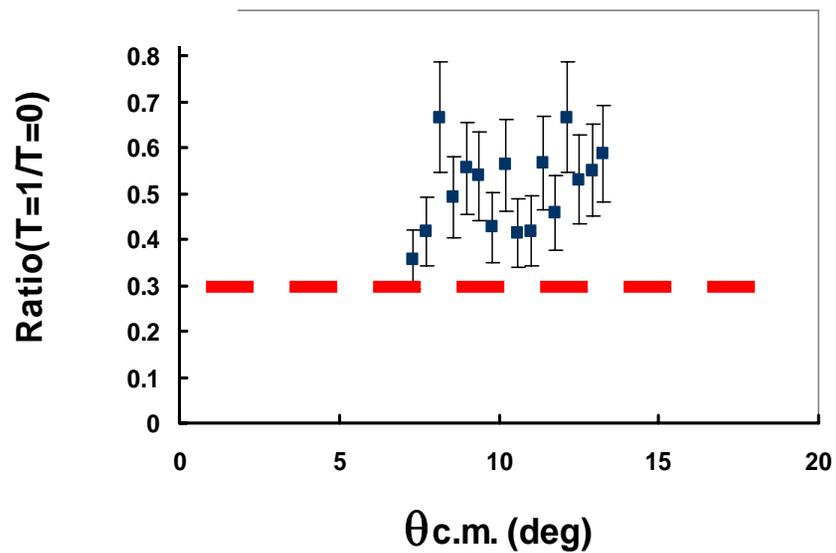
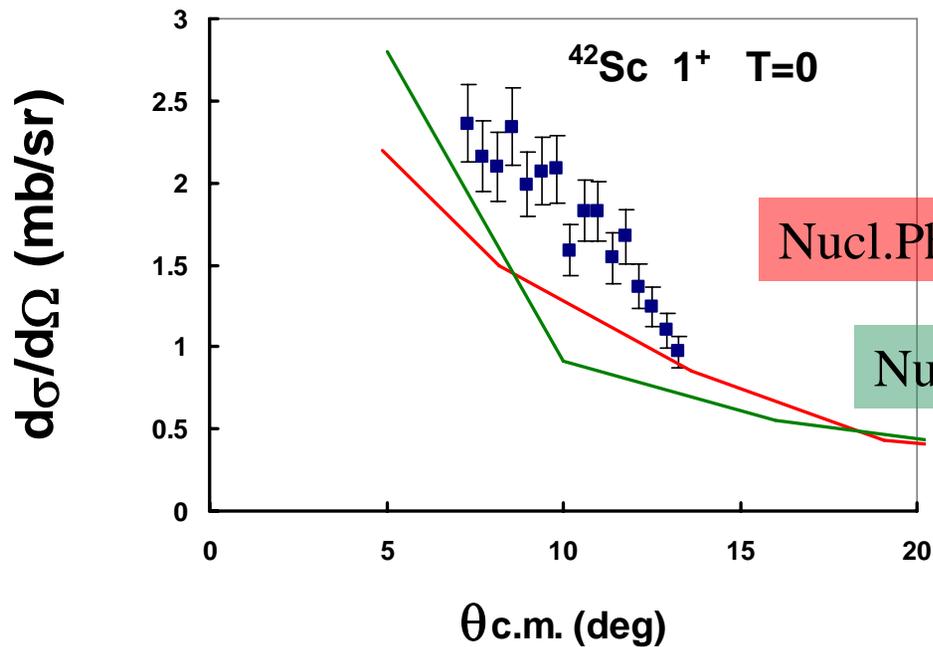
# Proof of principle

Inverse kinematics -  
Successful test experiment with stable beams



$^{40}\text{Ca}(^3\text{He},p)$  @ 220 MeV

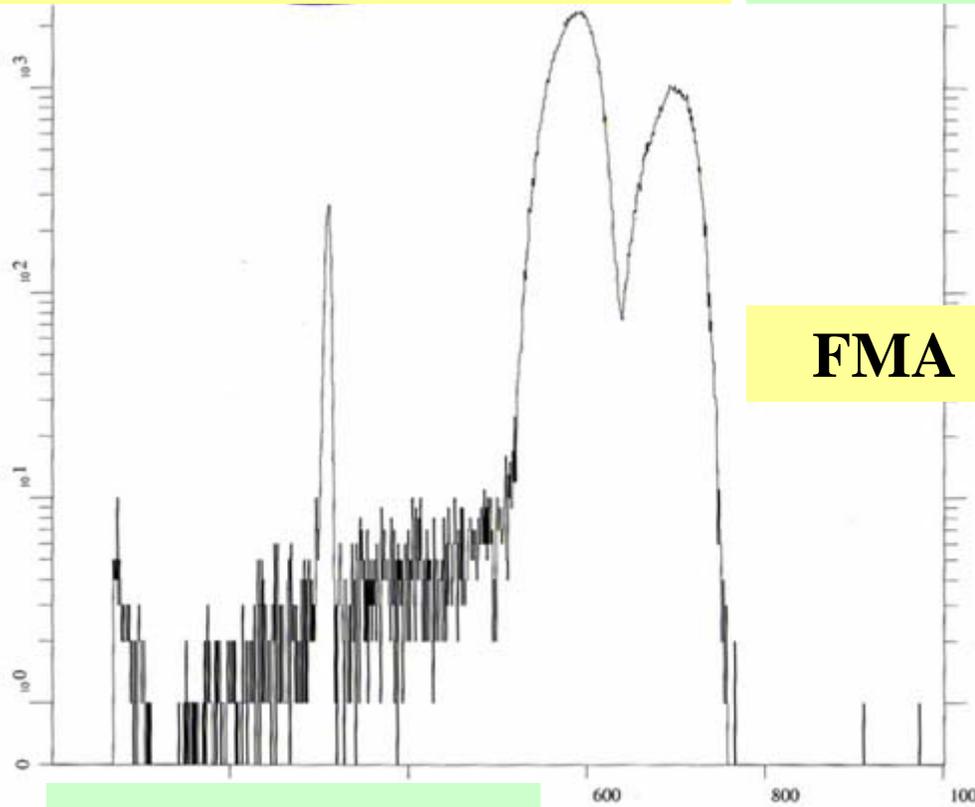




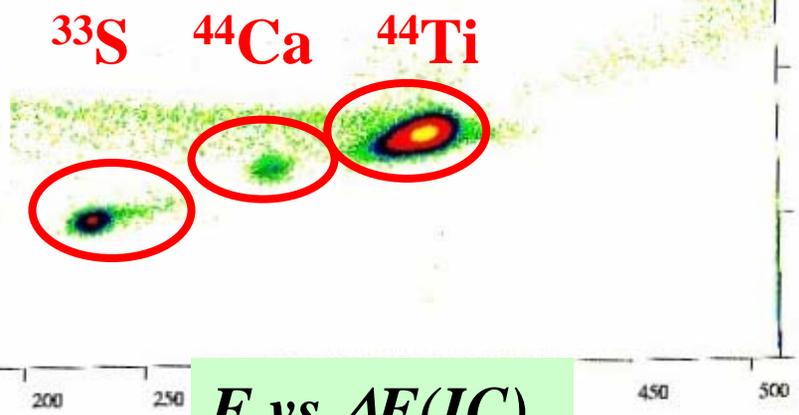
# 242 MeV $^{44}\text{Ti}$ Beam

$^{44}\text{Ti} / ^{44}\text{Ca} = 2.5 \sim 890/\text{sec}$

$\sim 1.1 \times 10^6 \text{ }^{44}\text{Ti}/\text{sec}$



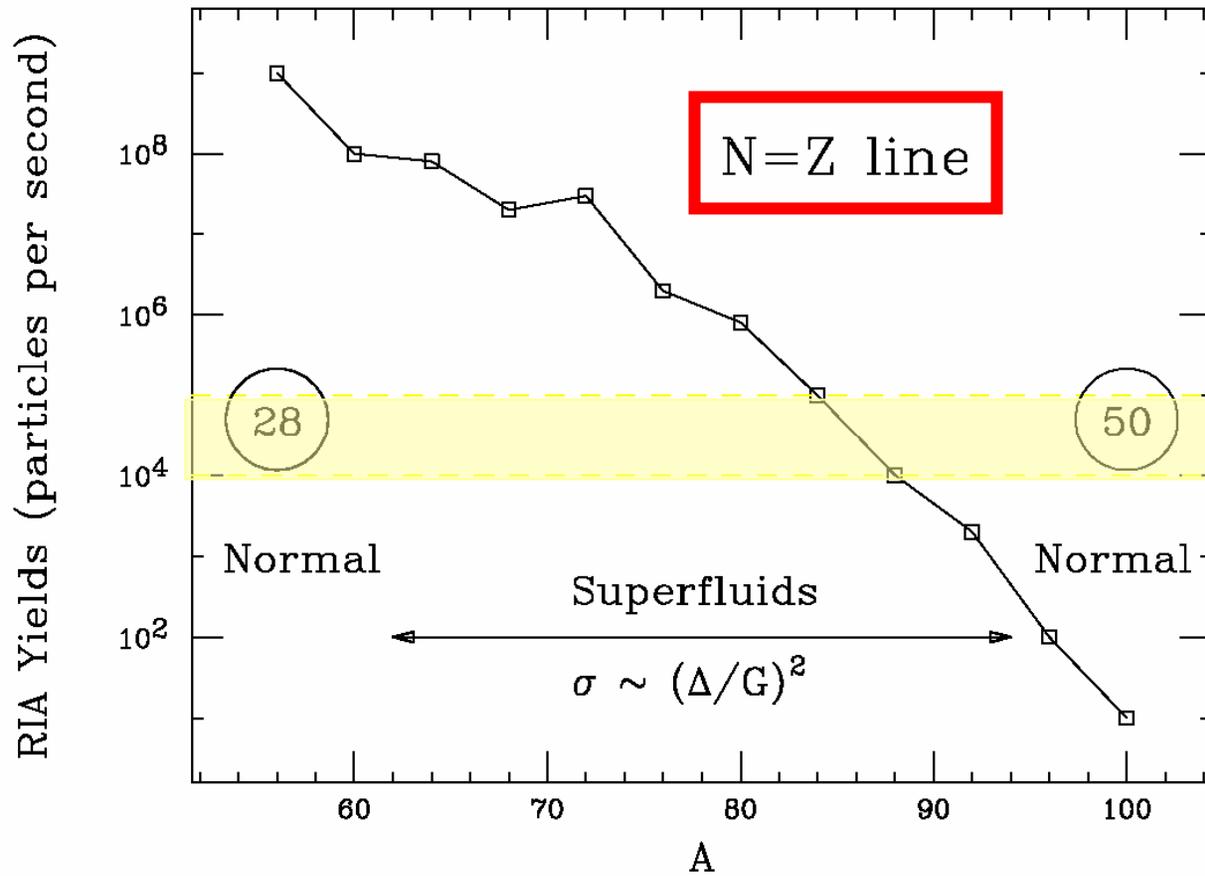
FMA Settings 107MeV,  $A=44$ ,  $q=20$



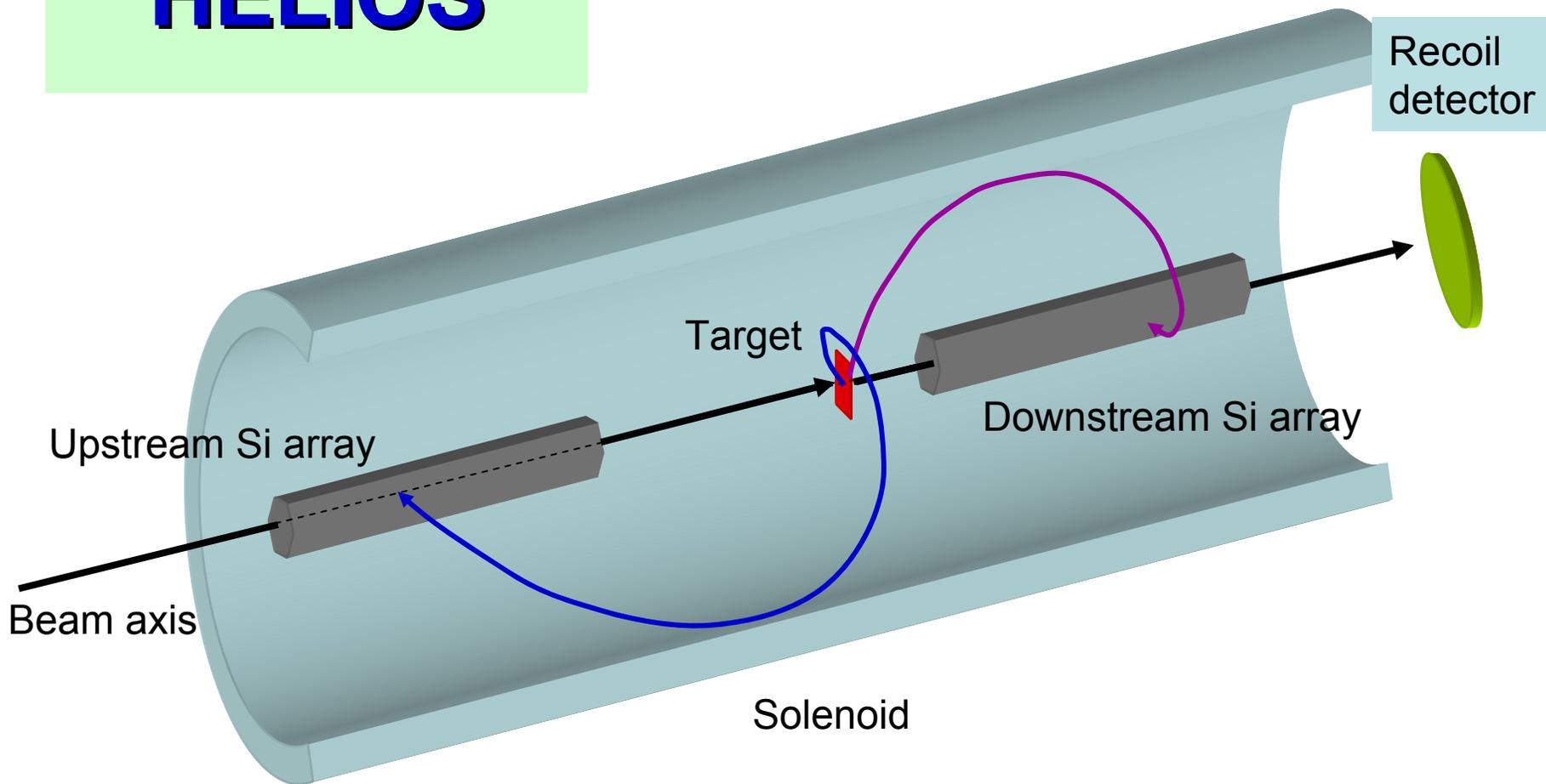
*E Monitor @  $0^\circ$*   
*X1800 Attenuator*

*E vs  $\Delta E(IC)$*

# Looking ahead ...



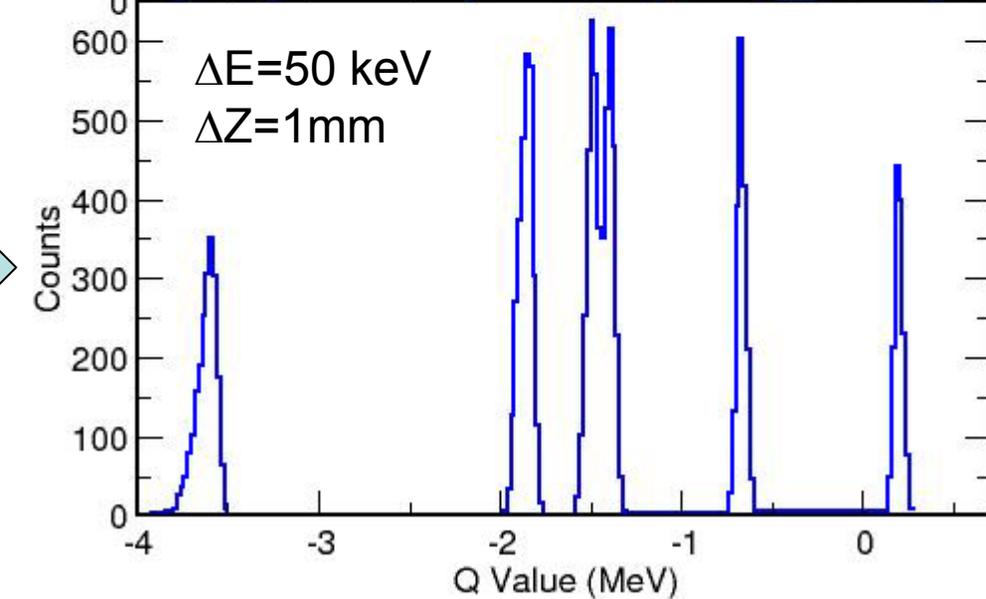
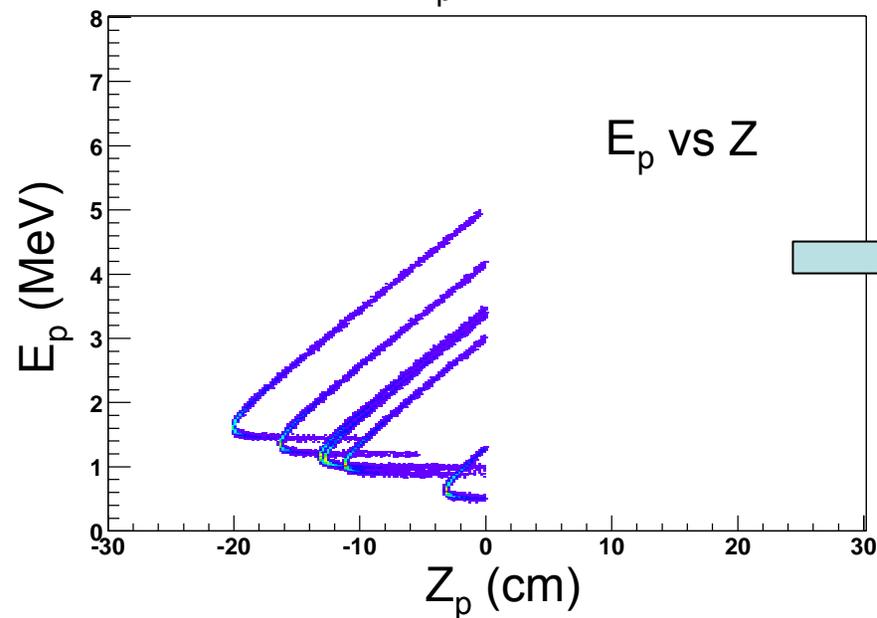
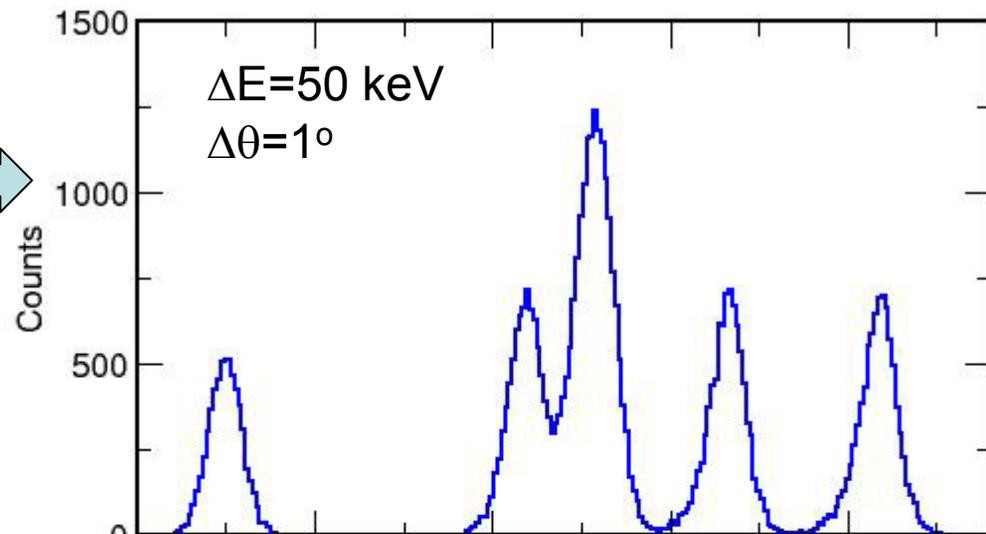
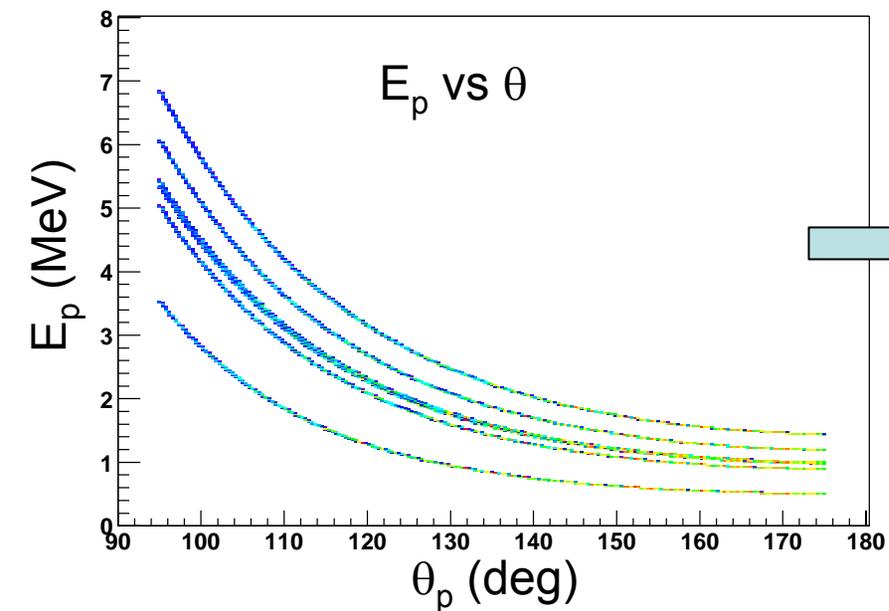
# HELIOS



Schematic  
design

# Improved resolution for (d,p)

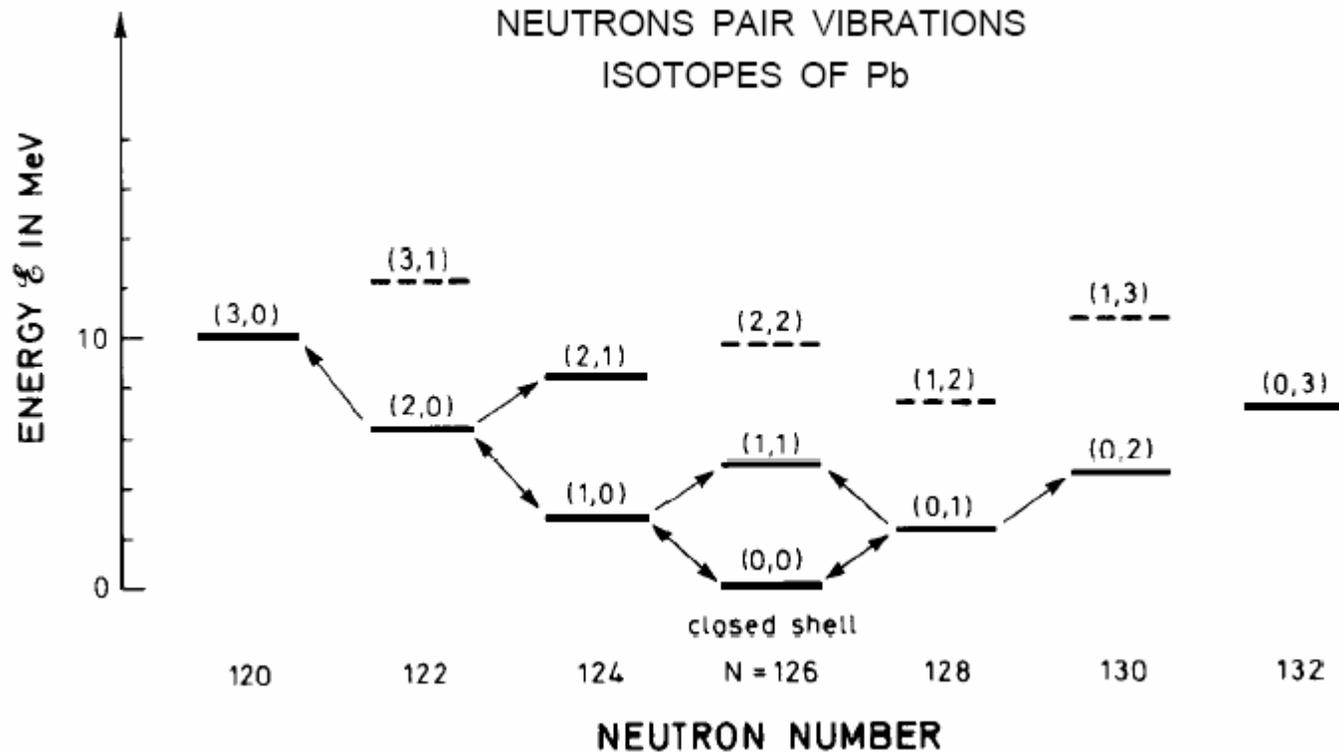
Q Value



# Pairing Phase Transition

# Pair-Vibrational Structures

(Nobel Lecture, Ben R. Mottelson, 1975 “Elementary Modes of Excitation in the Nucleus”)

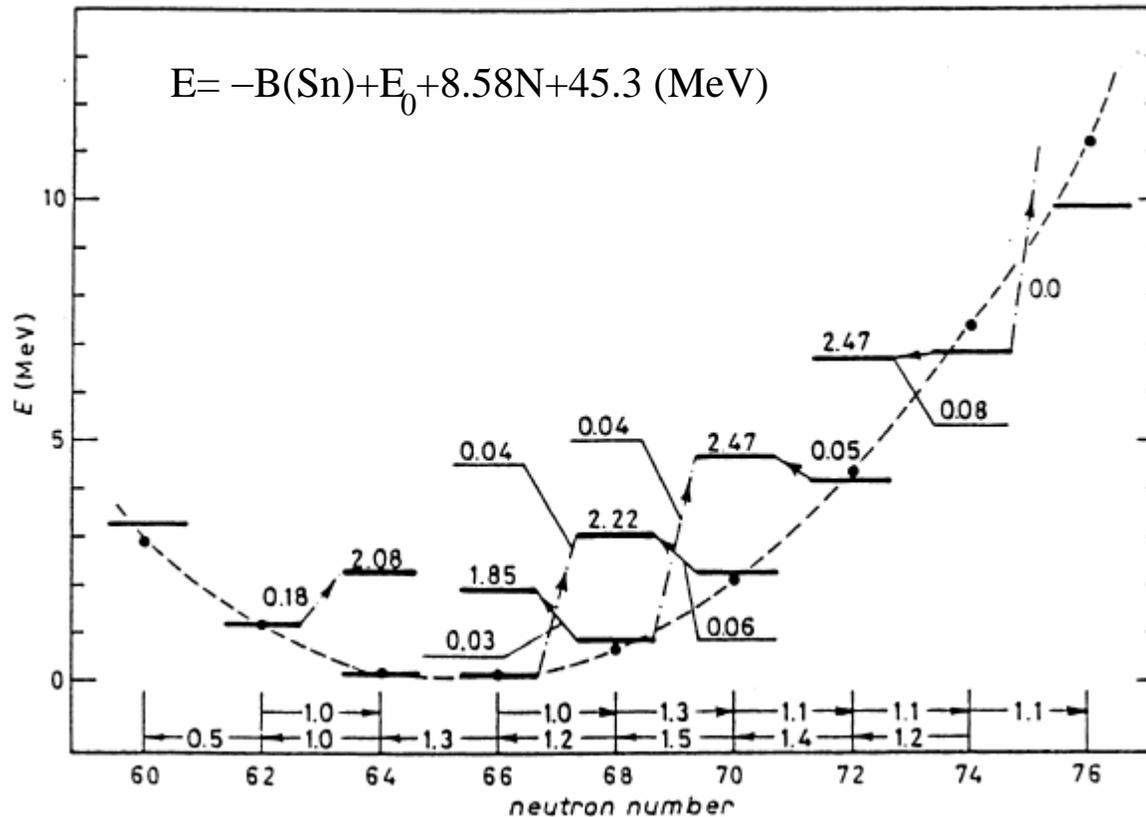


- Near closed shell nuclei (like  $^{208}\text{Pb}$ ) no static deformation of pair field.
- Corresponds to the “normal” nuclear limit.
- Fluctuations give rise to a vibrational-like excitation spectrum.
- Enhanced pair-addition and pair-removal cross-sections seen in (t,p) and (p,t) reactions (indicated by arrows).

⇒ • Large anharmonicities in spectrum must be accounted for.

# Pair-Rotational Structures

(R. A. Broglia, J. Terasaki, and N. Giovanardi, Phys. Rep. 335 (2000) 1)



- Many like-nucleon pairs outside a closed-shell configuration (e.g.  $^{116}\text{Sn}$ ) gives rise to a static deformation of the pair field.
- Corresponds to the “superconducting” limit.
- Rotational-like (parabolic – dashed line) spectrum formed by sequence of ground states of even- $N$  neighbors.
- Angular variable in rotational motion is gauge angle,  $\phi$ .

# Critical-Point Descriptions of Shape Transitions

F. Iachello, Phys. Rev. Lett. 85 (2000) 3580, Phys. Rev. Lett. 87 (2001) 052502.

PRL 96, 032501 (2006)

PHYSICAL REVIEW LETTERS

week ending  
27 JANUARY 2006

## Critical-Point Description of the Transition from Vibrational to Rotational Regimes in the Pairing Phase

R. M. Clark,<sup>1</sup> A. O. Macchiavelli,<sup>1</sup> L. Fortunato,<sup>2</sup> and R. Krücken<sup>3</sup>

<sup>1</sup>Nuclear Science Division, Lawrence Berkeley National Laboratory, Berkeley, California 94720, USA

<sup>2</sup>Dipartimento di Fisica Galileo Galilei, INFN, Sezione di Padova, Padova, Italy

<sup>3</sup>Physik Department E12, Technische Universität München, Garching, Germany

## The Collective Pairing Hamiltonian

D.R. Bès, R.A. Broglia, R.P.J. Perazzo, K. Kumar, Nucl. Phys. A 143 (1970) 1.

$$-\frac{\hbar^2}{2B} \frac{\partial^2 \psi}{\partial \alpha^2} - \frac{\hbar^2}{2B\alpha} \frac{\partial \psi}{\partial \alpha} + \left( \frac{\hbar^2 M^2}{8B\alpha} + V(\alpha) - E \right) \psi = 0$$

$\alpha$  is the deformation of the pair field (can be related to the gap parameter,  $\Delta$ ).

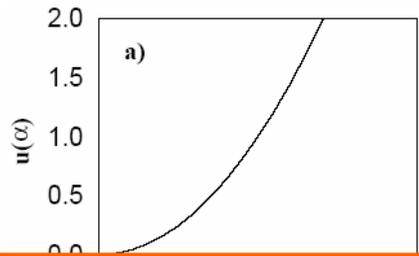
$B$  is a mass parameter.  $M=(A-A_0)$  (number of particles,  $A$ , relative to reference,  $A_0$ ).

$V(\alpha)$  is the potential.

## Analogy between Shapes and Pairing

$$\frac{\partial^2 \psi}{\partial \alpha^2} + \frac{1}{\alpha} \frac{\partial \psi}{\partial \alpha} + \left( \varepsilon - u(\alpha) - \frac{M^2}{4\alpha^2} \right) \psi = 0$$

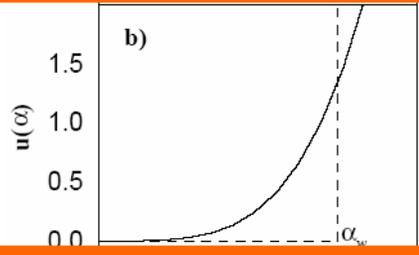
Vibrational Limit  
("Normal")



$$u(\alpha) = \alpha^2$$

**E(2)**

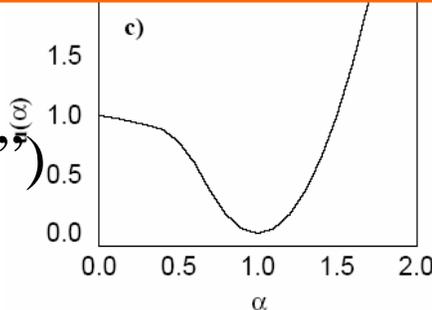
Transitional  
("Critical-Point")



$$u(\alpha) = 0, \alpha \leq \alpha_w$$

$$u(\alpha) = \infty, \alpha > \alpha_w$$

Rotational Limit  
("Superconducting")



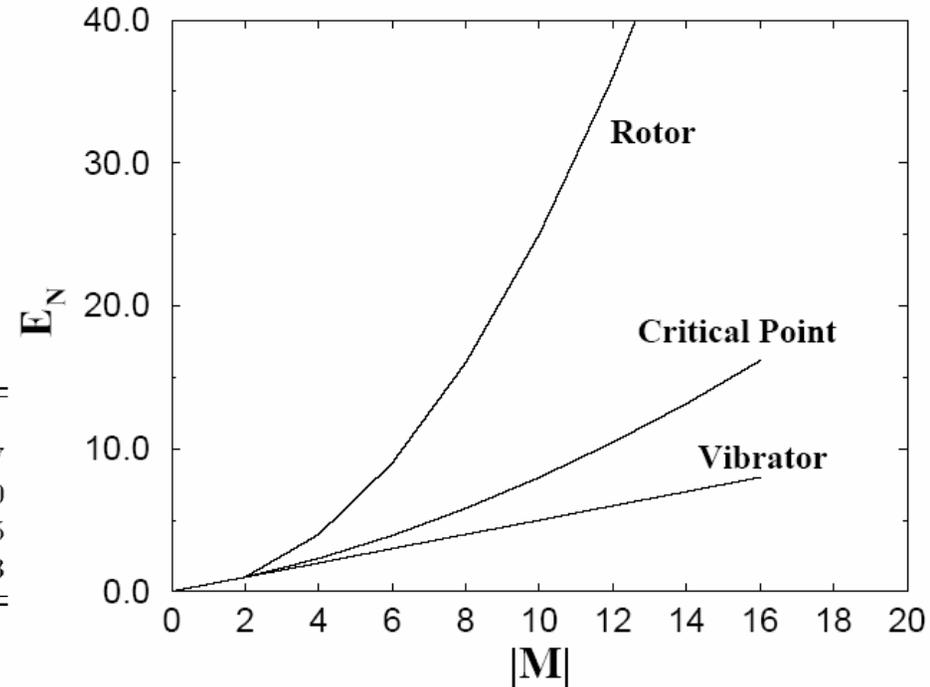
$$u(\alpha) = \delta_{\alpha_0}$$

# The Reduced Energy Spectrum

Normalizing energies of excited states to that of the first excited state:

$$E_N = \frac{x_{\xi, M}^2 - x_{1,0}^2}{x_{1,2}^2 - x_{1,0}^2}$$

	$\xi=1$	$\xi=2$	$\xi=3$	$\xi=4$
$ M =0$	0.00	2.77	7.77	14.97
$ M =2$	1.00	4.88	10.98	19.30
$ M =4$	2.31	7.31	14.52	23.95
$ M =6$	3.92	10.06	18.39	28.93

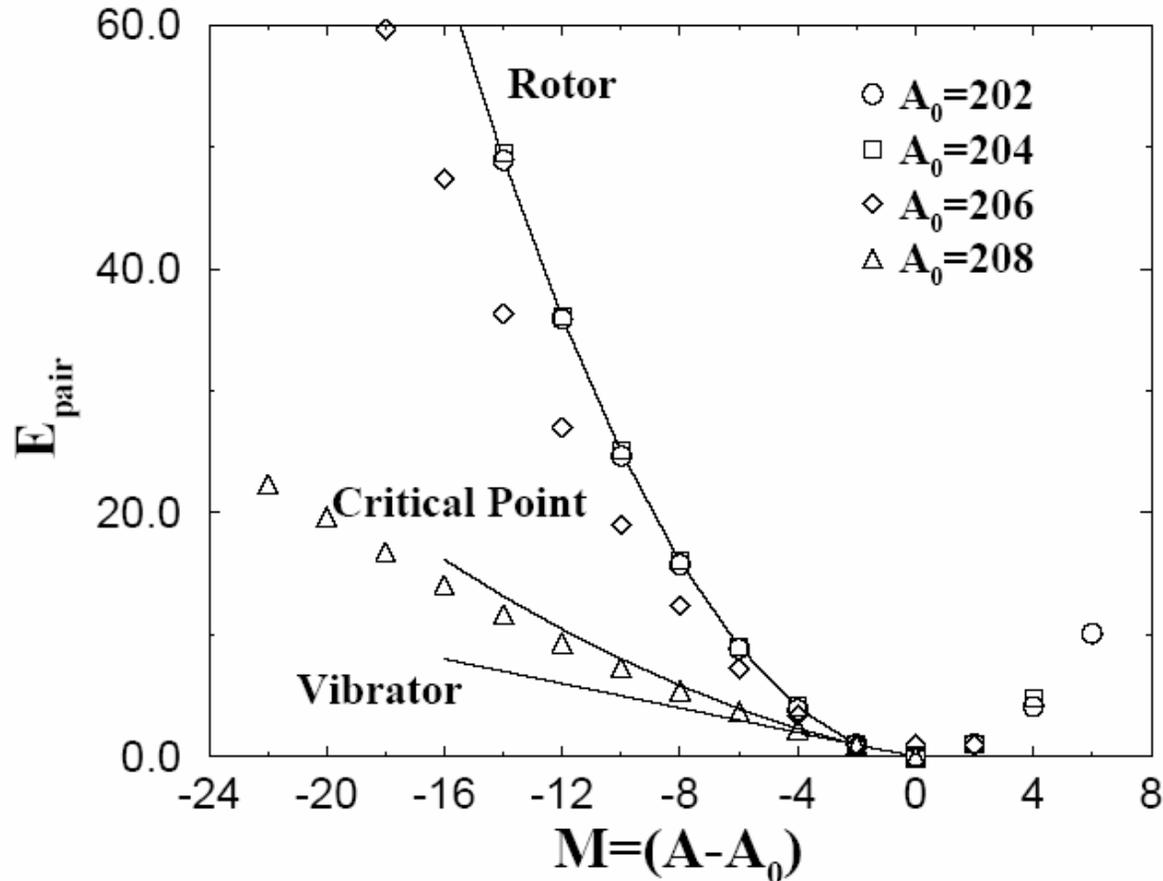


- $\xi=1$  sequence of states correspond empirically to the sequence formed by the  $0^+$  ground-states of neighboring even-even nuclei along isotopic or isotonic chain.
- $\xi=1$  sequence of states follows behavior between the linear dependence of a harmonic vibrator and the parabolic dependence of a deformed rotor, as expected.
- $\xi>1$  correspond to excited  $0^+$  states formed from pair excitations.

## Comparison with Data

Empirical neutron pairing energy defined as:  $E_{pair} = [\varepsilon(A) - \varepsilon(A_0)] - C \cdot (A - A_0)$

$\varepsilon(A) - \varepsilon(A_0)$  is difference in mass excess between isotope of mass A and reference nucleus of mass  $A_0$  (G. Audi et al., Nucl. Phys. A 729 (2003) 337)



Only a few nucleons outside of the closed shell are required for a static pair deformation (“superconductivity”) and pair rotational sequences develop.

## A comment on anharmonicities

From Bohr & Mottelson Vol. II pag. 646

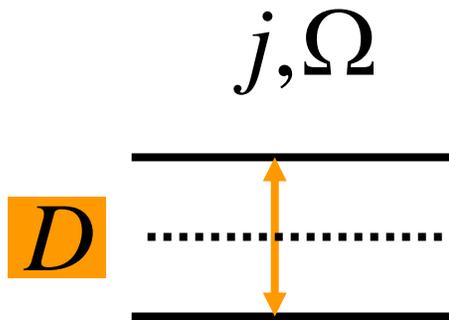
$$E \sim n\hbar\omega_0 + \frac{1}{2}V_{--}n(n-1)$$

$$V_{--} \sim 700\text{keV}$$

Our simple estimate is

$$V_{--} \sim (2.31 - 2)\hbar\omega_0 \sim 800\text{keV}$$

## A simple microscopic model: Two j-shells



$$H = \frac{1}{2} D(N_2 - N_1) - G\Omega(A_1^+ + A_2^+)(A_1 + A_2)$$

$$x = 2\Omega G / D$$

“Control parameter”

$$\Omega \sim A^{2/3}$$

$$G \sim 20/A$$

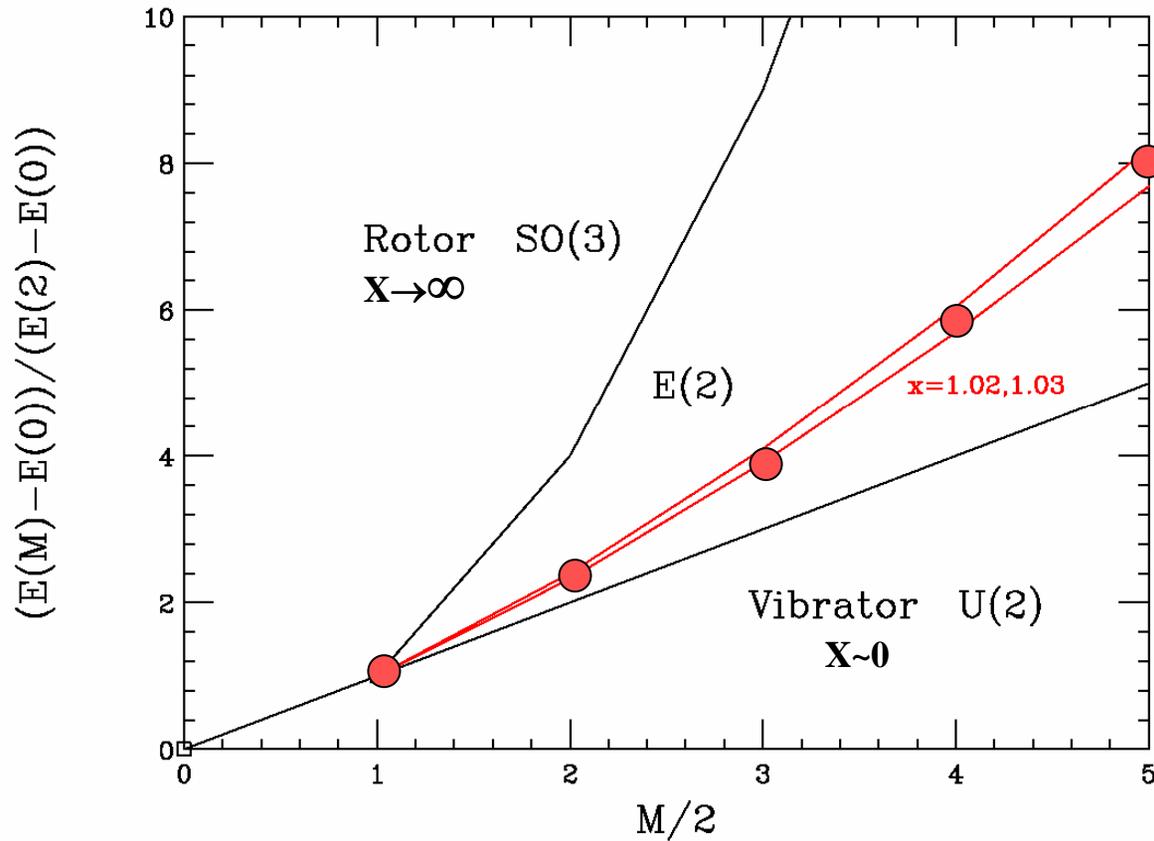
$$D \sim 41/A^{1/3}$$



$$x \sim 1$$

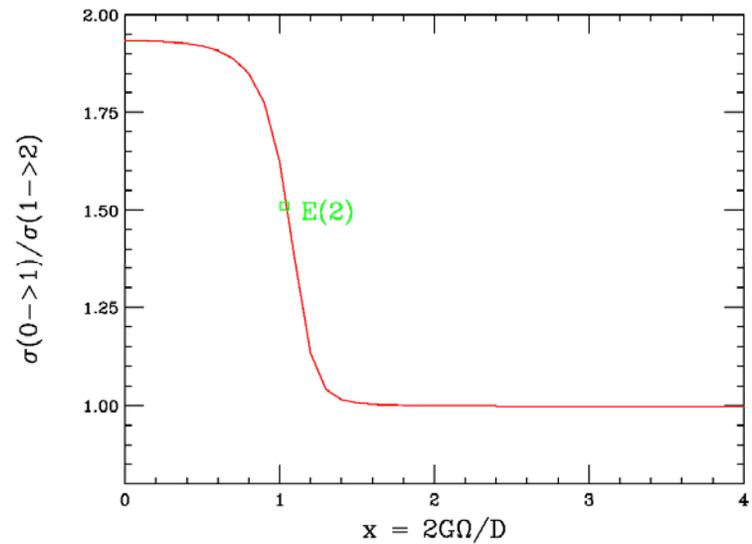
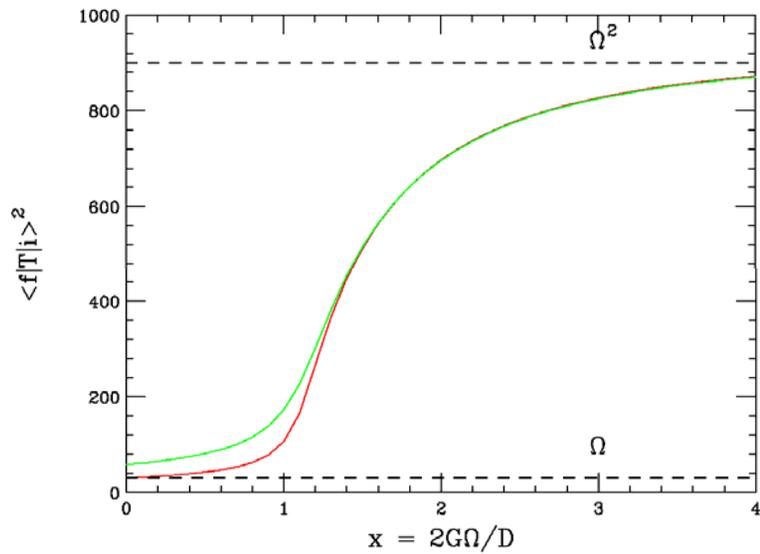
# A simple microscopic model: Two j-shells

## Energies



# A simple microscopic model: Two j-shells

## Transition probabilities



## Transition Probabilities

The transition matrix elements are related to two-nucleon transfer probabilities:

$$\langle \psi_{\xi', M'} | \hat{O} | \psi_{\xi, M} \rangle = \int \psi_{\xi', M'}^* \hat{O} \psi_{\xi, M} d\tau$$

where, the volume element is:  $d\tau = 2B\alpha \cdot d\alpha \cdot d\phi$

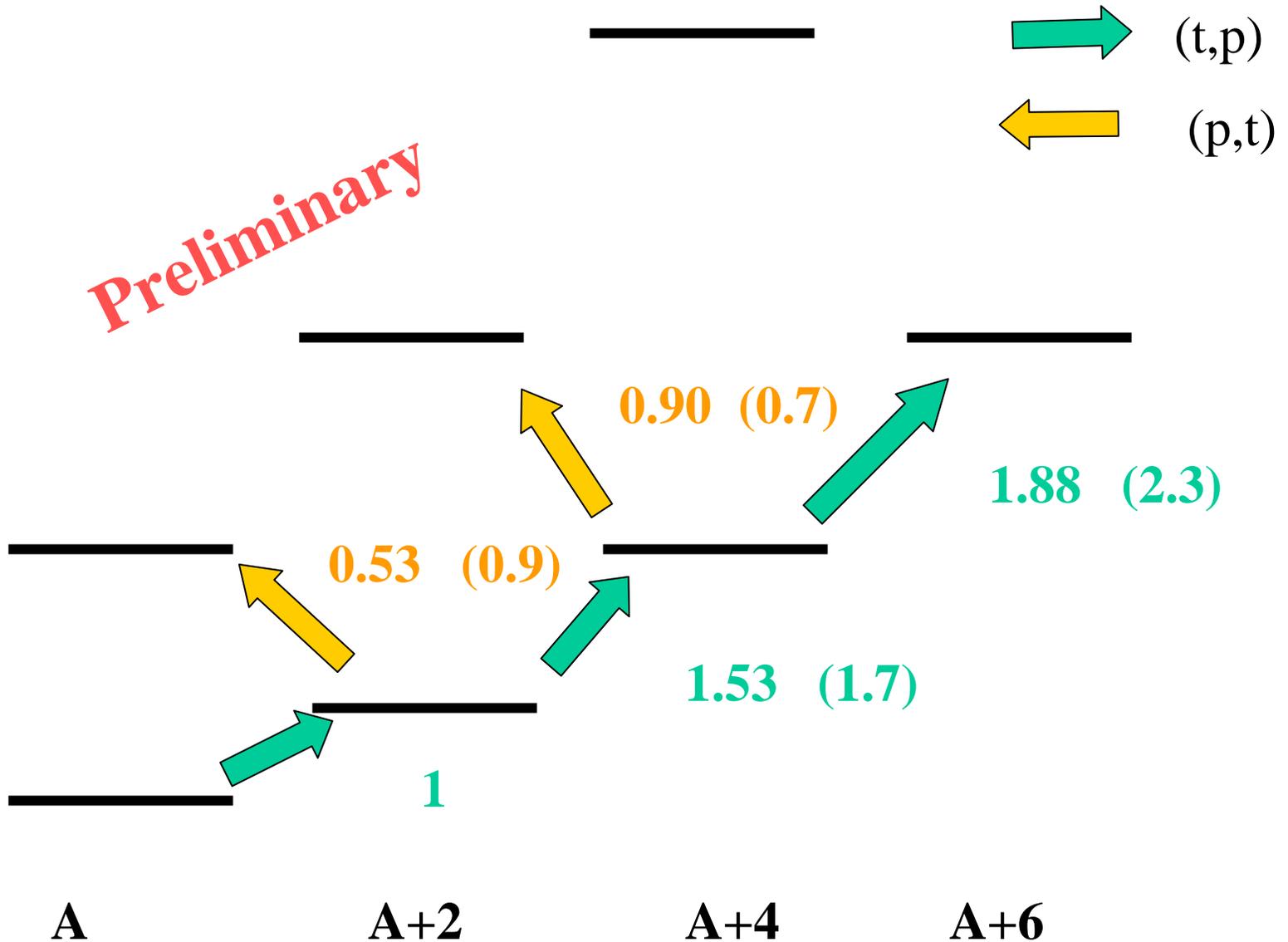
and the pair-transfer operator is:  $\hat{O} = \alpha_{\pm 2} = \alpha e^{\pm 2i\phi}$

One then finds:

$$\langle \psi_{\xi', M'} | \hat{O} | \psi_{\xi, M} \rangle \propto \int_0^{\alpha_w} \psi_{\xi', M'}^* \alpha^2 \psi_{\xi, M} d\alpha$$

The integrals can be solved numerically.

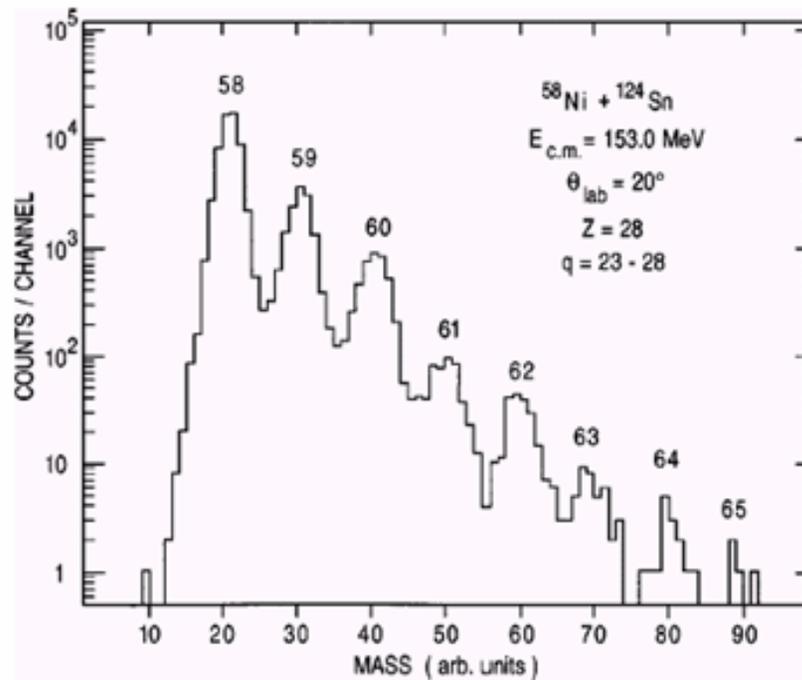
# Transition Probabilities



# Heavy Ion Transfer

# A Few Words on Heavy Ion Transfer Reactions

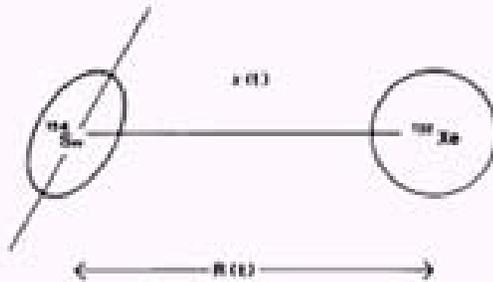
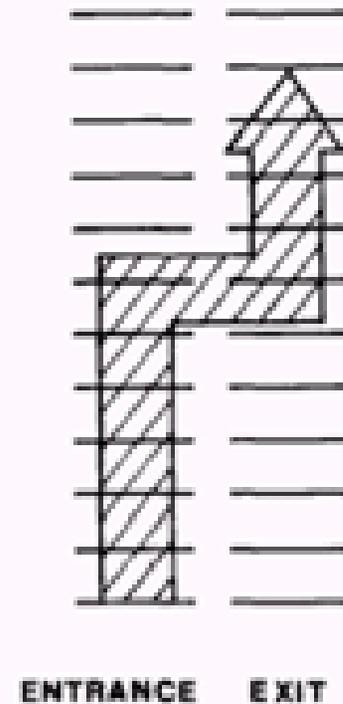
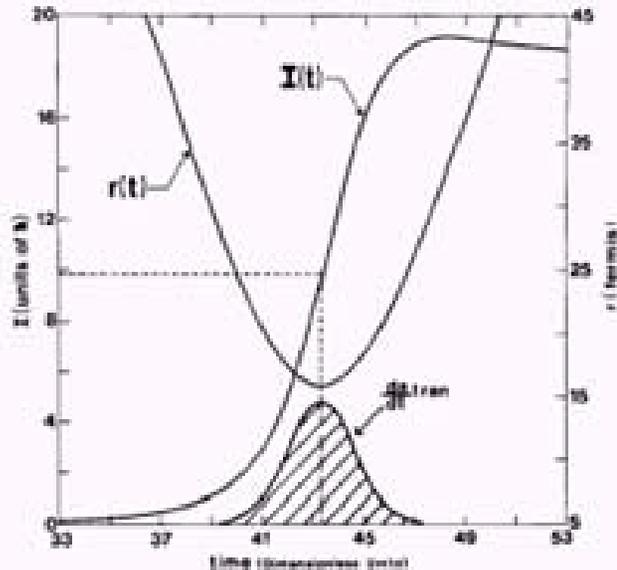
## Multiple nucleon transfer



*C.L.Jiang et al. PRC57(1998)*

$$\Delta M/M, \Delta L/L \text{ and } \Delta E/E \ll 1$$

Large Sommerfeld parameter  $\eta \gg 1$  allows for a semi-classical description.  $\Rightarrow$  Classical trajectory and tunneling



Population of high angular momentum and high excitation energy states

⇒ population mechanism, followed by  $\gamma$ -ray spectroscopy

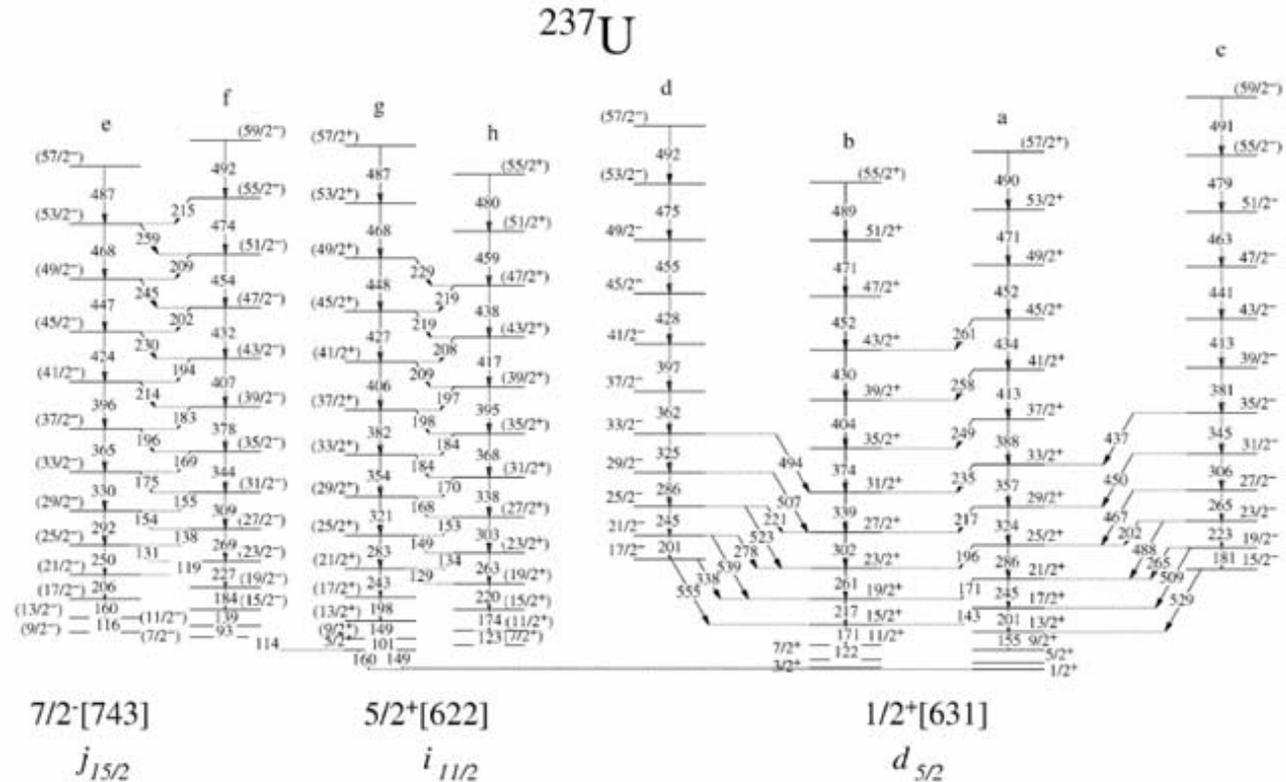


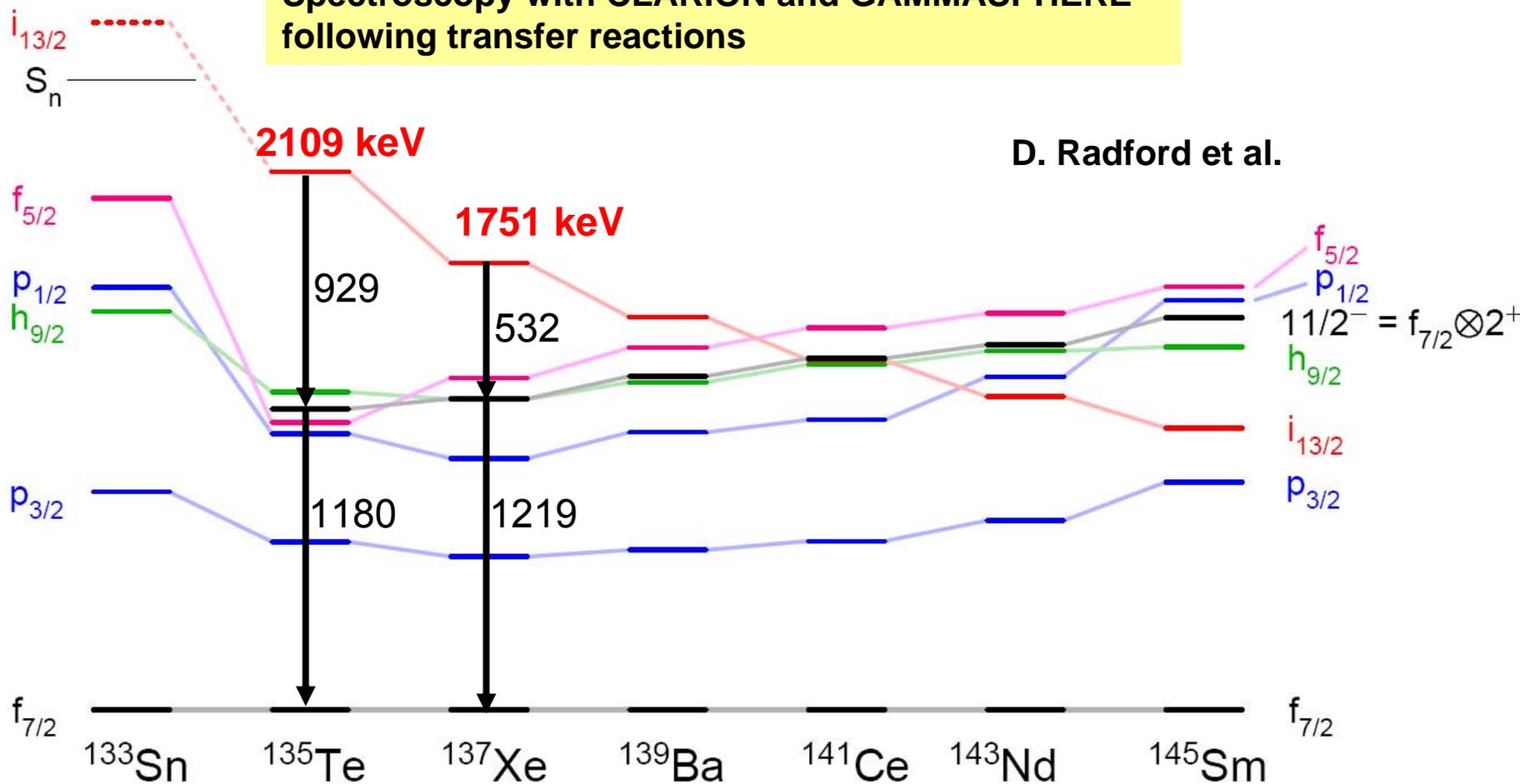
Fig. 3. Proposed level scheme for  $^{237}\text{U}$ . The energies of the transitions are given in keV. Bands c and d are the octupole bands under discussion in the present Letter.

One neutron pickup reaction,  $^{207}\text{Pb} + ^{238}\text{U}$  in GAMMASPHERE  
to study octupole correlations in actinides

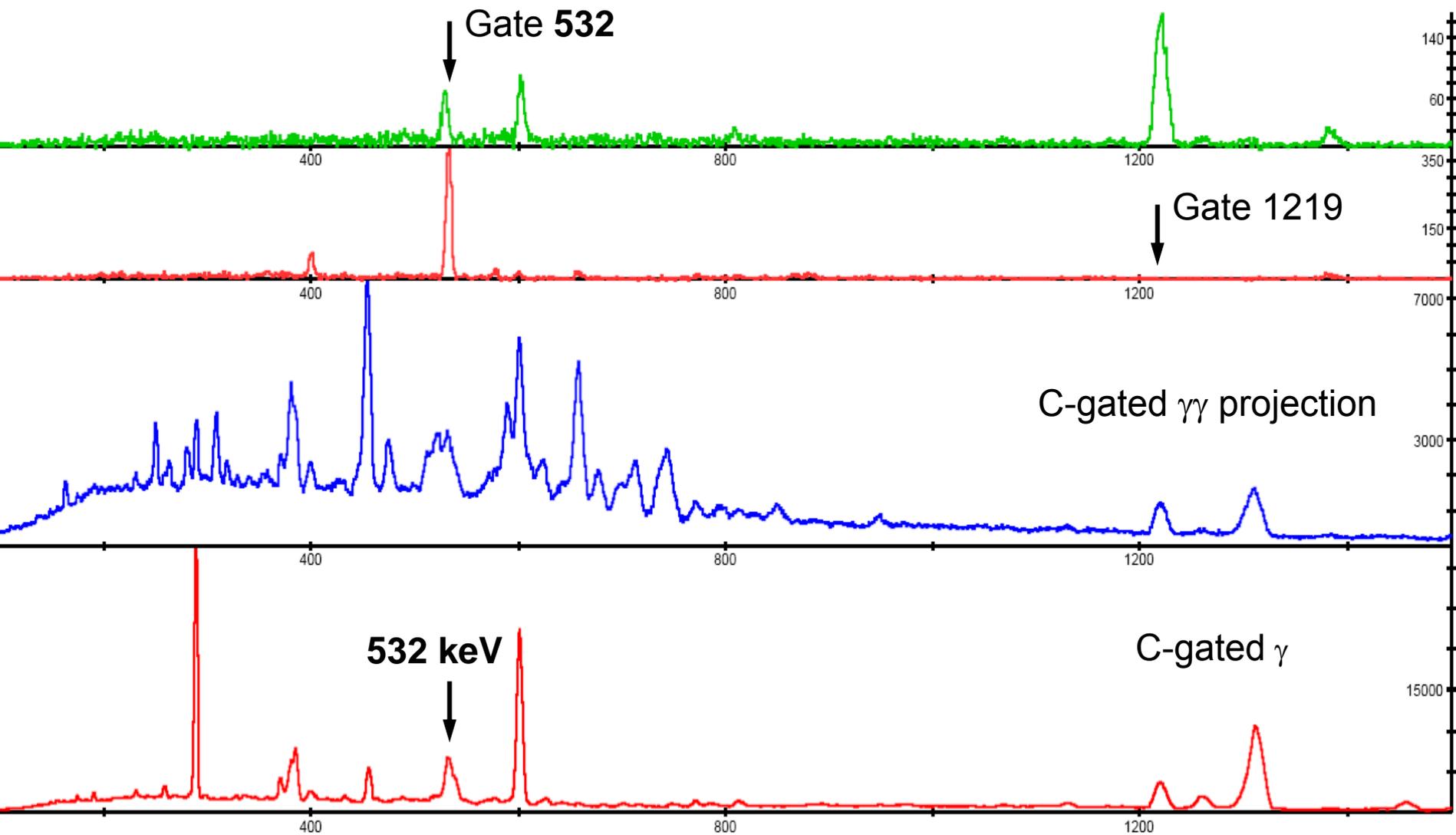
# Spin-orbit splitting again

## N = 83 level energy systematics

Spectroscopy with CLARION and GAMMASPHERE following transfer reactions



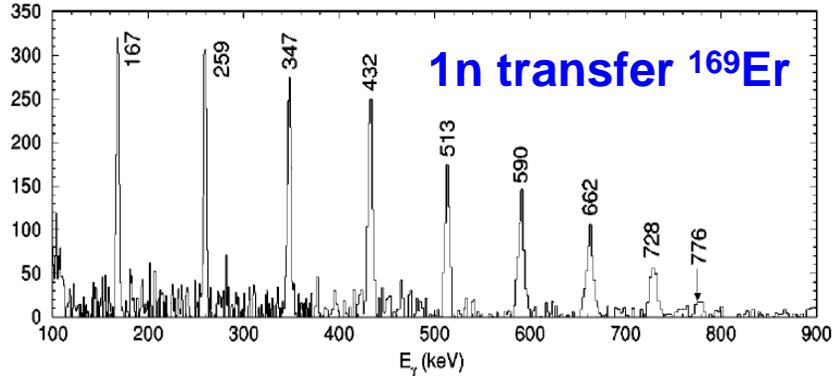
$^{13}\text{C}(^{136}\text{Xe}, ^{12}\text{C})^{137}\text{Xe}$  560 MeV



# GRETINA: $1\pi$ array

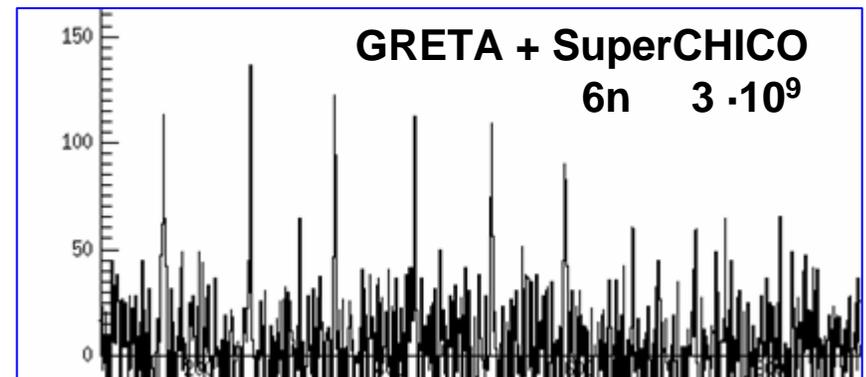
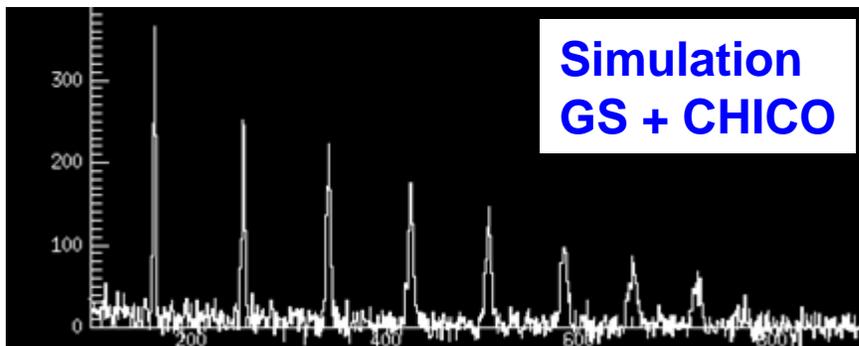
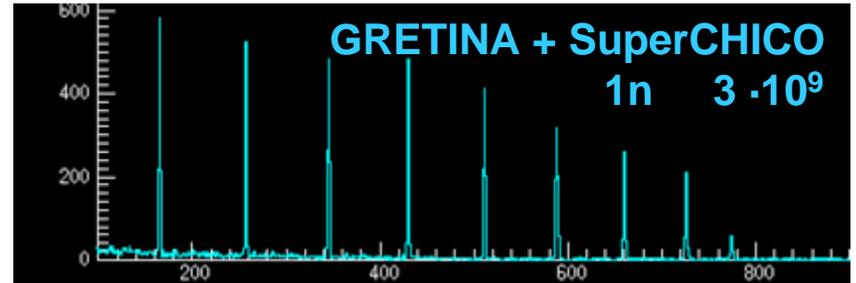
## Transfer reactions with re-accelerated RIBS and stable beams

$^{238}\text{U} + ^{170}\text{Er}$  5.7 MeV/u  
GS + CHICO  
 $3 \cdot 10^9$  p/s (0.5 pna), 0.5 mg/cm<sup>2</sup>  
3 days,  $\gamma$ - $\gamma$ - $\gamma$



C.Y. Wu et al., PRC 70, 014313 (2004)

Simulation  $^{170}\text{Er} + ^{238}\text{U}$  5.7 MeV/u



## Conclusions

**Transfer Reactions have provided a wealth of information that has shaped our current understanding of the structure of atomic nuclei. They will continue to provide a unique tool as we embark in our experimental study of very-neutron (proton) rich nuclei.**

**Existing and planned exotic beam facilities worldwide and new detector systems with increased sensitivity and resolving power not only will allow us to answer some burning questions we have today, but most likely will open up a window to new and unexpected phenomena.**

**As so eloquently expressed in the title of the conference, the nucleus and its structure are of paramount importance to many aspects of physics.**

**I believe the road ahead looks inviting!**

**Thank You!**